

Atmospheric Long-Lived Particle Searches

Pilar Coloma
IFIC, UV/CSIC

Based on:

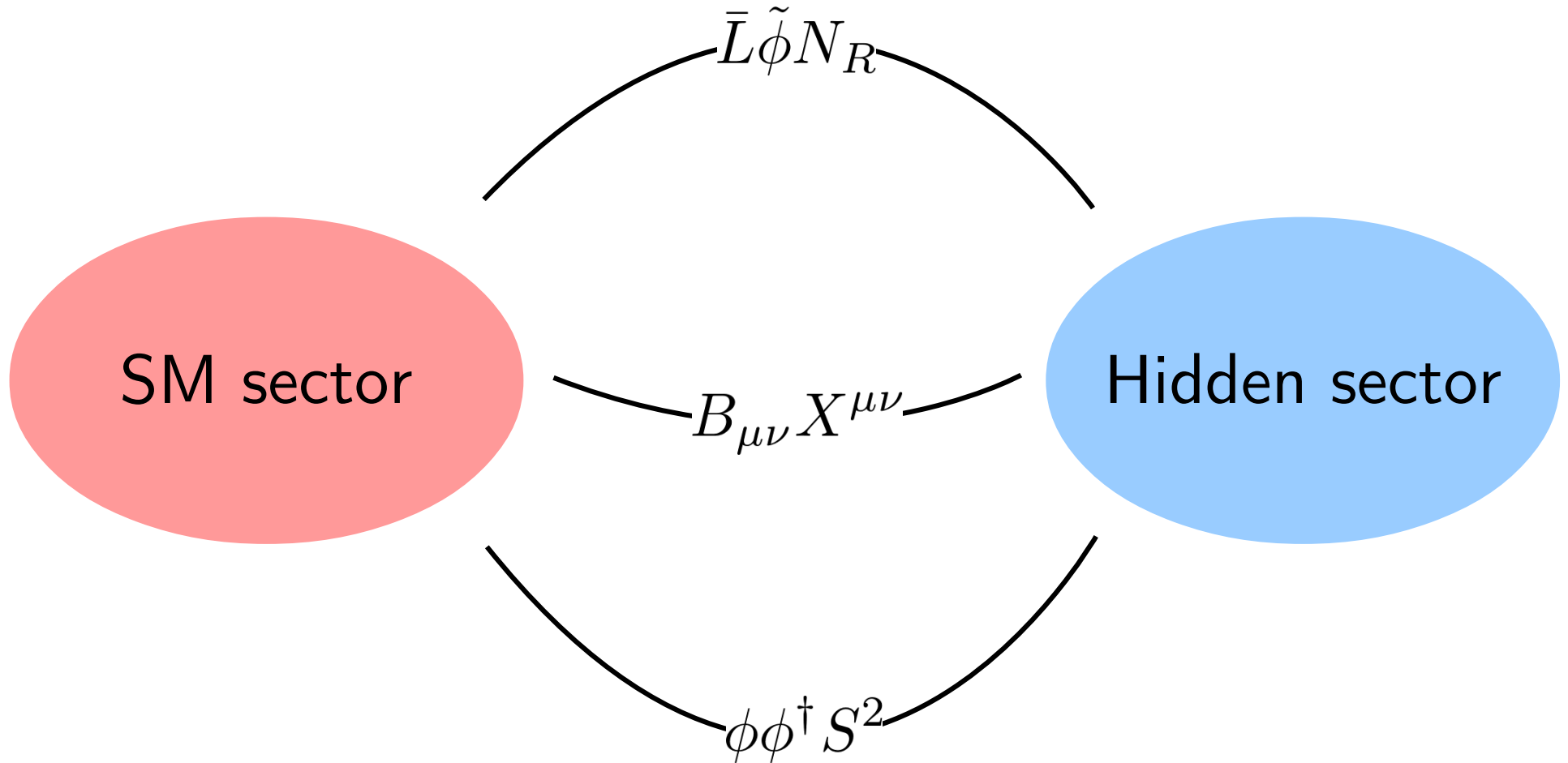
Argüelles, Coloma, Hernandez, Muñoz, 1910.12839
Coloma, Hernandez, Muñoz, Shoemaker, 1911.09129

Brookhaven Neutrino Theory Virtual Seminars
April 6th 2019

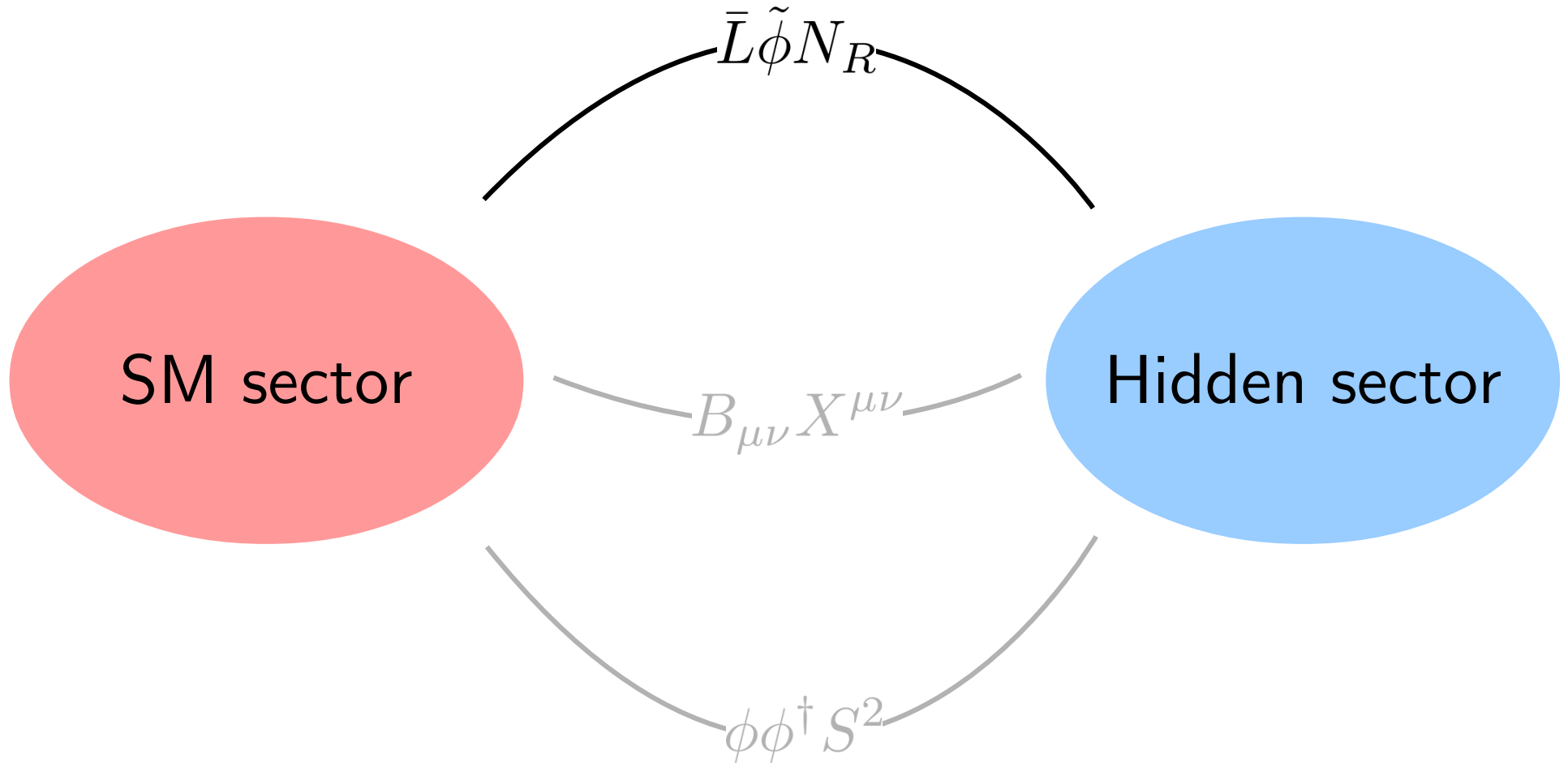
Outline

- 1) Introduction/motivation
- 2) Heavy neutral leptons
 - Sensitivity to general (non-minimal) models
 - Sensitivity to the minimal scenario
- 3) Dark photons, and a B-L model

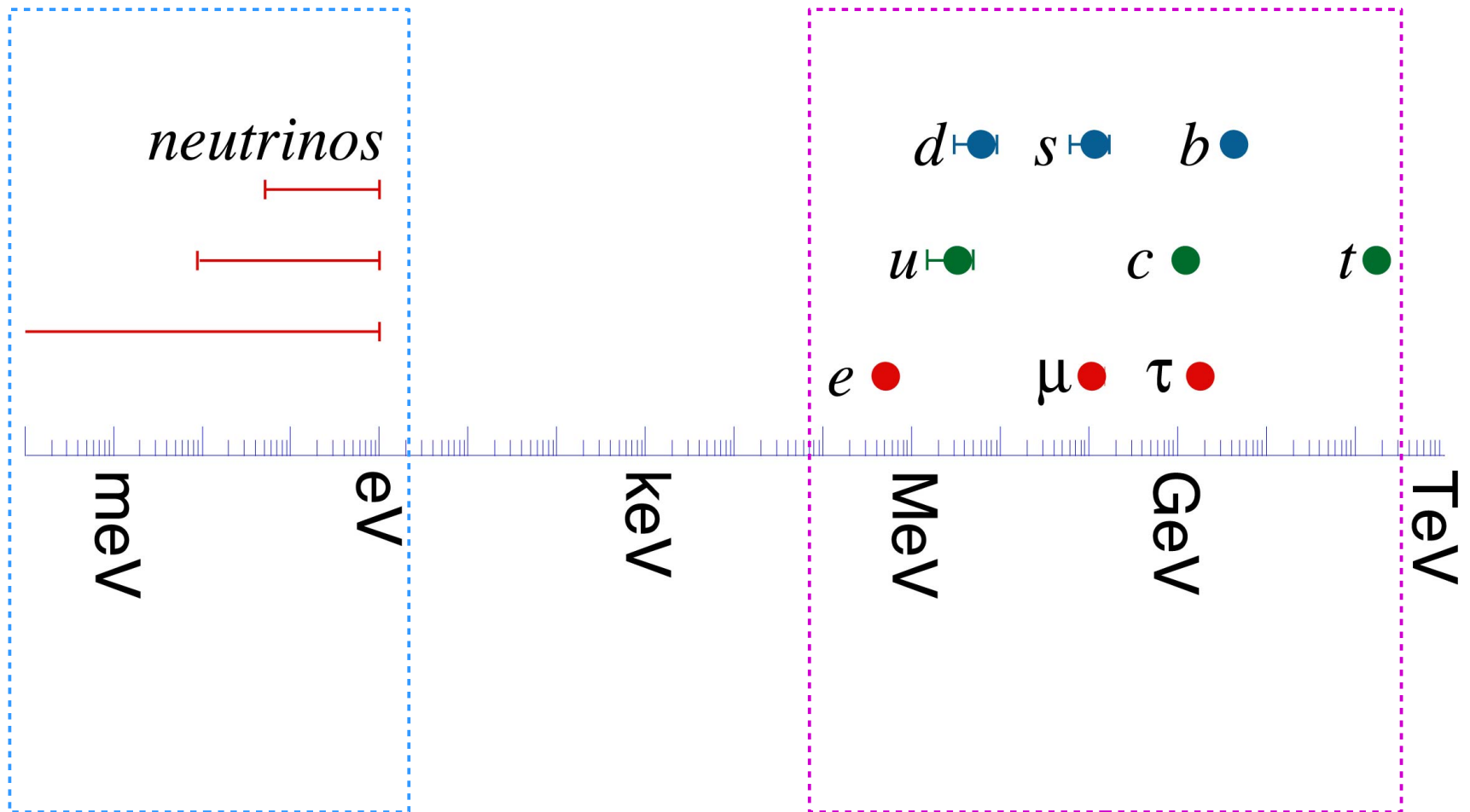
Where is the new physics?



Where is the new physics?



Neutrino masses



Majorana or Dirac?

New fields are required to give neutrinos a mass. Two main ways:

1) Dirac mass: as for the rest of fermions in the SM

$$Y_\nu \bar{L}_L \tilde{\phi} \nu_R \rightarrow m_\nu \bar{\nu}_L \nu_R$$

$$Y_\nu \lesssim 10^{-12}$$

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2) A Majorana mass. For example:

$$Y_\nu \bar{L}_L \tilde{\phi} \nu_R + \frac{1}{2} M \bar{\nu}_R^c \nu_R$$

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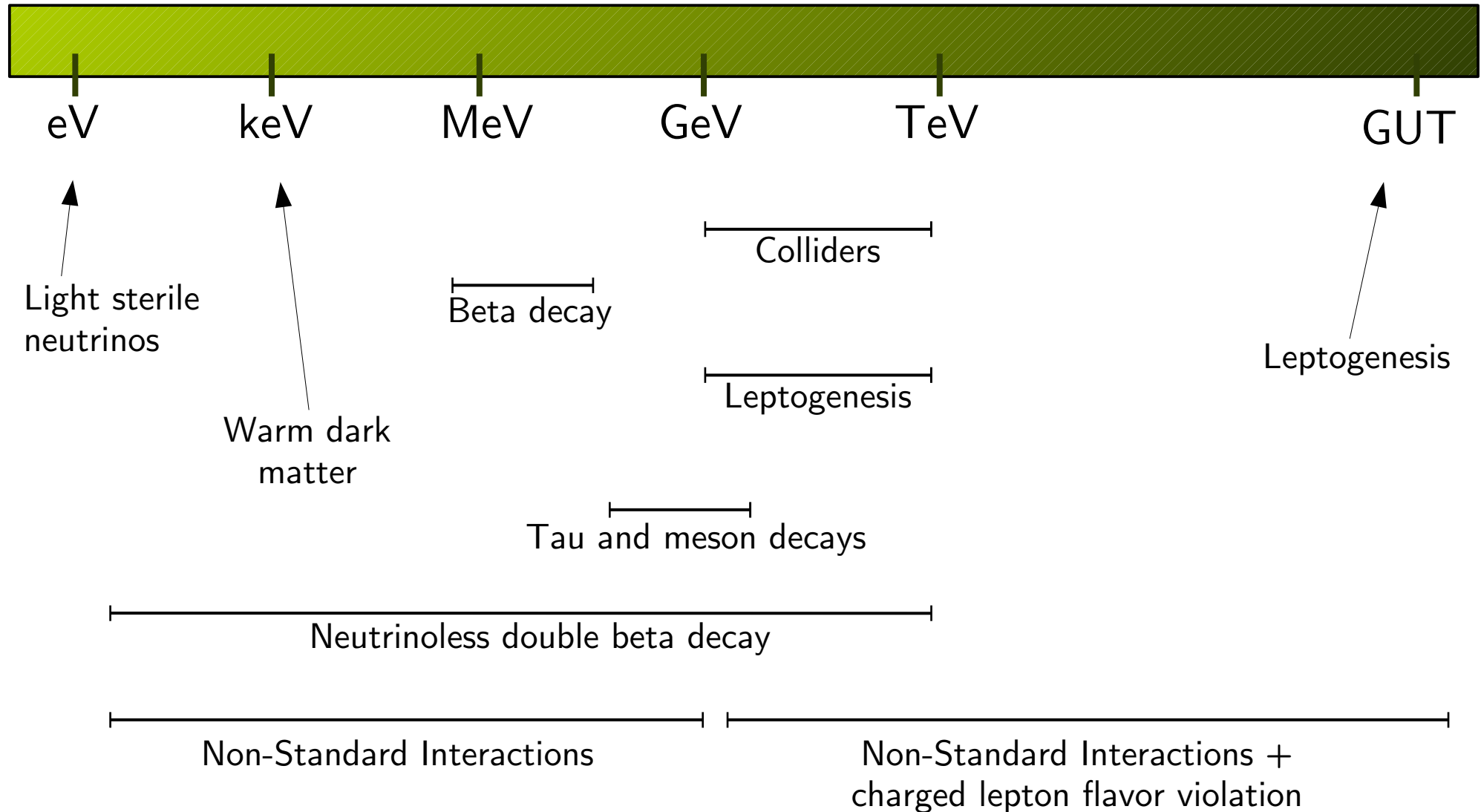
$$M \gg Y_\nu v$$

$$\rightarrow m_\nu = Y_\nu^\dagger M^{-1} Y_\nu v^2$$

Type I Seesaw:

Minkowski '77, Gell-Mann, Ramond, Slansky '79, Yanagida '79, Mohapatra, Senjanovic '80

Scale of new physics



Why the GeV scale?

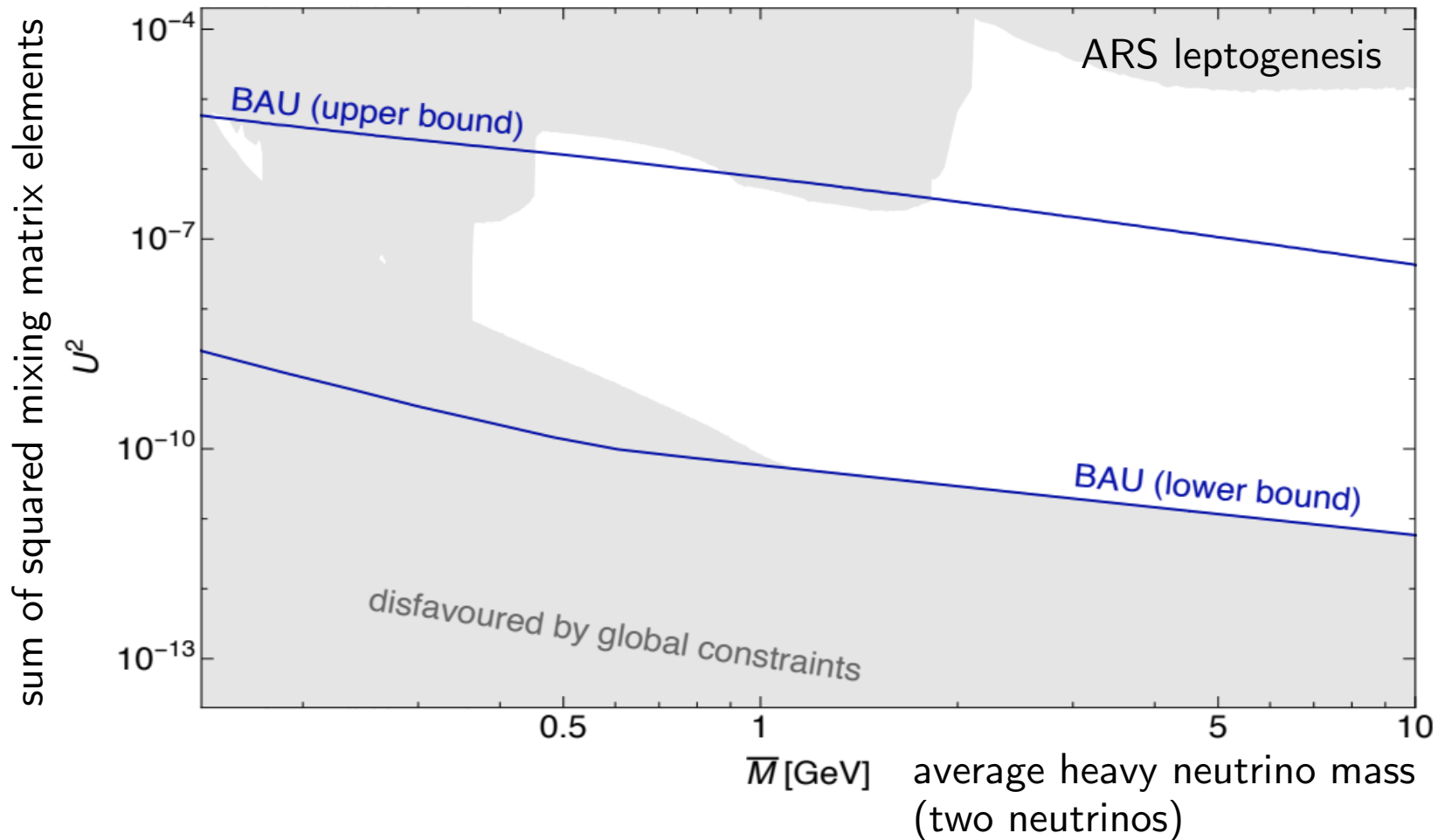
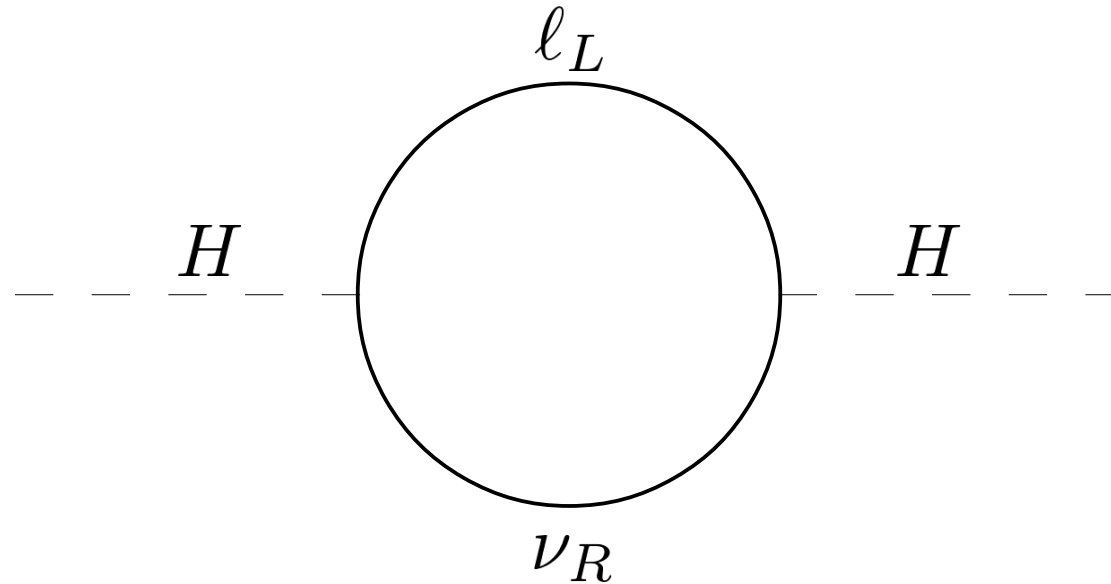


Figure from the review: Drewes, Garbretch, Hernandez, Kekic, Lopez-Pavon, Racker, Rius, Salvado, Teresi, 1711.02862

Why a low-scale seesaw?

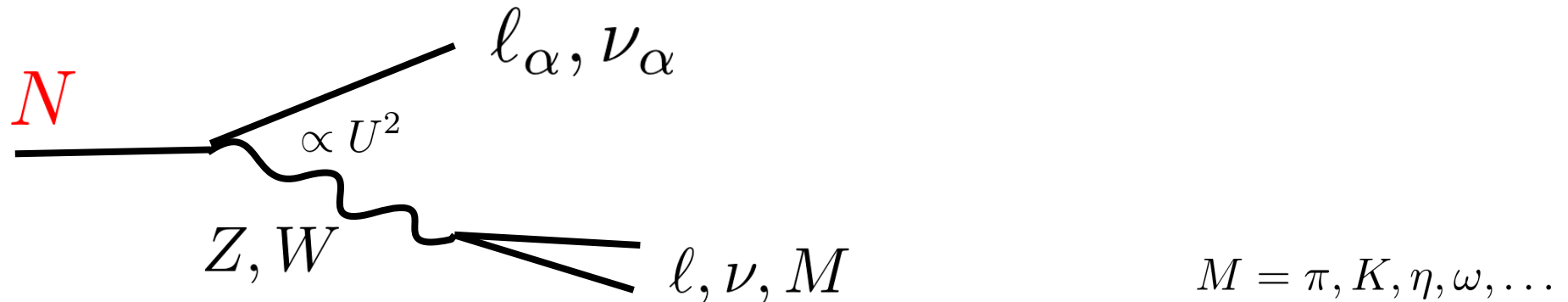


$$\delta\mu^2 \approx \frac{Y_\nu^2}{4\pi^2} M_R^2 \log \left(\frac{q}{M_R} \right)$$

Vissani, hep-ph/9709409

Casas, Espinosa, Hidalgo, hep-ph/0410298

Direct searches for GeV neutrinos



- Direct searches can be divided into two main categories:
 - Peak searches
 - Displaced decays: fixed target experiments, colliders, ...

$$c\tau \sim \text{few} \left(\frac{\text{GeV}}{m_N} \right)^5 \left(\frac{10^{-4}}{U^2} \right) \text{ m}$$

For detailed calculations of heavy neutrino decay channels see, e.g.:
 Ballett, Boschi, Pascoli, 1905.00284; Bondarenko et al, 1805.08567

Direct searches for GeV neutrinos

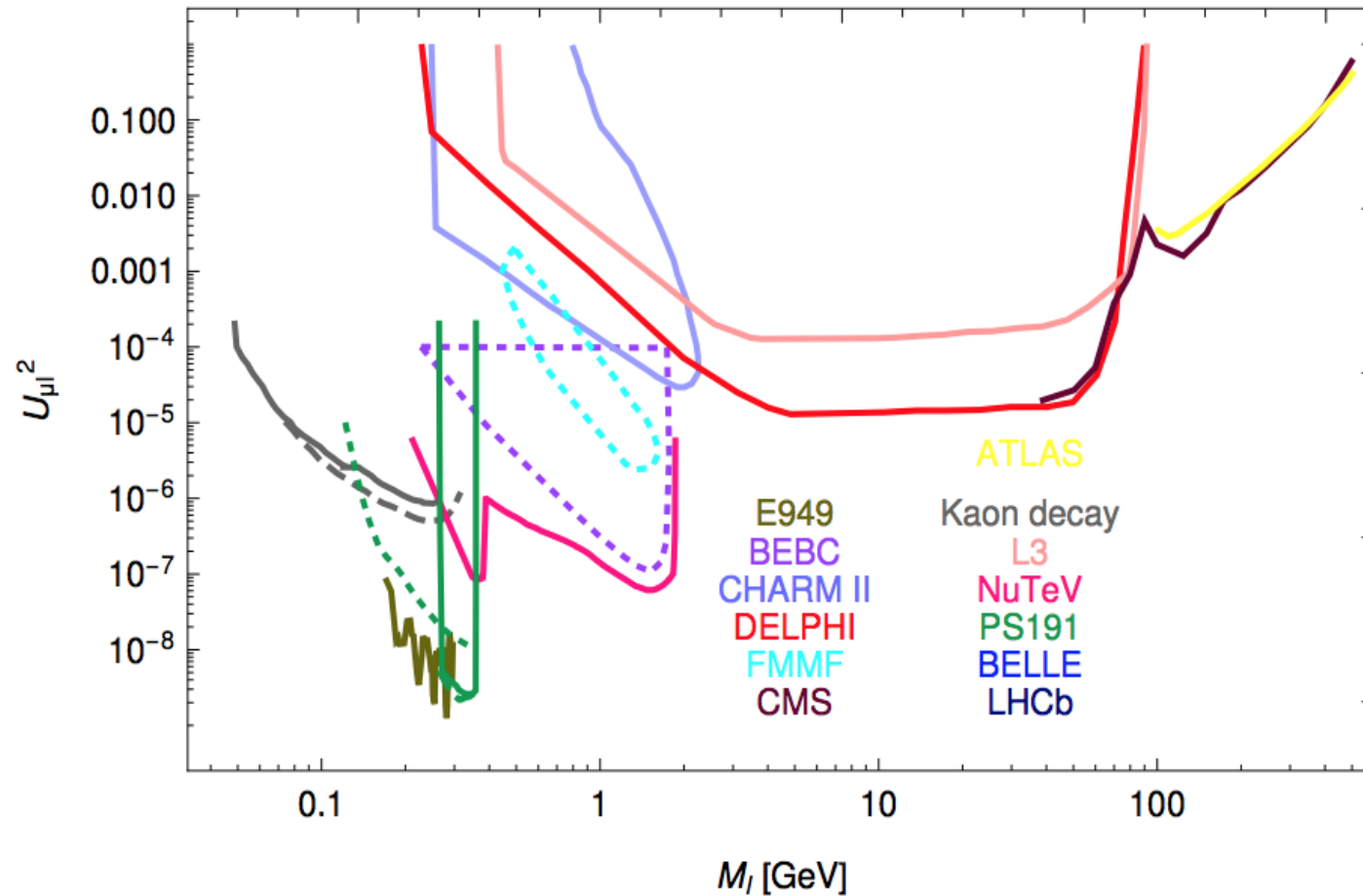
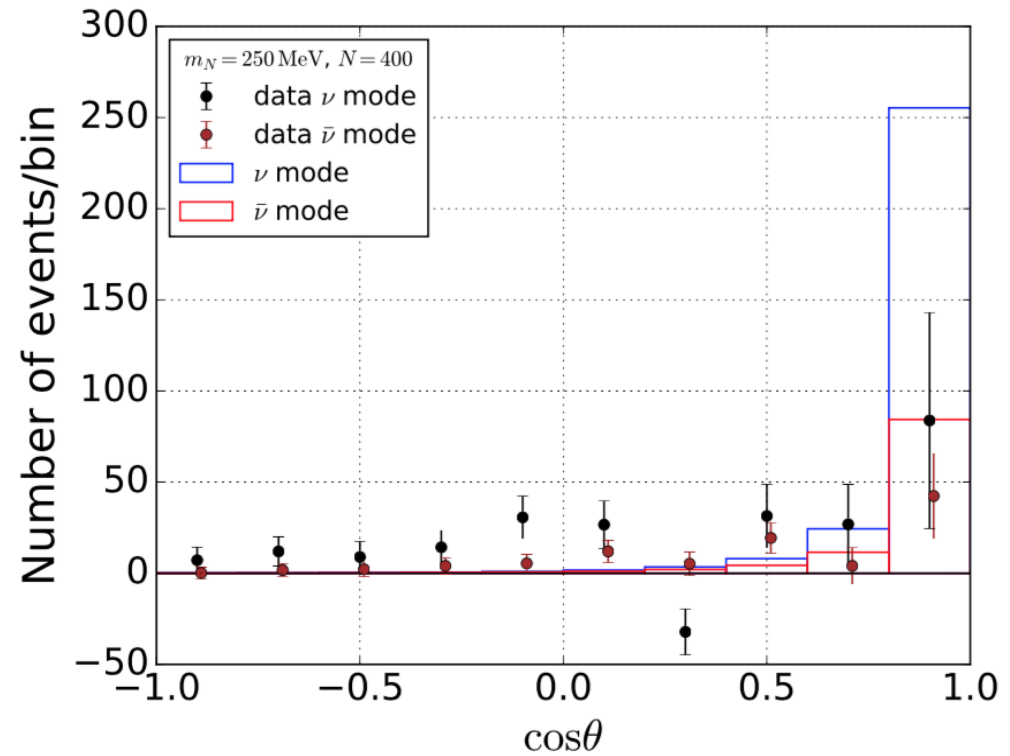
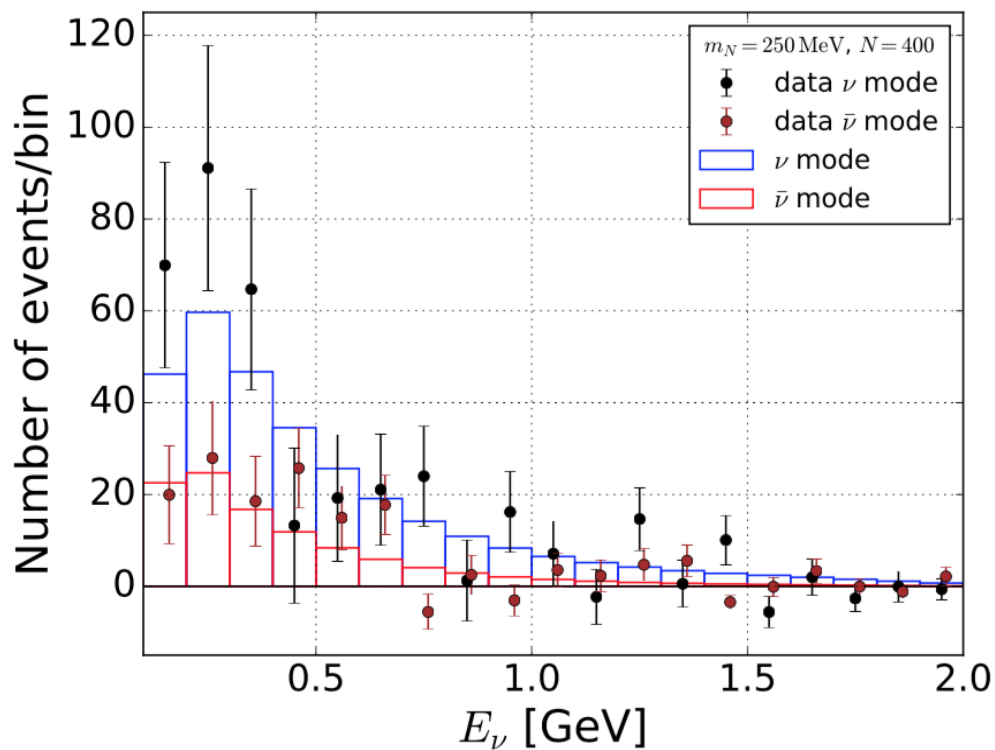


Figure from Drewes and Garbrecht, 1502.00477
 (See also Bryman and Shrock, 1904.06787 and 1909.11198,
 Ruchayskiy and Ivashko, 1112.3319, Atre et al, 0901.3589)

Further motivation...?

Figure from Fischer, Hernandez-Cabezudo, Schwetz, 1909.09561



See also:

Ballett, Pascoli, Ross-Lonergan, 1808.02915

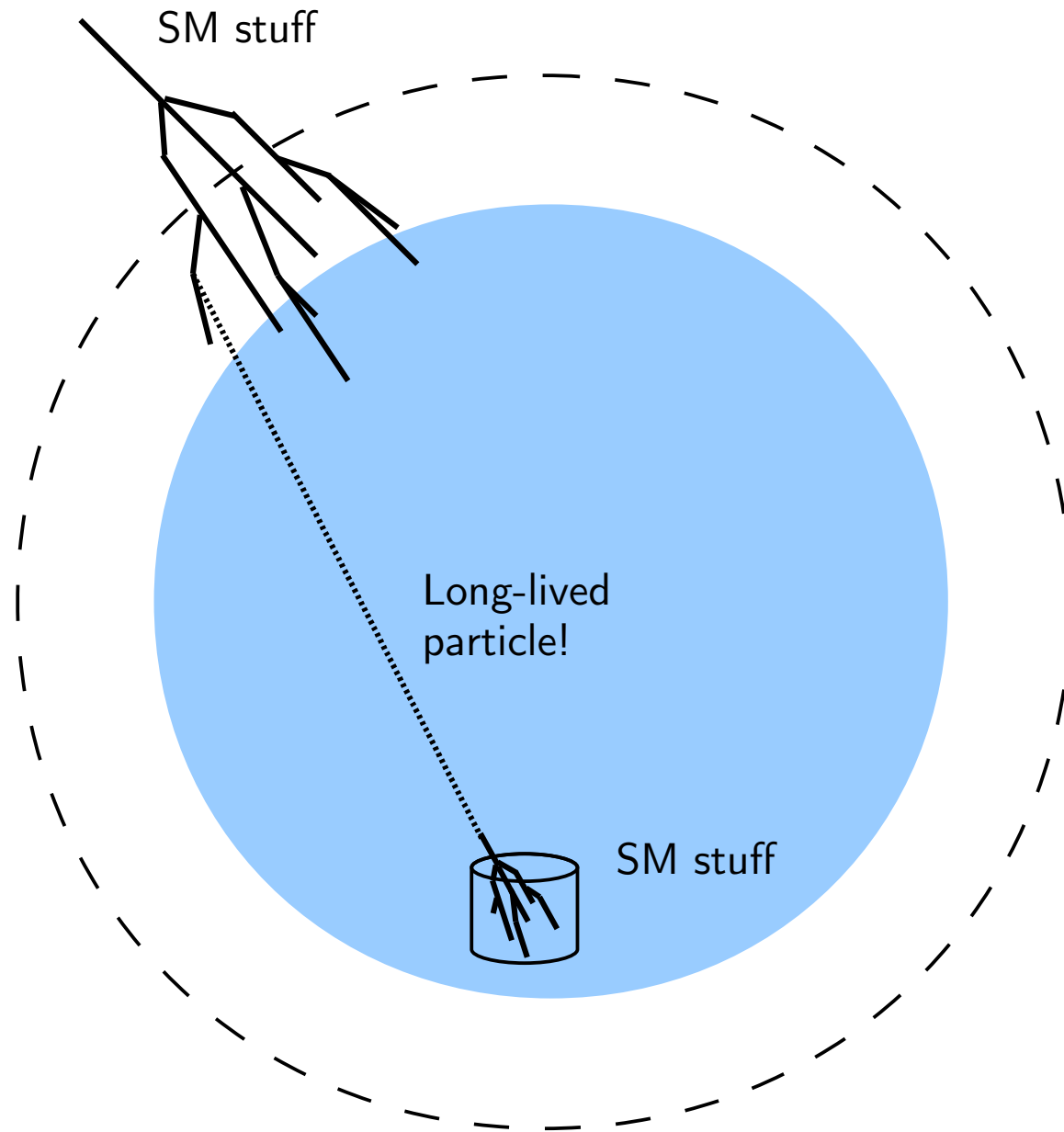
Bertuzzo, Jana, Machado, Zukanovich Funchal, 1807.09877

Dentler, Esteban, Kopp, Machado, 1911.01427

deGouvea, Peres, Prakash, Stenico, 1911.01447

Gninenko, 1009.5536 and 0902.3802

Palomares-Ruiz, Pascoli, Schwetz, hep-ph/0505216



Signal computation

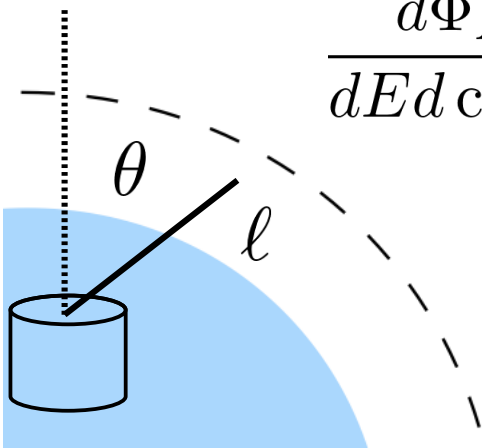
Production profile for the long-lived particle:

$$\frac{d\Pi_A}{dE d\cos\theta d\ell} = \sum_{ch} \int_{E_P^{\min}}^{E_P^{\max}} dE_P \frac{1}{\gamma_P \beta_P c \tau_P} \left[\frac{d\Phi_P(E_P, \cos\theta)}{dE_P d\cos\theta} \right] \frac{dn_{ch}(E_P, E)}{dE}$$

See, e.g., Gondolo, Ingleman and Thunman, hep-ph/9505417

From there, we get the flux as:

$$\frac{d\Phi_A}{dE d\cos\theta} = \int_0^{\ell_{\max}} d\ell \frac{d\Pi_A}{dE d\cos\theta d\ell} \left[e^{-\frac{\ell}{L_{\text{decay}}}} \right] \quad (A \equiv V, N)$$



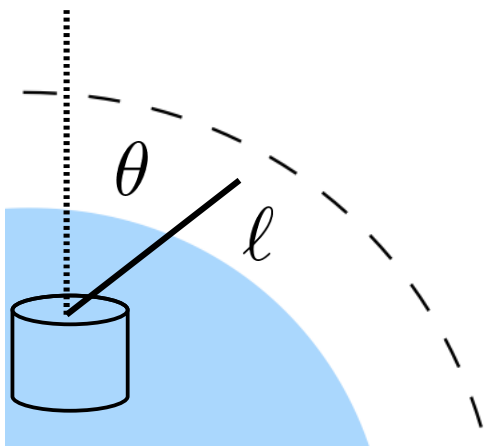
Parent meson fluxes extracted
from MCEq: <https://github.com/afedynitch/MCEq>

Signal computation

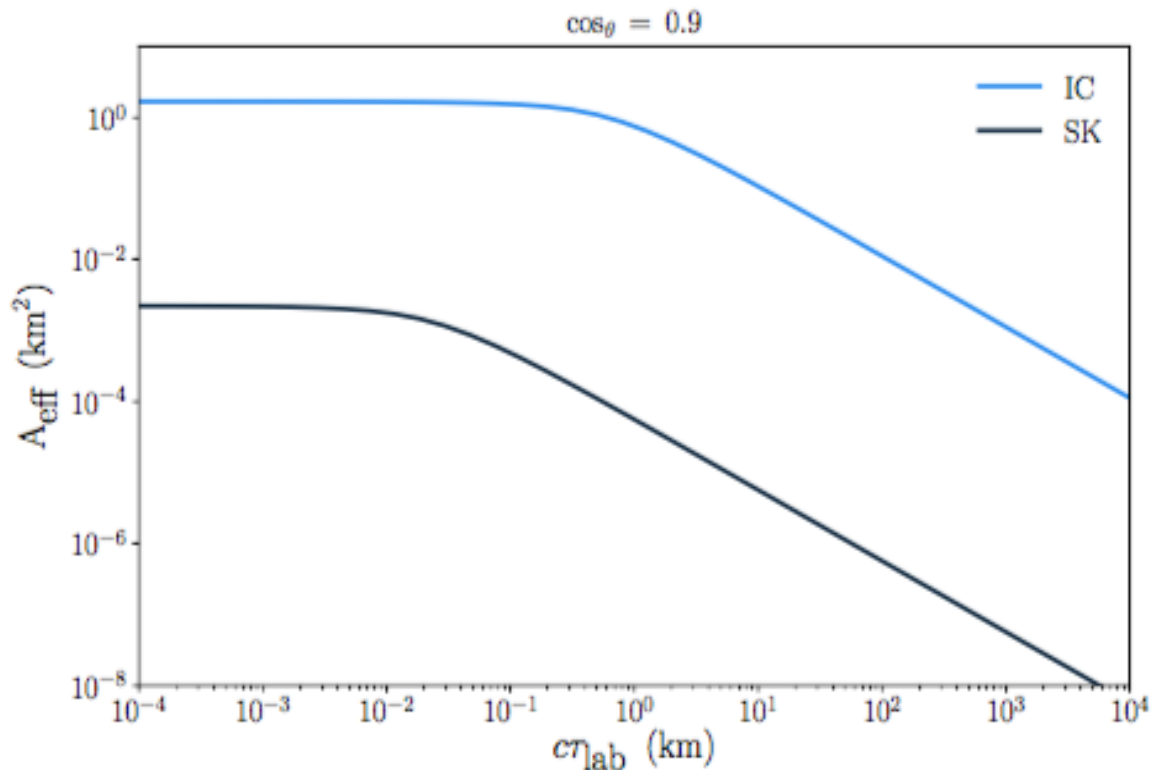
$$\frac{dN_i^\alpha}{dE d\cos\theta} \propto \text{Br}(\alpha\text{-like}) \underbrace{\epsilon^\alpha}_{\text{Detector efficiencies (estimated)}} \underbrace{\Delta T A_{\text{decay}}^{\text{eff}}(E, \cos\theta)}_{\text{Geometric acceptance}} \frac{d\Phi_A}{dE d\cos\theta}$$

Detector efficiencies
(estimated)

Geometric acceptance

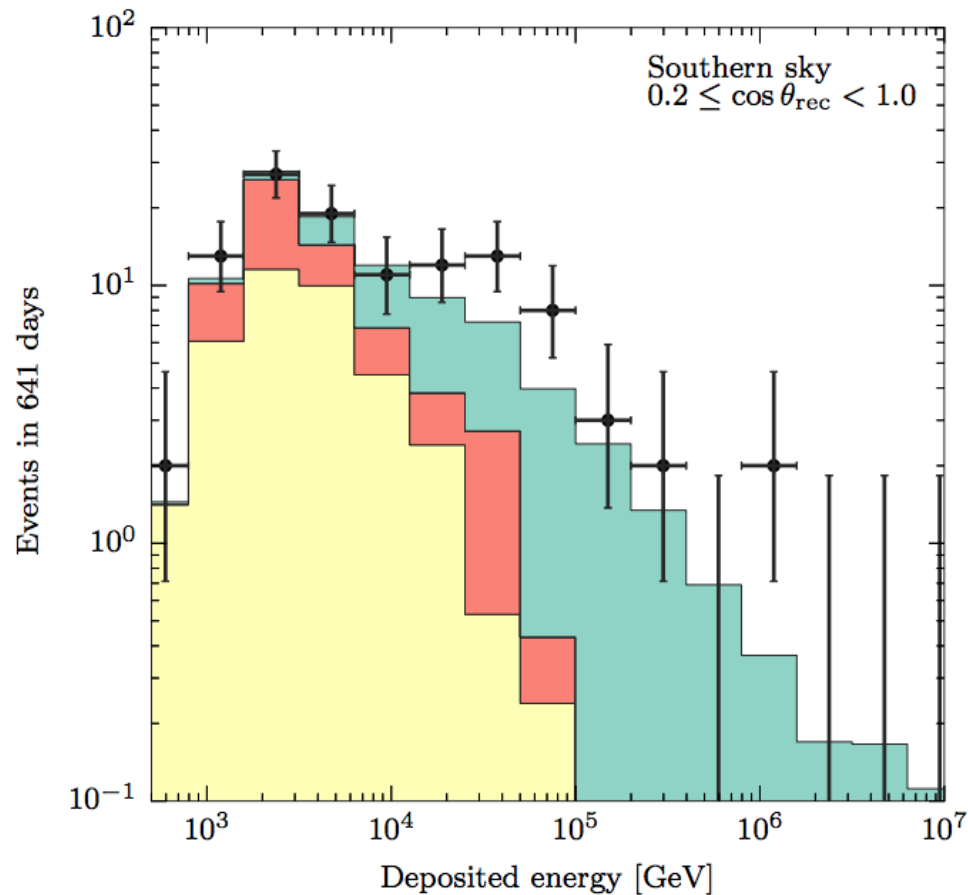


Pilar Coloma - IFIC



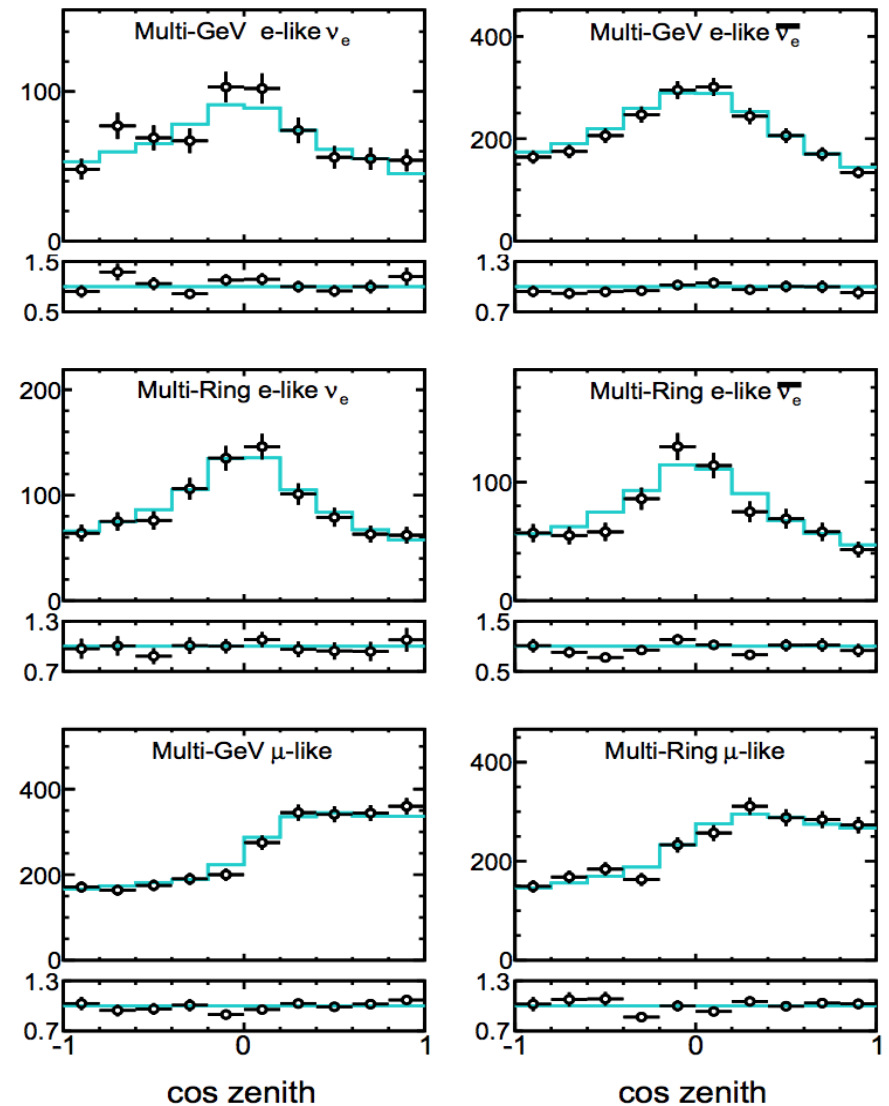
Backgrounds and data

MESE sample at Icecube (641 days)
Icecube collaboration, 1410.1749

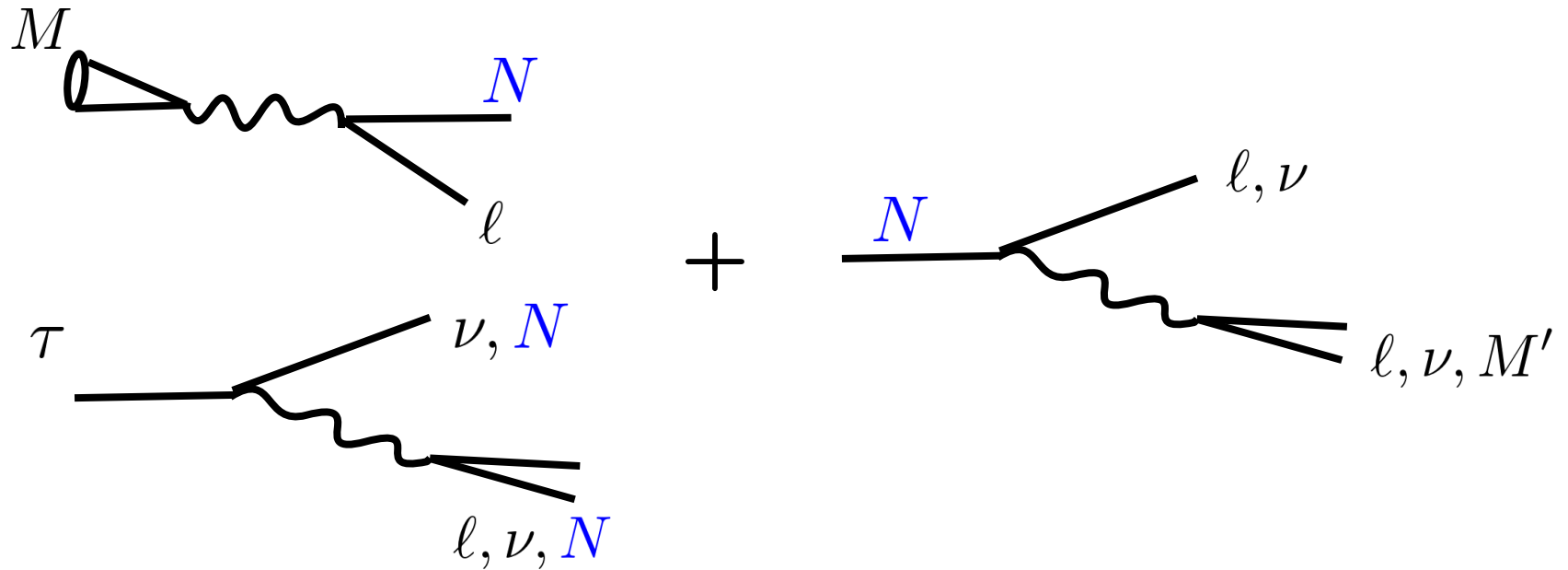


Pilar Coloma - IFIC

Data from SK I-IV (5,326 days)
Super-Kamiokande coll. 1710.09126

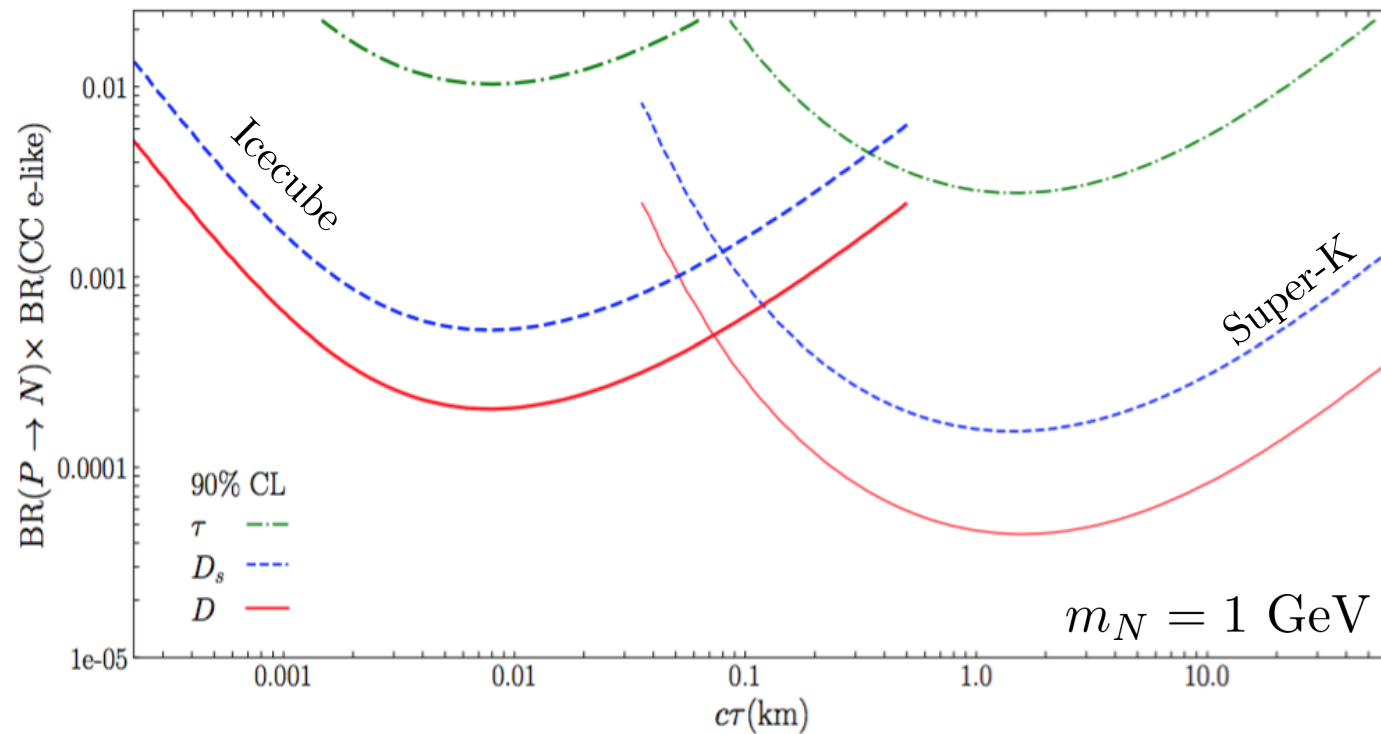
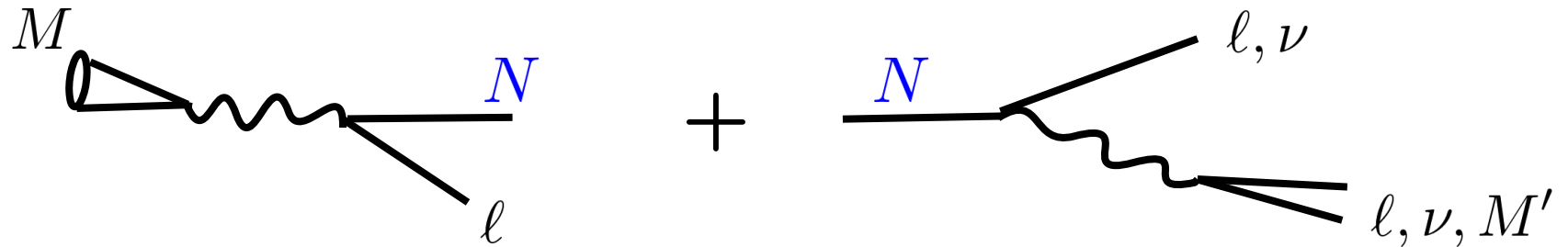


HNL production and decay



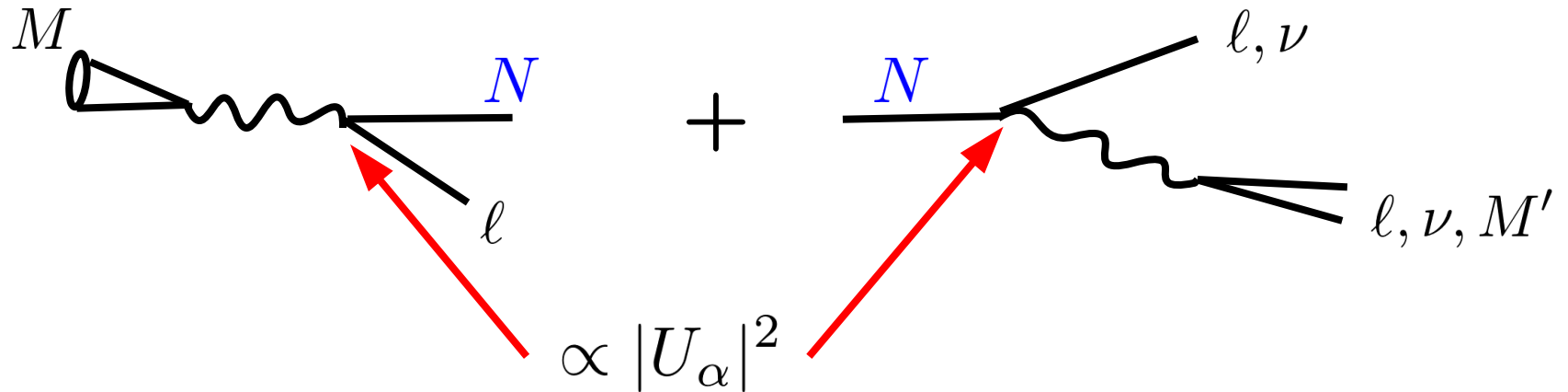
→ Maximum sensitivity if decay length in the lab frame is $O(10)$ km

HNL **above** the kaon mass



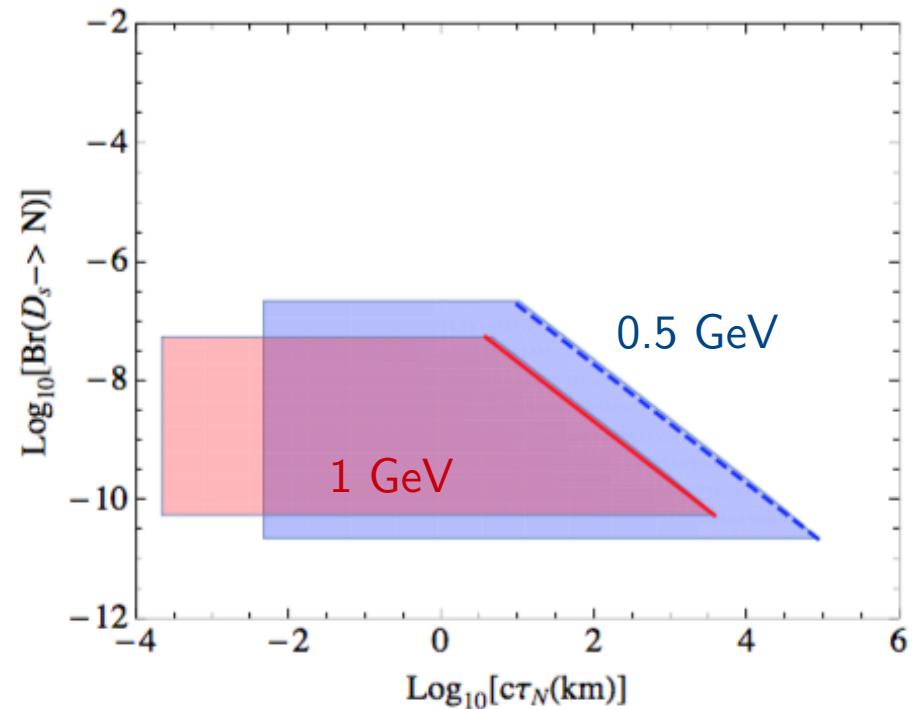
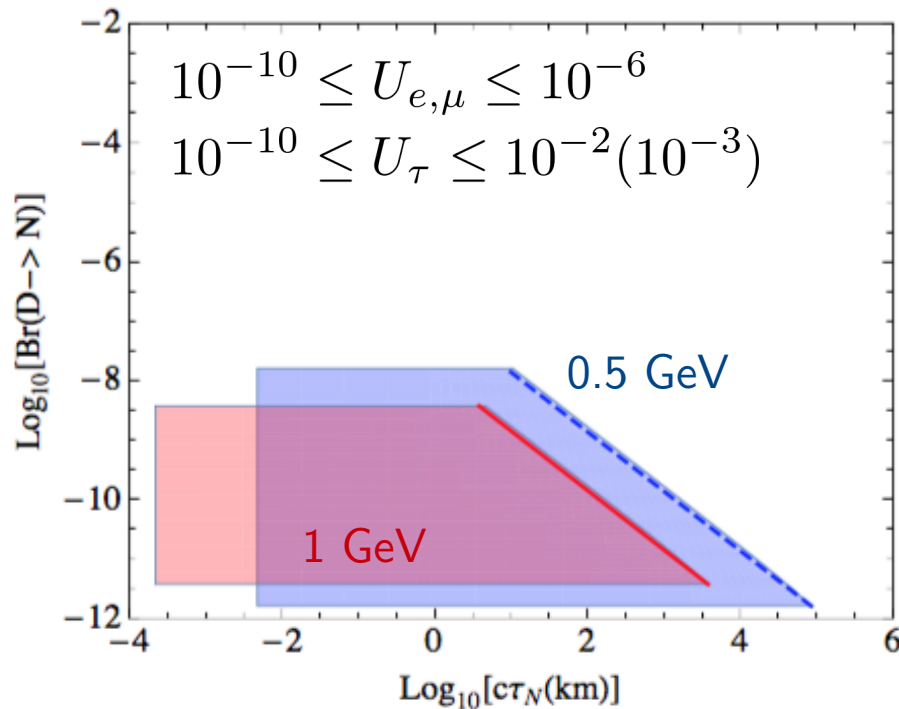
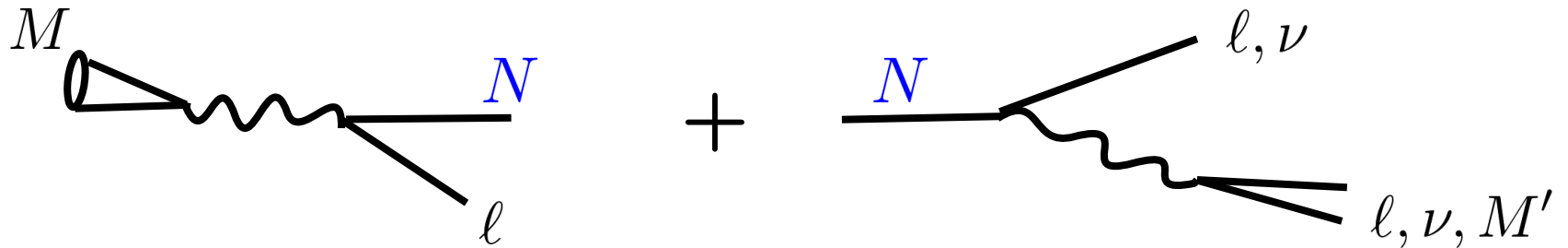
Argüelles, Coloma, Hernandez, Muñoz, 1910.12839

HNL: minimal scenario



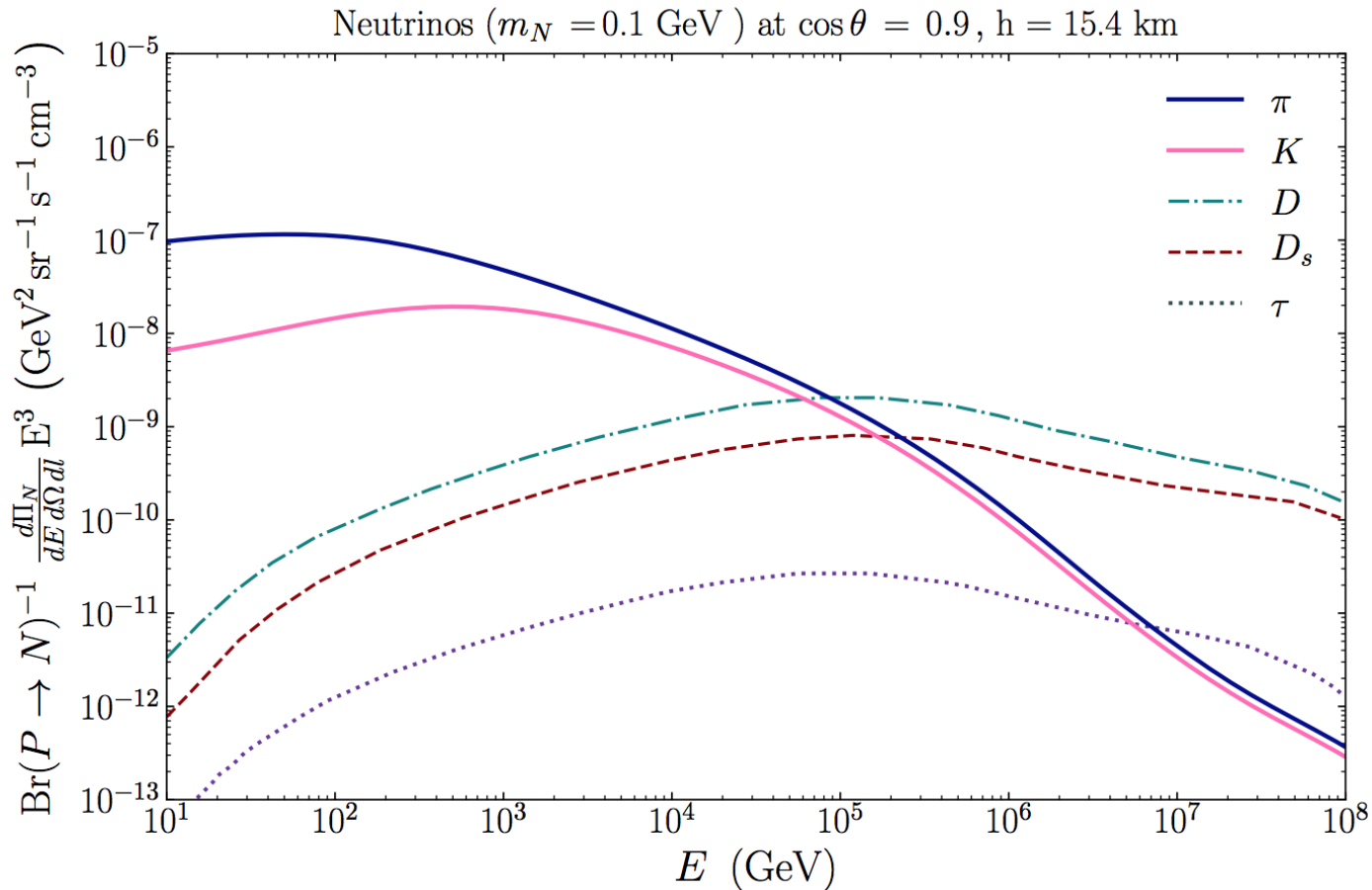
$$\mathcal{L}_N = \mathcal{L}_{SM} + \sum_j i\bar{N}_j \gamma^\mu \partial_\mu N_j - \left(Y_{\alpha j} \bar{L}_\alpha \tilde{\Phi} N_j + \frac{m_{N_j}}{2} \bar{N}_j N_j^c \right)$$

HNL: minimal scenario



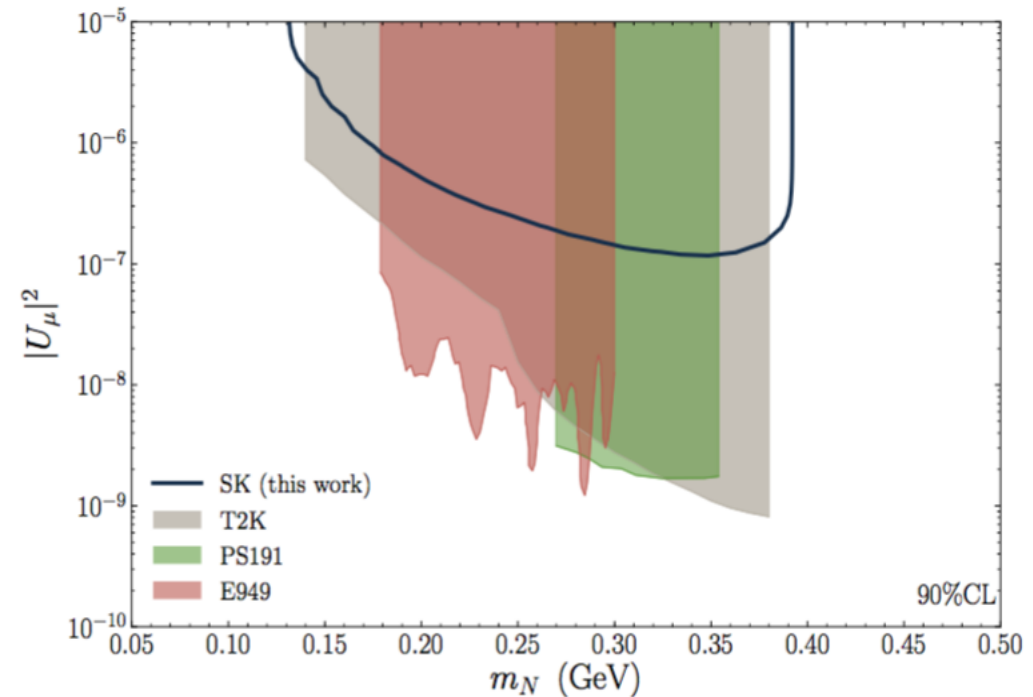
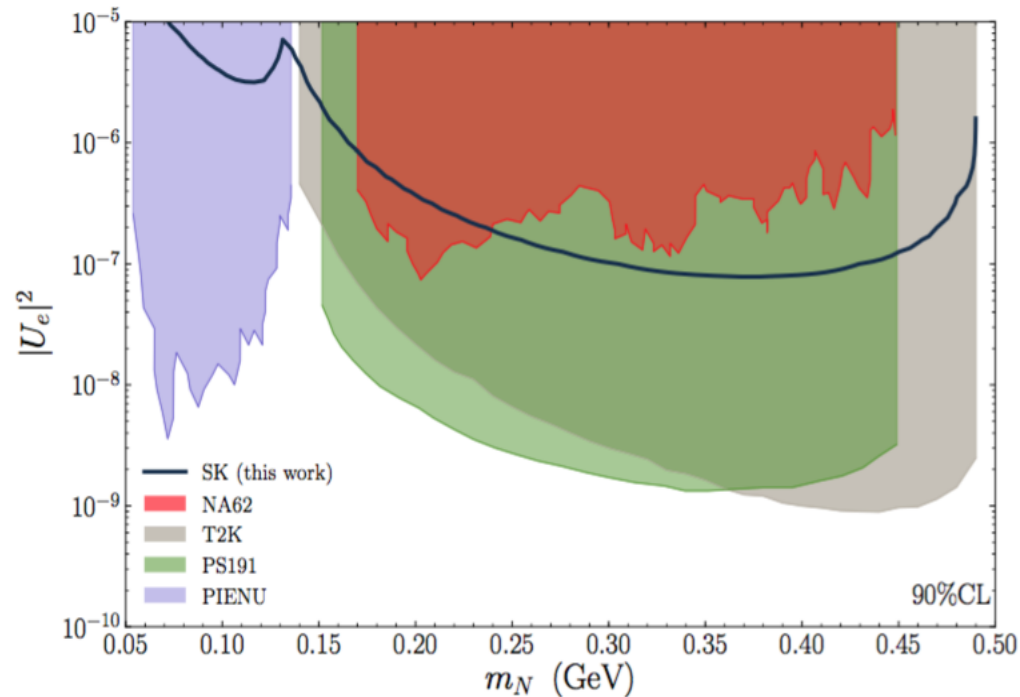
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HNL production per parent



Argüelles, Coloma, Hernandez, Muñoz, 1910.12839
Coloma, Hernandez, Muñoz, Shoemaker, 1911.09129

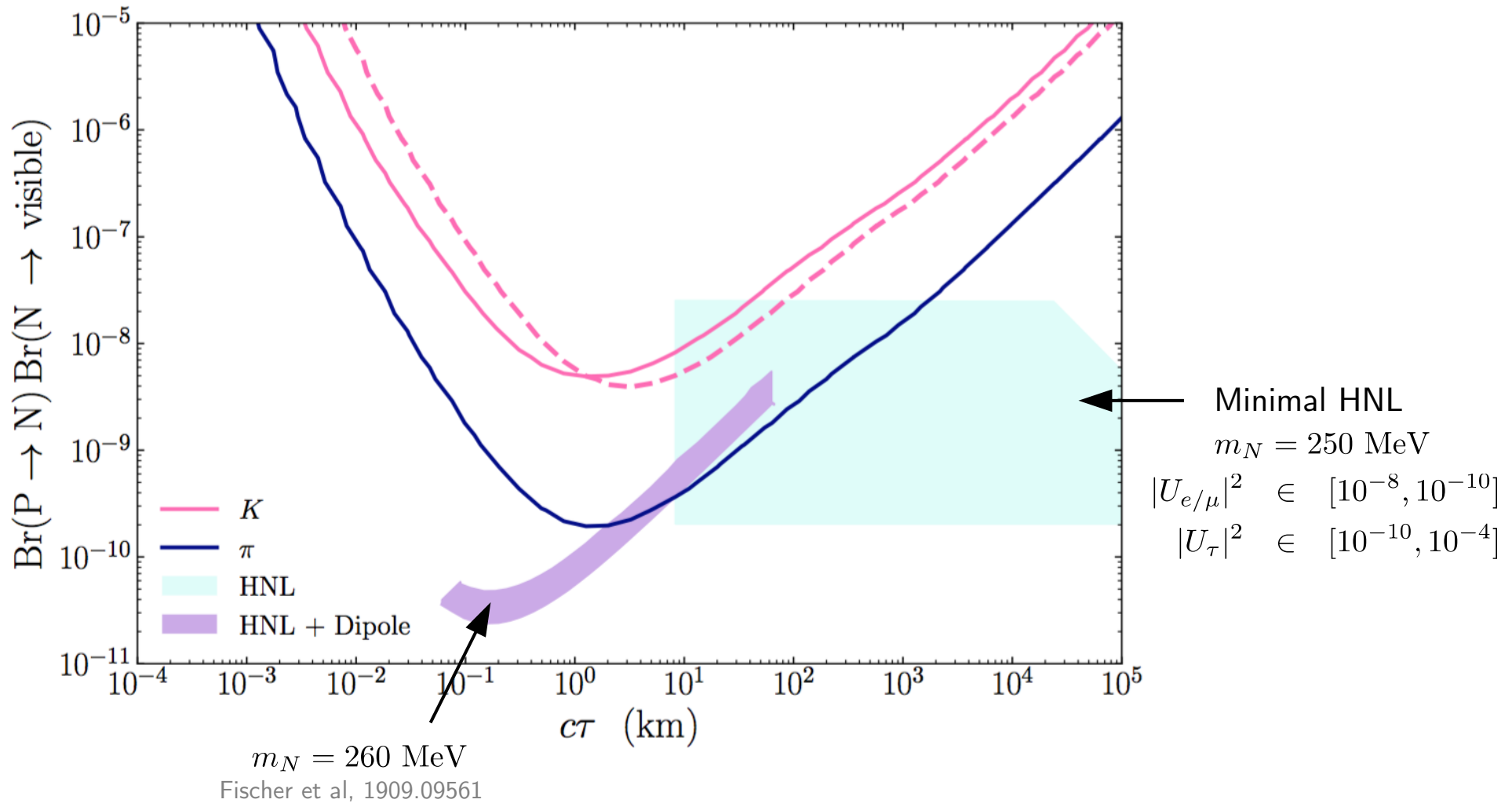
HNL below the kaon mass



Coloma, Hernandez, Muñoz, Shoemaker, 1911.09129
(see also Asaka and Watanabe, 1202.0725, Kusenko, Pascoli and Semikoz, hep-ph/0405198)

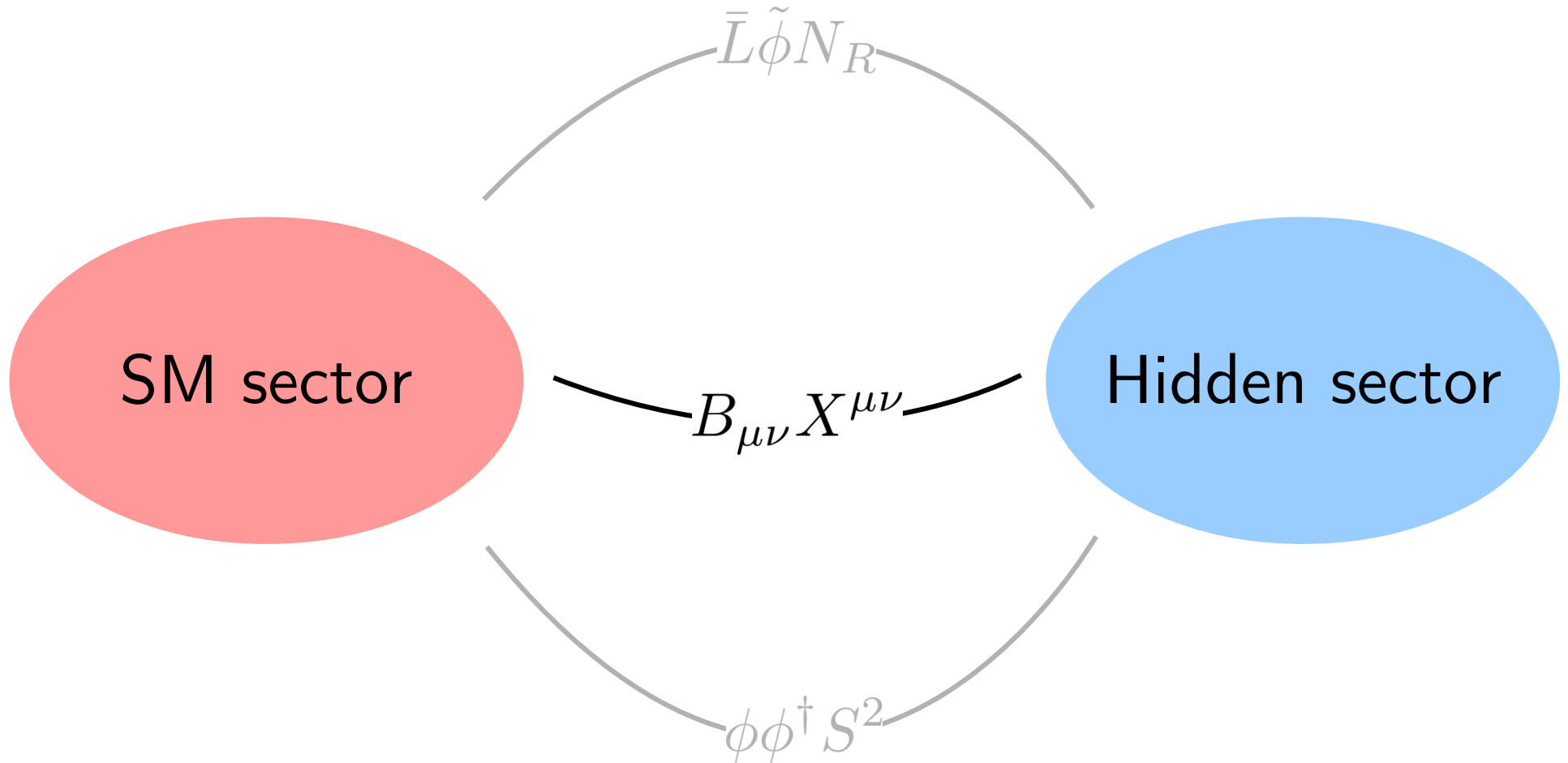
HNL below the kaon mass

Coloma, Hernandez, Muñoz, Shoemaker, 1911.09129



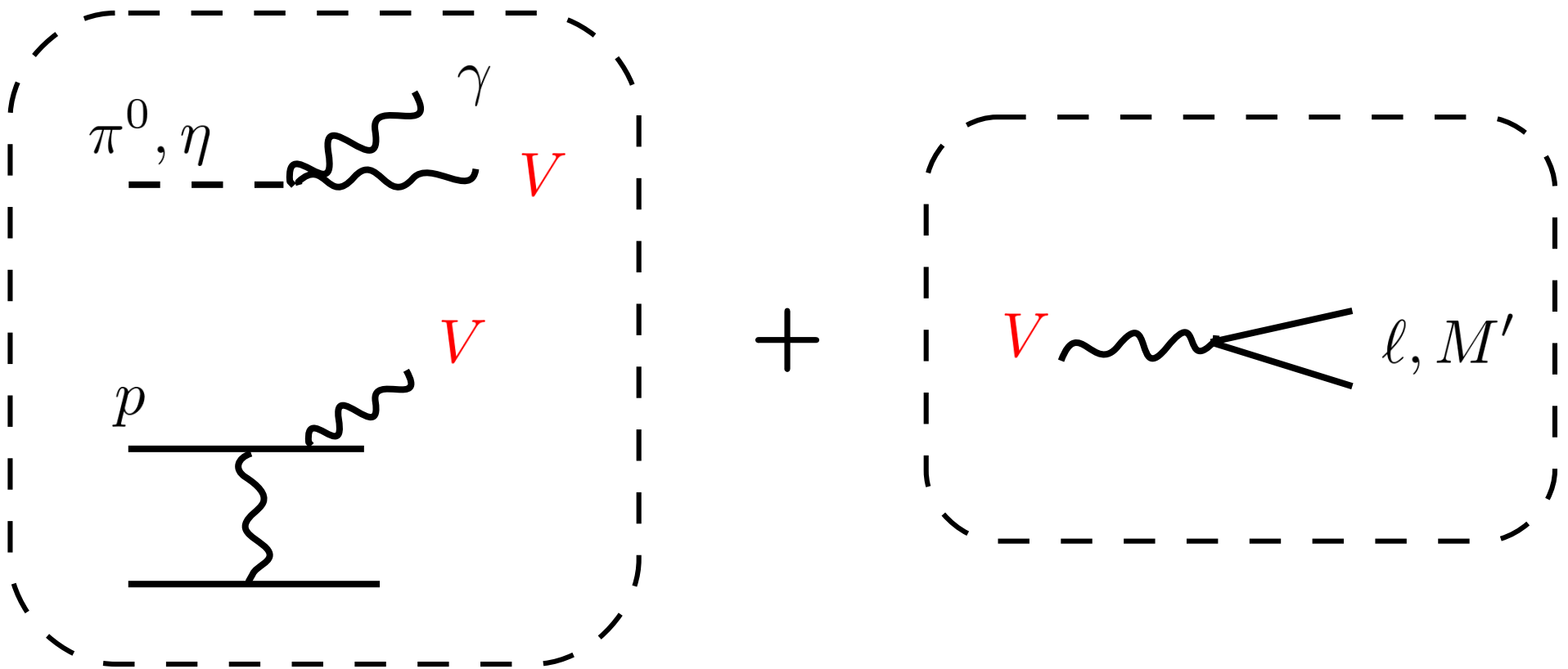
Dark photons

Where is the new physics?

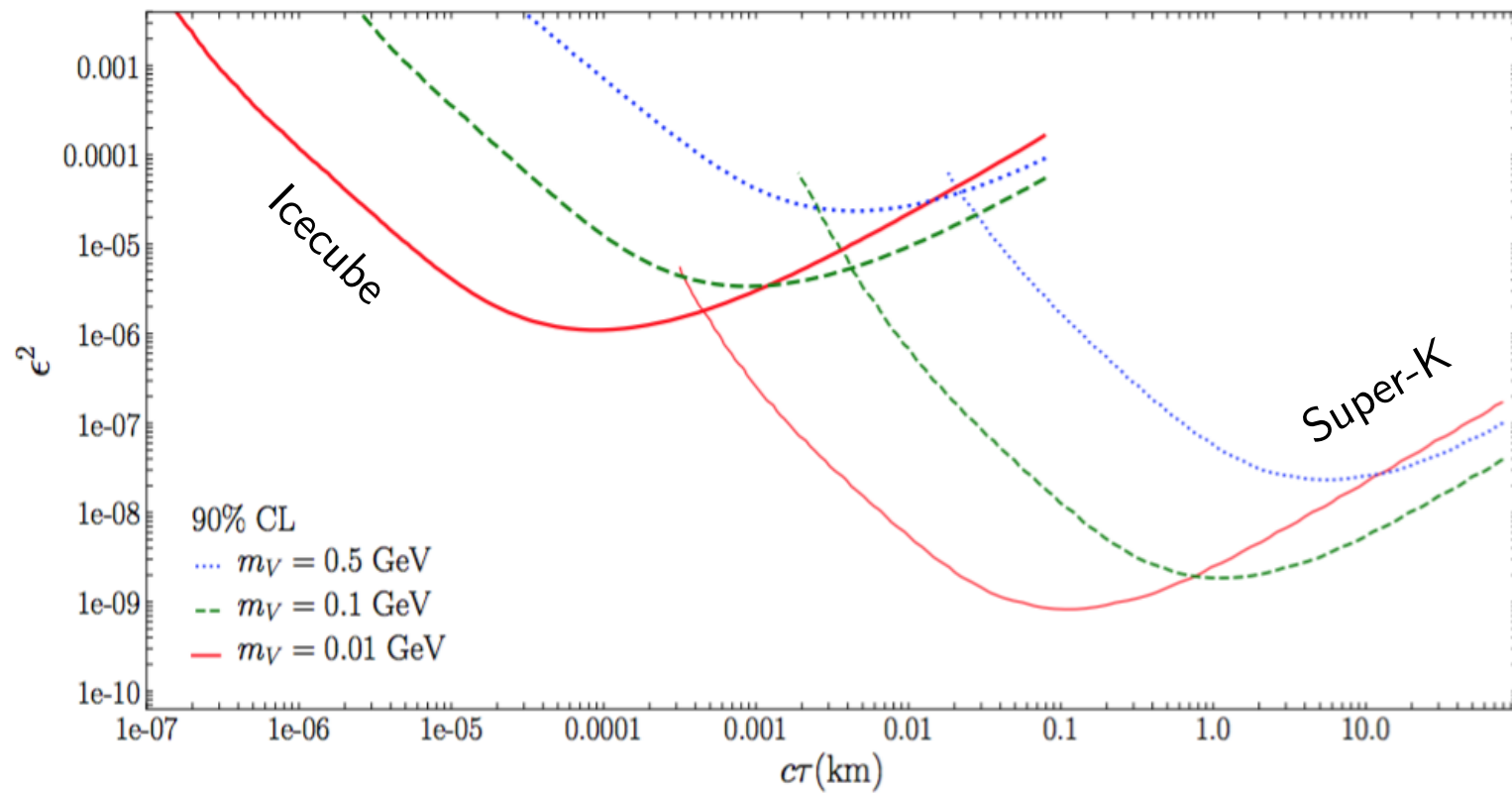


Dark photons: kinetic mixing

$$\mathcal{L}_V = \mathcal{L}_{SM} - \frac{1}{4}V_{\mu\nu}V^{\mu\nu} + \frac{\epsilon}{2}B_{\mu\nu}V^{\mu\nu} - \frac{1}{2}m_V^2 V_\mu V^\mu$$

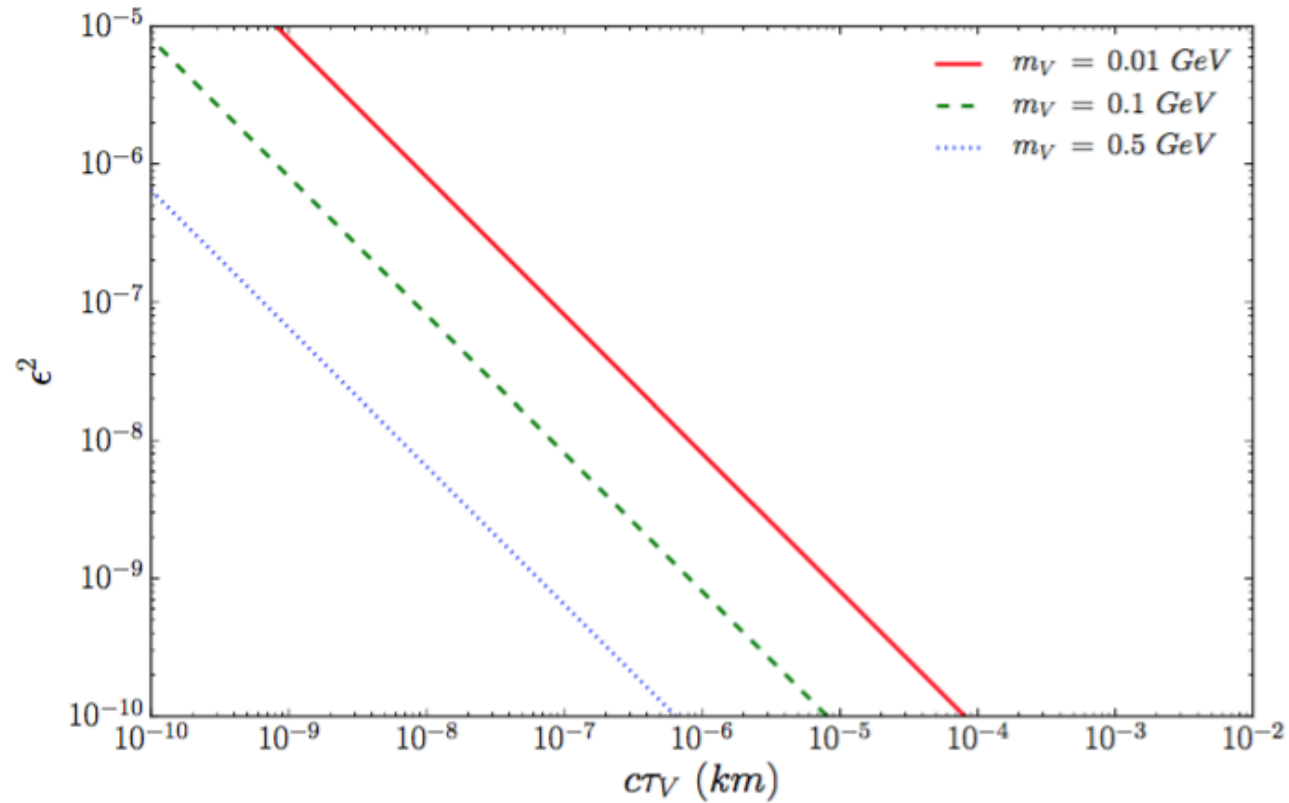


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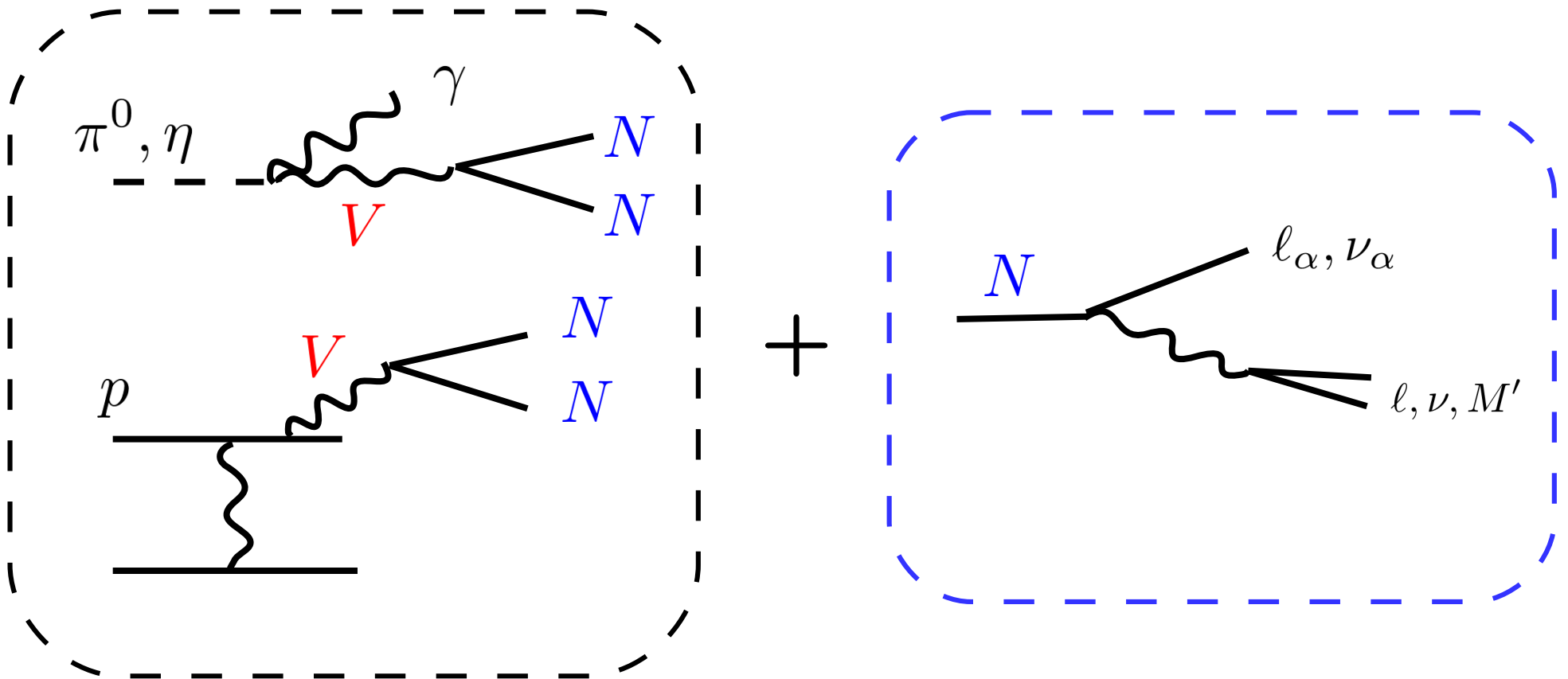
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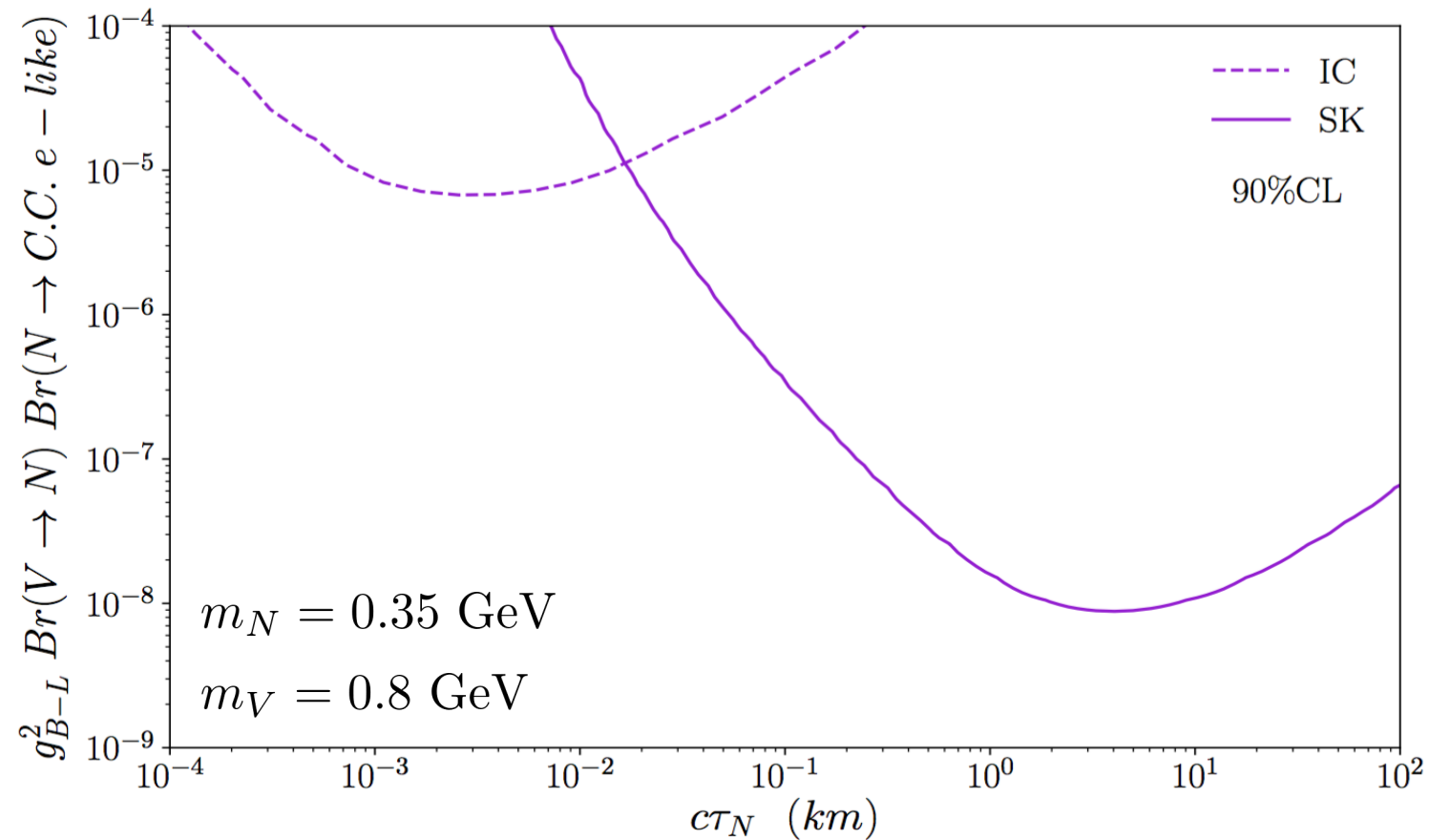
Argüelles, Coloma, Hernandez, Muñoz, 1910.12839

Dark photons: B-L model

$$\mathcal{L} = \mathcal{L}_N - \frac{1}{4}V_{\mu\nu}V^{\mu\nu} - \frac{1}{2}m_V^2 V_\mu V^\mu + g_{B-L} \left(\sum_f Q_{B-L}^f \bar{f} \gamma^\mu f V_\mu - \bar{N} \gamma^\mu N V_\mu \right)$$



Dark photons: B-L model



Argüelles, Coloma, Hernandez, Muñoz, 1910.12839

Summary

- Atmospheric neutrino detectors are well-suited to search for the decay products of long-lived particles produced **in the atmosphere**
 - Optimal sensitivity for decay lengths in the range of tens of kilometers
- We have derived exclusion limits using Icecube and SK data
 - We have studied HNL, dark photons and a B-L model
 - Our limits are shown in the $BR - c\tau$ plane, allowing to be easily interpreted for non-minimal scenarios
 - SK able to constrain the minimal HNL model as well, for m_N below the kaon mass

Thank you!!