

## Abstract

Terrestrial and solar neutrino experiments have a variety of anomalous data that has resisted clarification. Recently, it has appeared that measurements of neutrinos from intense sources on gallium have passed 5 and other hints from MicroBooNE and elsewhere remain interesting. I will present the latest update of these anomalies. I will then explain the primary reasons why these cannot be simply interpreted as a 1 eV sterile neutrino due to constraints from other experimental probes, notably solar neutrinos and cosmological data sets. I will present a novel, simple model that evades many of these constraints by adding in one new particle, which is the dark matter, beyond a sterile neutrino leading to shape-shifting sterile neutrinos.

# Light Sterile Neutrinos: A Modern Picture and a Model to Evade Cosmology

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WIN

July 7, 2023



**Brookhaven**<sup>™</sup>  
National Laboratory



# Overview

1. Sterile neutrino picture through 2020
  - ▶ Cosmology!
2. MicroBooNE
3. Evading cosmology

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  - ▶ Cosmology!
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Data is confusing  
Up to you to decide

# Neutrinos have mass

- ▶ Can get usual Dirac mass term via Higgs
  - ▶  $\Rightarrow$  three new right-handed neutrinos
- ▶ Steriles can have additional mass terms
  - ▶ Seesaw?

H. Fritzsch, M. Gell-Mann, P. Minkowski [PLB 1975](#)  
P. Minkowski [PLB 1977](#)  
W. Konetschny, W. Kummer [PLB 1977](#)  
D. Wyler, L. Wolfenstein [NPB 1983](#)  
R. Foot, H. Lew, X. He, G. Joshi [ZPC 1989](#)

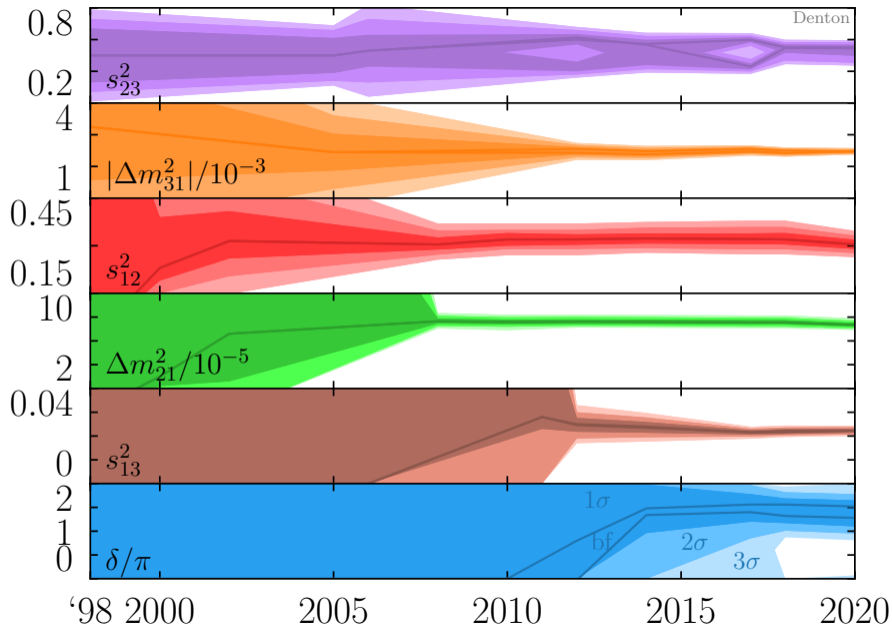
- ▶ Pseudo-Dirac?

L. Wolfenstein [NPB 1981](#)  
S. Bilenky, S. Petcov [RMP 1987](#)

- ▶ Some options have no sterile neutrinos, but other new particles
  - ▶ E.g. type-II seesaw

Interesting mass ranges are often  $10^{13}$  GeV,  $10^3$  GeV, or  $10^{-26}$  GeV, not  $10^{-9}$  GeV

# Three flavor oscillation picture



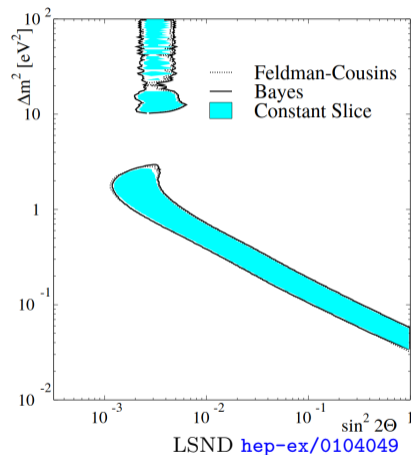
Three flavor oscillation picture: looks good

Let's check many  $\Delta m^2$ 's!



# Accelerator: LSND

- ▶ LSND ran from 1993-1998
  - ▶  $E_{\bar{\nu}_\mu} \in [20, 53]$  MeV
  - ▶  $L = 30$  m
  - ▶ Looked for  $\bar{\nu}_\mu \rightarrow \bar{\nu}_e$  appearance
  - ▶ Excess of:  $87.9 \pm 22.4 \pm 6.0 \Rightarrow 3.8\sigma$  (1 dof)
  - ▶ Interesting region:
    - ▶  $\Delta m_{41}^2 \sim 1$  eV<sup>2</sup>
    - ▶  $\sin^2 2\theta_{\mu e} = 4|U_{e4}|^2|U_{\mu4}|^2 \sim 0.002$
- OPERA, ICARUS disfavor  $\sin^2 2\theta_{\mu e} \gtrsim 0.02$



# Accelerator: MiniBooNE

- ▶ MiniBooNE ran from 2002 to 2019
- ▶ Built to test LSND, higher energy, longer baseline, similar  $L/E$ , both  $\nu, \bar{\nu}$
- ▶  $E_{\nu_\mu} \sim 500$  MeV
- ▶  $L = 541$  m
- ▶ Excesses:
  - ▶  $\nu_e$ :  $381.2 \pm 85.2 \Rightarrow 4.5\sigma$  (1 dof)
  - ▶  $\bar{\nu}_e$ :  $79.3 \pm 28.6 \Rightarrow 2.8\sigma$  (1 dof)
  - ▶ Combined:  $4.7\sigma$  (1 dof)
  - ▶ Excesses consistent with LSND under sterile hypothesis
  - ▶ Combined with LSND:  $\Rightarrow 6.0\sigma$  (1 dof)

MiniBooNE [1805.12028](#)

# Accelerator experiment caveats

- ▶ Neither LSND nor MiniBooNE is particularly well fit by a sterile
  - ▶ The excess grows at lower energies faster than it should
  - ▶ Not necessarily a huge problem
- ▶ LSND result may not be robust under cut assumptions

J. Hill [hep-ex/9504009](#)

- ▶ Not a problem for MiniBooNE

MiniBooNE [2006.16883](#)

- ▶  $\nu_e$  appearance requires both  $\nu_\mu$  disappearance and  $\nu_e$  disappearance
  - ▶ Since  $|U_{\mu 4}|^2 |U_{e 4}|^2 > 0$  and  $|U_{\alpha i}| \in [0, 1]$ ,  $\exists$  lower limits on both  $|U_{\mu 4}|$  and  $|U_{e 4}|$

# The gallium experiments

- ▶ Low energy solar neutrino experiments measure the  $pp$  flux

GALLEX: 1991-1997, GNO: 1998-2003 [1001.2731](#)  
SAGE: 1989-2007 [0901.2200](#)

- ▶ Consistent with KamLAND

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- ▶ Calibrate detectors with intense radioactive sources
- ▶ See fewer neutrinos than expected:

$3.0\sigma$ : C. Giunti, M. Laveder [1006.3244](#)  
 $2.3\sigma$ : J. Kostensalo, et al. [1906.10980](#)  
 $> 4\sigma$ : BEST [2109.11482](#)  
 $\rightarrow > 5\sigma$ : C. Giunti, et al. [2212.09722](#)

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- ▶ Prefers:
  - ▶  $\Delta m_{41}^2 \gtrsim 0.5 \text{ eV}^2$
  - ▶  $\sin^2 2\theta_{ee} = 4|U_{e4}|^2(1 - |U_{e4}|^2) \sim 0.4$

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- ▶ Attempts to explain with standard physics: unsuccessful

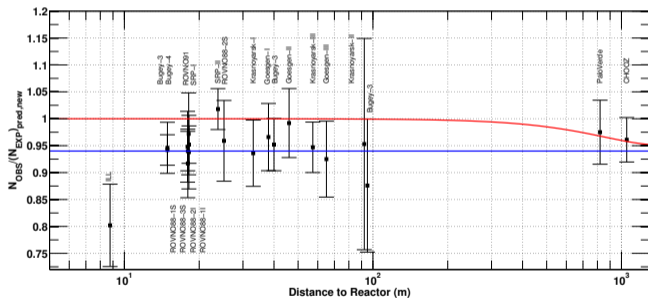
C. Giunti, et al. [2212.09722](#)  
V. Brdar, J. Gehrlein, J. Kopp [2303.05528](#)  
S. Elliott, V. Gavrin, W. Haxton [2303.13623](#)

# Reactor rates

Deficit relative to prediction

P. Huber [1106.0687](#)

T. Mueller, et al. [1101.2663](#)

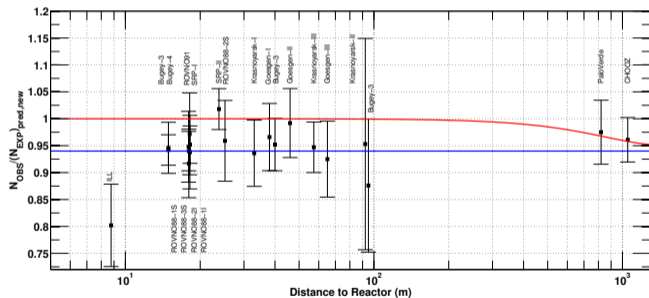


G. Mention, et al. [1101.2755](#)



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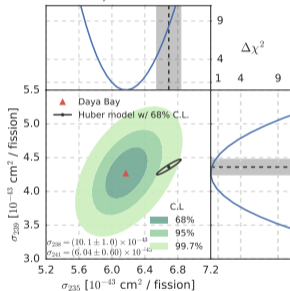
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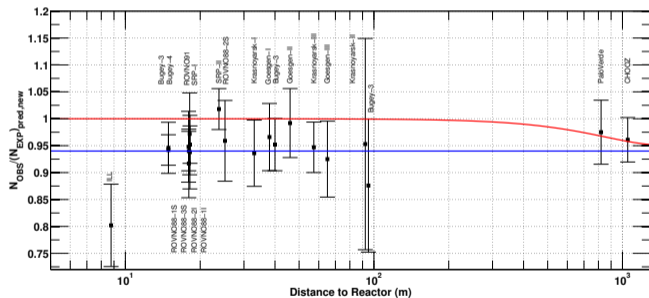
Daya Bay [1704.01082](#)

RENO [1806.00574](#)

Daya Bay, PROSPECT [2106.12251](#)

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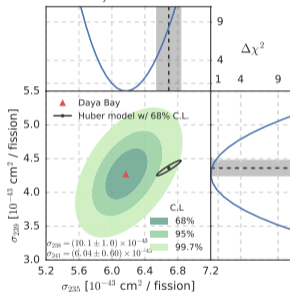
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Probably nuclear physics

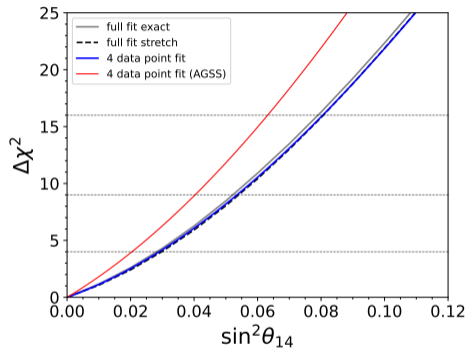
## Short baseline spectral

- ▶ NEOS, DANSS see some spectral anomalies
  - ▶  $\Delta m_{41}^2 = 1.26 \text{ eV}^2$  and  $\sin^2 2\theta_{14} = 0.044$  at  $3.3\sigma$
- ▶ Mixings larger than  $\sin^2 2\theta_{14} \sim 0.01$  disfavored by spectral data
- ▶ Neutrino-4 also sees spectral anomalies
  - ▶  $\Delta m_{41}^2 = 7.32 \text{ eV}^2$  and  $\sin^2 2\theta_{14} = 0.31$
  - ▶ In tension with other reactor data
  - ▶ Analysis issues

J. Berryman, P. Huber [2005.01756](#)

# Solar

1. Use gallium and Borexino for  $pp$  data
2. Use SNO and SK for  $^8\text{B}$  data  
No Borexino data?
3. Use KamLAND data to set  $\Delta m_{21}^2$
4. Fix  $\theta_{13}$  to best fit
5. Vary  $\theta_{12}$  and  $\theta_{14}$
6. Consider impact on  $U_{e4}$  ( $\theta_{14}$ ) only
7. Applies for  $\Delta m_{41}^2 \gtrsim 10^{-3} \text{ eV}^2$
8. Is effectively a unitary violation analysis
9. Checked Wilks' theorem with MC



K. Goldhagen, et al. [2109.14898](#)

So far:

Have anomalous  $\nu_\mu \rightarrow \nu_e$

LSND, MiniBooNE

Might have anomalous  $\nu_e \rightarrow \nu_e$

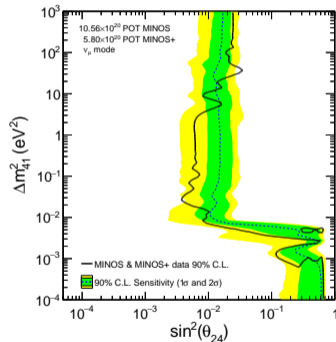
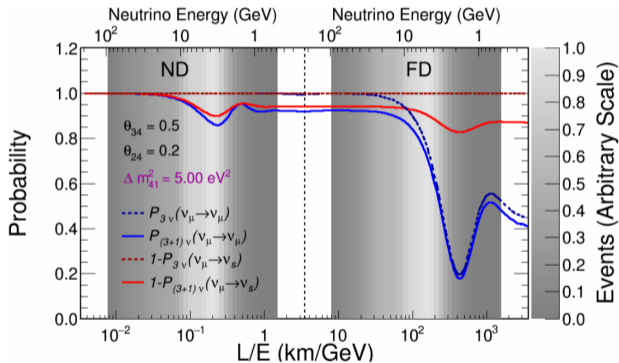
Yes: Gallium, ~~Reactor rate~~

No: Reactor spectral, solar

Do we have anomalous  $\nu_\mu \rightarrow \nu_\mu$ ?

# MINOS/MINOS+

- ▶ MINOS ran from 2005-2012, MINOS+ (higher energy) ran from 2013-2016
- ▶ Leverage near- and far-detectors simultaneously



MINOS [1710.06488](#)

Some concerns, e.g. W. Louis [1803.11488](#)

# IceCube

At  $E \sim 1$  TeV and  $\Delta m_{41}^2 \sim 1$  eV<sup>2</sup>,

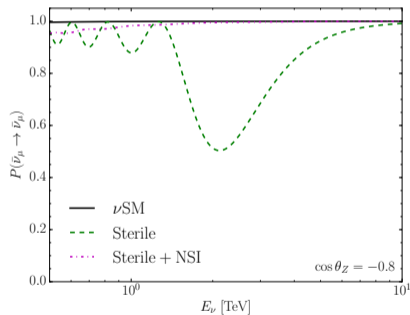
$\bar{\nu}_\mu \rightarrow \bar{\nu}_\mu$  experiences large disappearance through the Earth's core

H. Nunokawa, O. Peres, R. Funchal [hep-ph/0302039](#)

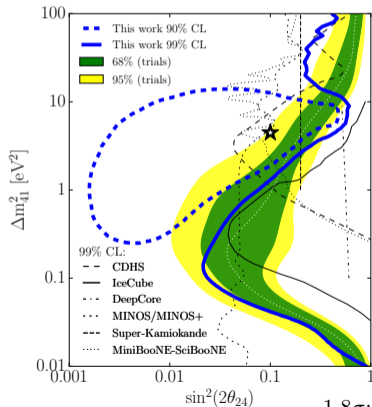
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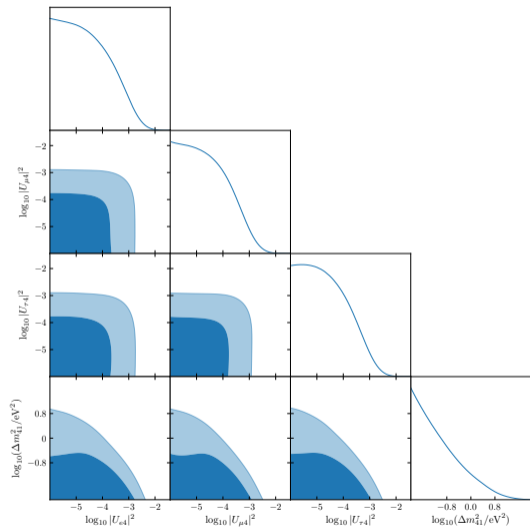
PBD, Y. Farzan, I. Shoemaker [1811.01310](https://arxiv.org/abs/1811.01310)



1.8 $\sigma$ : IC [2005.12942](https://arxiv.org/abs/2005.12942)



# Cosmological bounds

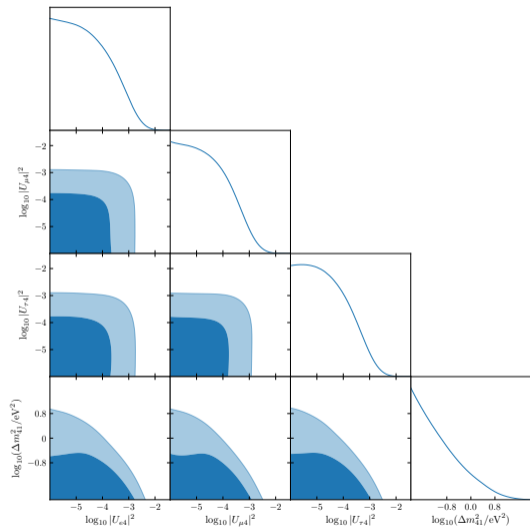


1 $\sigma$ , 2 $\sigma$

S. Hagstotz, et al. [2003.02289](#)

- ▶ Includes CMB temperature, polarization, and lensing, and BAO
- ▶ No local  $H_0$  constraint
- ▶ Bounds independent of flavor
- ▶ To be consistent with data must have small mixing **and** small mass

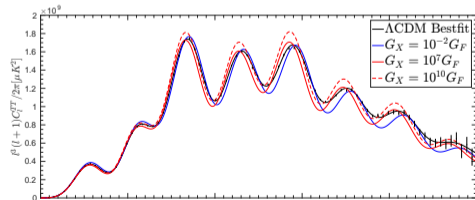
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- ▶ No local  $H_0$  constraint
- ▶ Bounds independent of flavor
- ▶ To be consistent with data must have small mixing **and** small mass
- ▶ Much more than just  $N_{\text{eff}}$  and  $\sum m_\nu$
- ▶ Just adding a new interaction is not straightforward



N. Song, M. Gonzalez-Garcia, J. Salvado [1805.08218](#)

# Cosmological bounds with an interaction

- ▶ Include  $H_0$  and  $\sigma_8$  tensions
- ▶ Data prefers:  $N_{\text{eff}} = 4.02 \pm 0.29$  and  $G_X \sim 10^8 G_F$ 
  - C. Kreisch, F. Cyr-Racine, O. Doré [1902.00534](#)
  - G. Barenboim, [PBD](#), I. Oldengott [1903.02036](#)
- ▶ Large self-interaction is constrained by:
  - ▶  $Z \rightarrow$ invisible for large couplings
  - ▶ BBN+CMB for light masses
  - ▶ Kaon decays for all remaining parameter space for  $\nu_e, \nu_\mu$
- ▶ Viable space persists  $m_X \sim 10$  MeV if the self interaction is in the  $\nu_\tau$  sector
  - N. Blinov, et al. [1905.02727](#)
- ▶ Testable by IceCube looking for dips due to  $C\nu B$ 
  - G. Barenboim, [PBD](#), I. Oldengott [1903.02036](#)
  - C. Creque-Sarbinowski, J. Hyde, M. Kamionkowski [2005.05332](#)
  - I. Esteban, et al. [2107.13568](#)

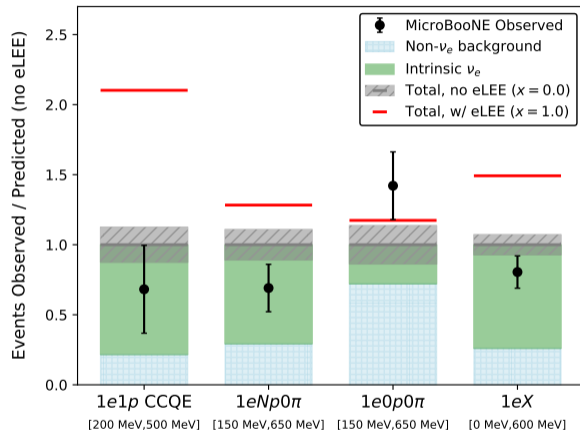
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Not a great fit to the cosmological data  
Other new physics (cosmo) scenarios fit the data better

Let's resolve this terrestrially

# MicroBooNE results



- ▶ Three analysis teams:
  1. Wire-Cell
  2. Deep Learning
  3. Pandora
    - ▶ With 0 protons
    - ▶ With 1+ protons
- ▶ Underfluctuation compared to no-oscillations
- ▶ Disfavors MiniBooNE's best fit LEE hypothesis at  $3.75\sigma$

MicroBooNE [2110.14054](#)

## MicroBooNE disappearance

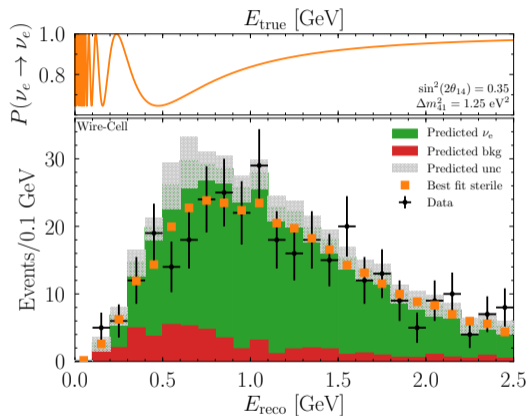
MicroBooNE is focused on  $\nu_e$  appearance  
Can do  $\nu_\mu$  and  $\nu_e$  disappearance too!

See also D. Cianci, et al. [1702.01758](#)

MiniBooNE backgrounds too big, plus anomaly

# Dip hunting

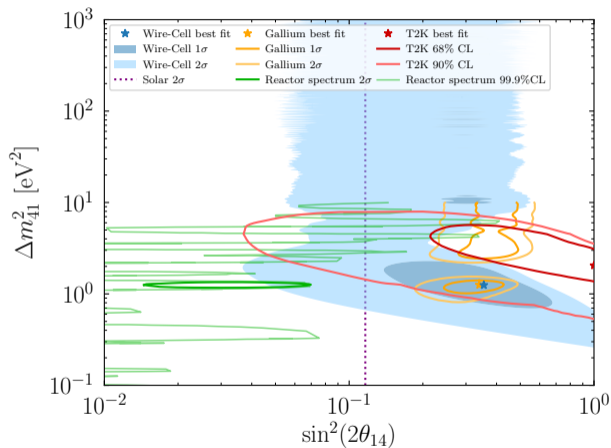
- ▶ 4 analysis channels
  - ▶ Wire cell has most statistics
  - ▶ Analyses not fully independent
- ▶ Dip appears in multiple analyses





# Global $\nu_e$ disappearance picture

MicroBooNE prefers  $\nu_e$  disappearance at  $\sim 2.4\sigma$



Cosmology disfavors entire plane!

# What does it take to evade cosmology?

[2301.09651](#) with Hooman Davoudiasl

See also:

Y. Farzan [1907.04271](#)

V. Brdar, J. Gehrlein, J. Kopp [2303.05528](#)

## Shape-shifting sterile neutrinos

- ▶ Sterile neutrinos seem to act differently in different places:
  - ▶ Earth's surface
  - ▶ Sun
  - ▶ Early universe
- ▶ Suppose sterile neutrino talk to nucleons via long-range scalar  $\phi$
- ▶ Suppose  $m_\phi \sim 5 \times 10^{-15}$  eV  $\Rightarrow 1/m_\phi \sim 40,000$  km  $\sim 6R_\oplus$

A broad range of values works

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- ▶  $\phi$  - nucleon coupling below fifth force limits:  $g_n \sim 5 \times 10^{-25}$

MICROSCOPE [2209.15487](#)

- ▶ At Earth's surface, field has non-zero value:

$$\phi^\oplus \approx -\frac{g_n N_n^\oplus}{4\pi R_\oplus} e^{-m_\phi R_\oplus} = -4 \times 10^{12} \text{ eV}$$

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- ▶ Take  $m_0 = 1$  eV and  $g_s \sim 5 \times 10^{-14} \Rightarrow |g_s \phi^\oplus| = 0.2$  eV

- ▶ Also need a bare mass term for the new mass state:  $m_s(\vec{x}) = m_0 + g_s \phi(\vec{x})$

$m_0 \neq 0$  needed for cosmological  $\sum m_\nu$

# Shape-shifting sterile neutrinos

Dirac mass matrix:

$$M_\nu = \begin{pmatrix} m_\nu & m_D \\ 0 & m_s(\vec{x}) \end{pmatrix} \begin{matrix} \nu_{e,\mu,\tau} \\ \nu_s \end{matrix}$$

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$$\tan 2\theta_{14} \simeq \frac{2m_D m_s(\vec{x})}{m_s^2(\vec{x}) - m_D^2 - m_\nu^2}$$

$$m_1 \simeq m_\nu \frac{m_s(\vec{x})}{\sqrt{m_s^2(\vec{x}) + m_D^2}}$$

$$m_{2,3} \simeq m_\nu$$

$$m_4 \simeq \sqrt{m_s^2(\vec{x}) + m_D^2}$$

Set  $m_\nu = 0.03$  eV and  $m_D = 0.3$  eV

## Self interacting scalar

Start with a potential:

$$V(\phi) = \frac{1}{2}m_\phi^2\phi^2$$

But loops will add in a term:

$$V(\phi) = \frac{1}{2}m_\phi^2\phi^2 + \frac{\lambda}{4!}\phi^4$$

$$\delta\lambda \sim \frac{g_s^4}{16\pi^2} \sim 4 \times 10^{-56}$$



# Shape-shifting sterile neutrinos in the early universe

To avoid cosmological constraints, need:

1.  $m_s \gtrsim \text{keV}$  ( $\theta_{14} \lesssim 10^{-3}$ )
2. Thus  $\phi_{\text{BBN}} \gtrsim 10^{16}$  eV
3. At minimal reheating temp  $T_{\text{rh}} \sim 10$  MeV, need  $\phi_i \gtrsim \text{few} \times 10^{16}$  eV
4. Various ways to do this, e.g. thermal misalignment

D. Marsh [1510.07633](#)

5. At  $\lambda \sim 4 \times 10^{-56}$   $\phi$  is initially quartic dominated
6. Transitions to  $m_\phi^2 \phi^2$  dominated at  $\sim \text{keV}$  with  $m_\phi^2 \phi^2 \sim 0.2 \text{ eV}^2$
7. Thus  $\phi$  contributes  $\sim 10^{-9}$  of DM

# Ultralight dark matter

1.  $\phi$  needs to transition to matter-like by  $T \sim \text{keV}$

S. Das, E. Nadler [2010.01137](#)

2. Need  $\phi \sim 10^{19}$  eV at  $T \sim \text{keV}$  to get the relic abundance

$10^5$  higher than before

3. Assuming quartic dominates before keV, need  $\phi_i \sim 10^{23}$  eV at  $T_{\text{rh}} \sim 10$  MeV

4. So  $\lambda \sim 3 \times 10^{-66}$  (or smaller)

Fine tuning

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Fine tuning

Can have a shape-shifting sterile neutrino  
that evades cosmology by adding one particle: dark matter

# Shape-shifting sterile neutrinos in vacuum

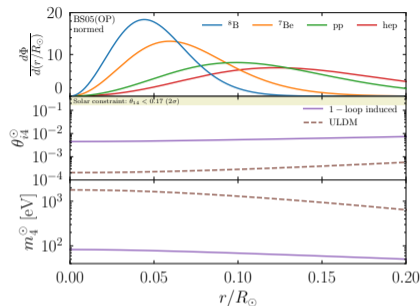
Back to low energy pheno:

- ▶ Vacuum of space:  $m_4 \sim 1$  eV,  $\theta_{14} \sim 0.3$
- ▶ Active neutrinos as expected
- ▶  $\sum m_\nu \lesssim 0.1$  eV comes mostly from  $z \in [10, 100]$

C. Lorenz, et al. [2102.13618](#)

# Shape-shifting sterile neutrinos at the Earth

1. Earth's surface: things are nearly the same as vacuum
2. Center of sun:  $m_4 \sim 10^3$  eV,  $\theta_{14} \sim 3 \times 10^{-4}$



## Other phenomena of shape-shifting sterile neutrinos

- ▶  $\nu_s$ 's will be resonantly produced in the early universe in small bursts as  $\phi$  oscillates past 0
  - ▶ Effect is small
- ▶ The sterile neutrino is too heavy to affect supernova dynamics
- ▶ The Sun's potential could lead to an annual (and daily) modulation in sterile signals
  - ▶ Depends on  $m_\phi$  which is flexible
  - ▶ No such search has been performed
- ▶ Could lead to a modification of atmospheric constraints on steriles
- ▶ Doesn't address surface constraints e.g. reactor spectral, KARMEN, etc.

# 1 eV sterile summary

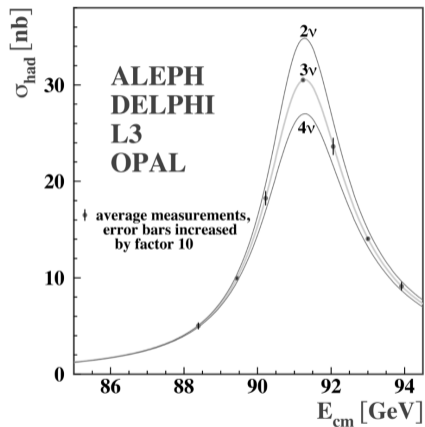
- ▶ Hints for  $\sim 1$  eV steriles persist
  - ▶ RAA is essentially gone
  - ▶ Gallium is back
- ▶ Constraints for  $\sim 1$  eV steriles persist
- ▶ Cosmological constraints are strong and robust
  - ▶ Maybe Hubble parameter tension?
  - ▶ Testable with IceCube upgrade
- ▶ MicroBooNE does not see appearance
- ▶ MicroBooNE might be seeing disappearance
  - ▶ Consistent with gallium
  - ▶ Inconsistent with other constraints
- ▶ Possible to evade cosmology with: 1 sterile neutrino and ultra-light DM

Thanks!



# Backups

Any new light neutrinos must be sterile: SM gauge singlets



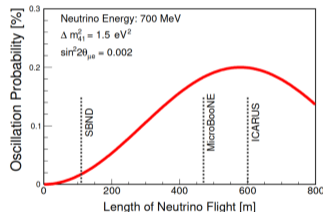
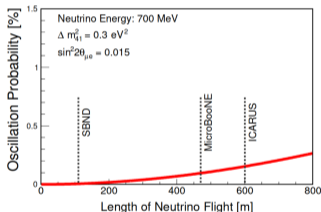
Fun fact: pre-LEP upper limit on  $N_\nu \sim 6000!$

# Short baseline program

1. Leverage LAr to discriminate photons from electrons

MicroBooNE [1910.02166](#)

2.  $L$  is easier to measure than  $E$



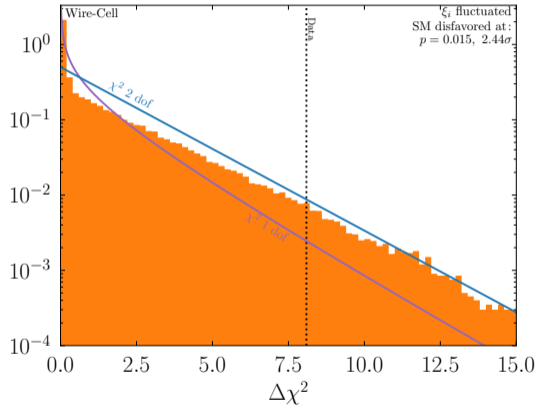
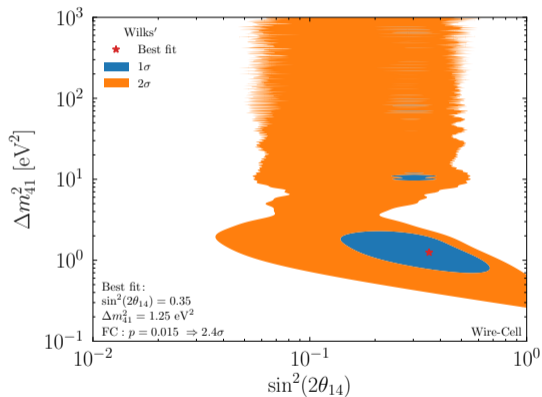
P. Machado, O. Palamara, D. Schmitz [1903.04608](#)

3. Beam is mostly  $\nu_{\mu}$ , but some  $\nu_e$  too
4. Test bed for LAr technology

# Analysis procedure

1. Take systematics as fully uncorrelated bin to bin
2. Unfold predicted spectrum to spectrum in true energy
  - ▶ Use a derivative regulator
3. Apply oscillation probability
4. Reapply energy smearing
5. Compare to data with LLR-Poisson with pull terms
6. Apply Feldman-Cousins
  - ▶ Fluctuate systematics
  - ▶ Literature suggests this is conservative
  - ▶ Verified that it is conservative in this case
7. Get contours via Wilks'
  - ▶ FC contours are very similar

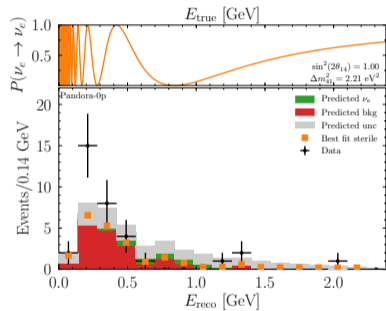
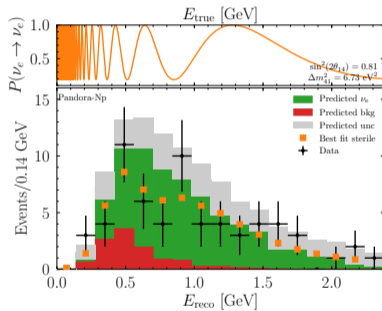
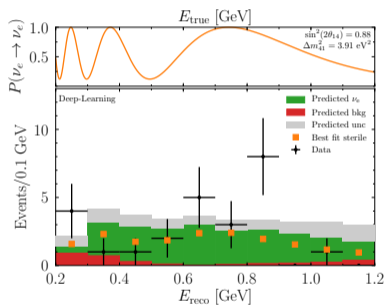
# Results and Monte Carlo significance



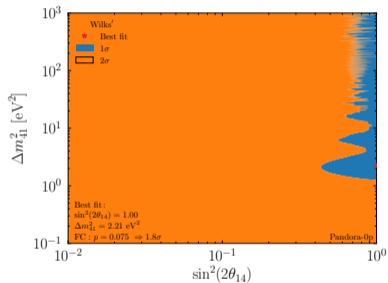
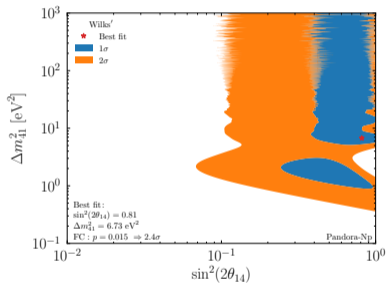
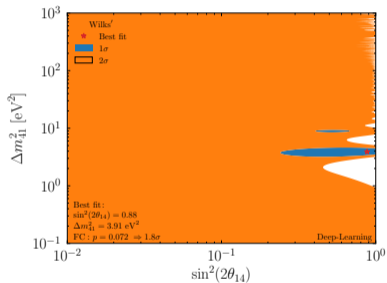
## Other MicroBooNE analysis channels

Analysis	$\sin^2(2\theta_{14})$	$\Delta m_{41}^2$ (eV <sup>2</sup> )	$N\sigma$ (FC)
Wire-Cell	$0.35^{+0.19}_{-0.16}$	$1.25^{+0.74}_{-0.39}$	2.4
Deep-Learning	$0.88^{+0.12}_{-0.41}$	$3.91^{+0.40}_{-0.40}$	1.8
Pandora-Np	$0.81^{+0.19}_{-0.47}$	$[1.28, 2.44]$ $6.73^{+1.75}_{-0.90}$ $\vdots$	2.4
Pandora-0p	$1_{-0.29}$	$2.21^{+0.82}_{-0.60}$ $\vdots$	1.8

# MicroBooNE data in other analyses

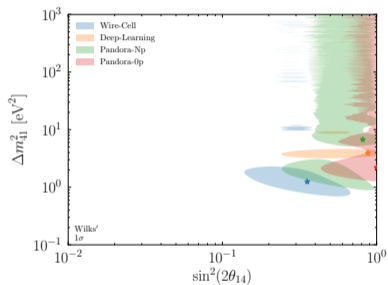


# MicroBooNE contours in other analyses





# MicroBooNE contours in other analyses



# MicroBooNE analyses overlap

Events in multiple analyses:

Analysis	W-C	D-L	Pan-Np	Pan-0p
Wire-Cell	606	15	45	7
Deep-Learning	15	25	9	0
Pandora-Np	45	9	64	0
Pandora-0p	7	0	0	35

# Unitarity constraints

Unitary violation: the study of how  $U_{3\times 3}$  is not unitary independent of  $m_4, m_5, \dots$   
Constraints vary considerably among “global” analyses:

$$1 - |U_{e1}|^2 - |U_{e2}|^2 - |U_{3e}|^2 < \begin{cases} 0.05 \\ 0.001 \end{cases} \quad \text{at } 2\sigma$$

S. Parke, M. Ross-Lonergan [1508.05095](#)

Z. Hu, et al. [2008.09730](#)

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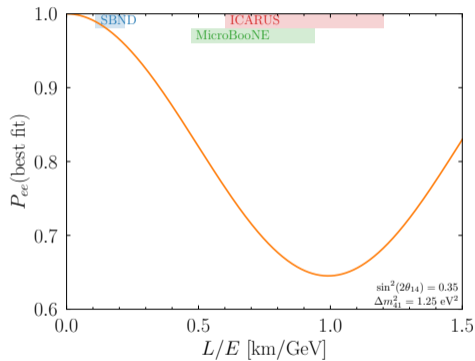
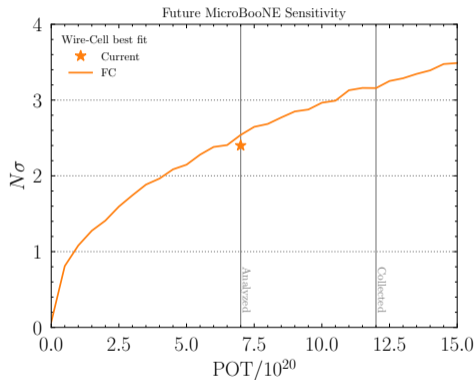
$$1 - |U_{e1}|^2 - |U_{e2}|^2 - |U_{3e}|^2 < \begin{cases} 0.05 \\ 0.001 \end{cases} \quad \text{at } 2\sigma$$

All analyses *assume* unitarity  
Throw out LSND, MiniBooNE, RAA, gallium, etc.

S. Parke, M. Ross-Lonergan [1508.05095](#)

Z. Hu, et al. [2008.09730](#)

# To the future



## Other analyses

- ▶ Evidence for appearance is still there with MiniBooNE, but lower significance
- ▶ Don't see  $> 2\sigma$  evidence for disappearance but very similar best fit

C. Argüelles, et al. [2111.10359](#)

- ▶ Evidence for appearance is still there, but lower significance

MiniBooNE [2201.01724](#)

- ▶ Analysis depends on whether focused on disappearance or both
- ▶ Others handle fully-/partially-contained better
- ▶ Also doesn't see high evidence for disappearance

MicroBooNE [2210.10216](#)

None discuss cosmological constraints

# 3+1+NSI

A new interaction can mitigate IceCube constraints

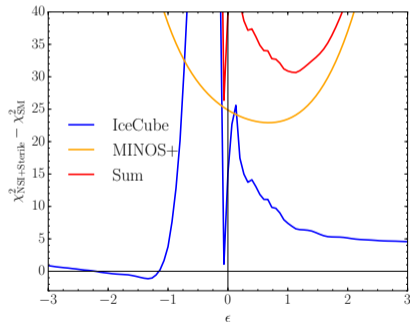
$\epsilon_{\mu\mu}, \epsilon_{\tau\tau}$ : J. Liao, D. Marfatia [1602.08766](#)

Can it also help with MINOS?

## A new interaction can mitigate IceCube constraints

 $\epsilon_{\mu\mu}, \epsilon_{\tau\tau}$ : J. Liao, D. Marfatia [1602.08766](#)

Can it also help with MINOS?



- ▶ Built UV complete model with  $\epsilon_{SS}$
- ▶ IceCube: 3+1+NSI is preferred over SM
- ▶ MINOS: No preference for 3+1 even with NSI

PBD, Y. Farzan, I. Shoemaker [1811.01310](#)



# Self interacting long range forces sourced by celestial objects

1. A general long range force will have a self interaction
2. Need to solve:

$$-2\frac{\phi'}{r} - \phi'' + m^2\phi + \frac{1}{3!}\lambda\phi^3 + n(r) = 0$$

3. Analytic solutions don't exist, and is a boundary problem

$$\phi'(0) = 0, \phi(\infty) = 0$$

4. Need to employ shooting method to infer  $\phi(0)$
5. Depending on shape of  $n(r)$  and value of  $\lambda$ , can hit double precision limit

