Abstract

The nature of CP violation in the lepton sector is one of the biggest open questions in particle physics. Long-baseline accelerator neutrino experiments have the opportunity to determine if CP is violated in the mass matrix. I will look at the most recent NOvA and T2K data which show a slight and very interesting tension. While this tension possibly indicates a flipping in the mass ordering, it is better fit by new physics such as non-standard neutrino interactions (NSI) with an additional source of CP violation. The strength of this NSI can be easily estimated analytically and I will present a numerical analysis of the preferred regions which are generally consistent with other constraints.

CP Violation at Long-Baseline Neutrino Experiments

Peter B. Denton

Mitchell Conference

May 27, 2022

Brookhaven National Laboratory

Brookhaven NDI
Neutrino Discovery Initiative
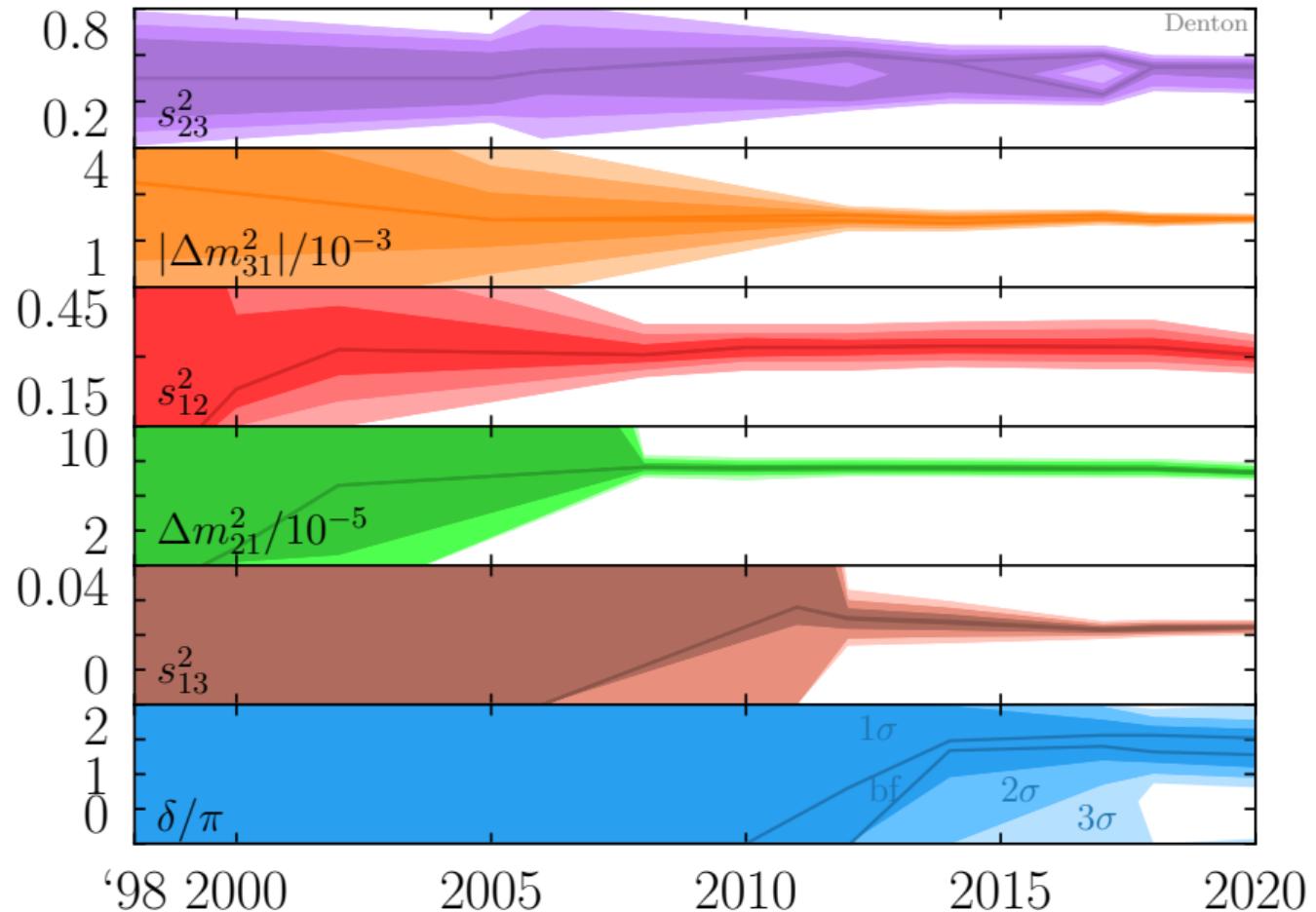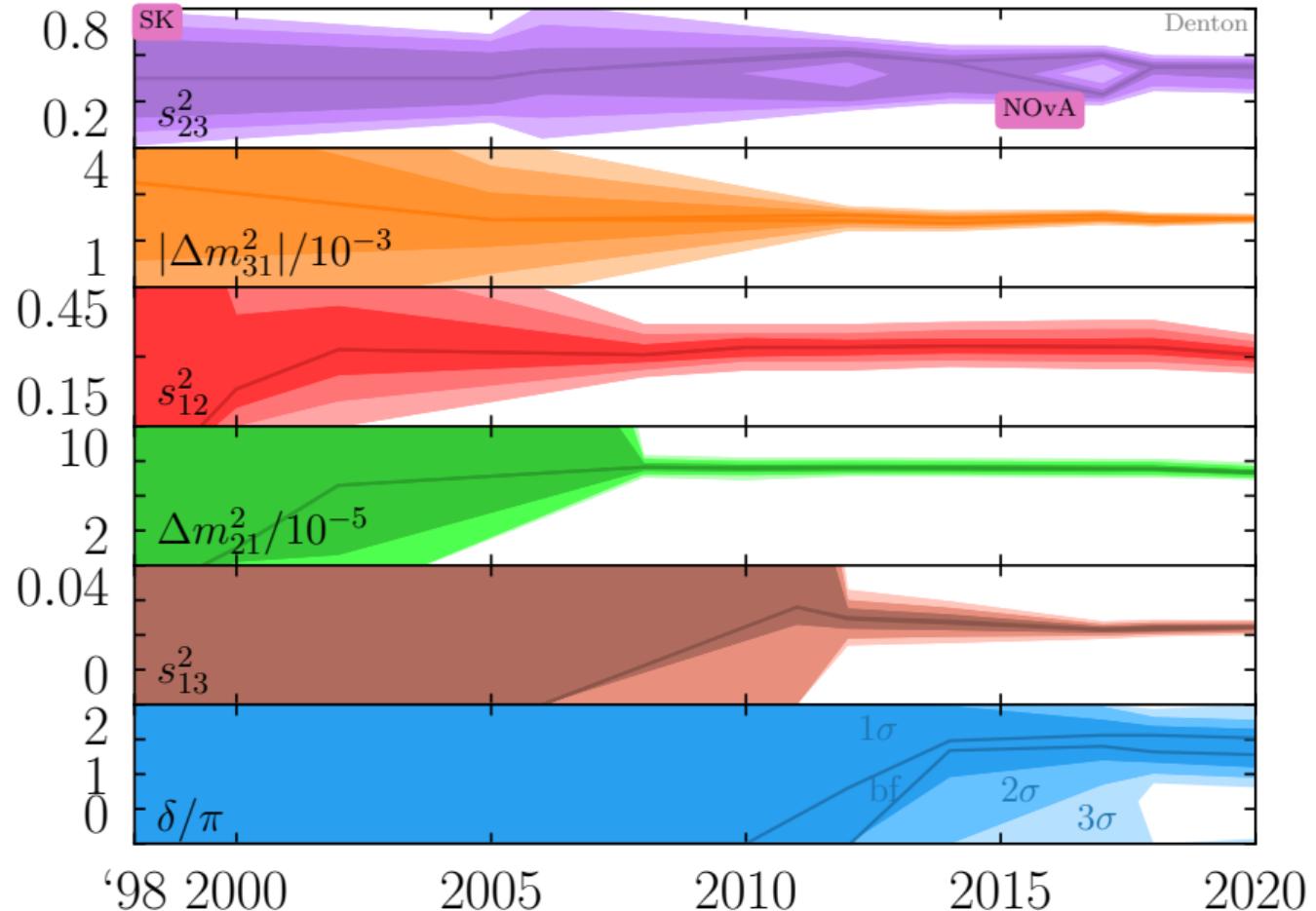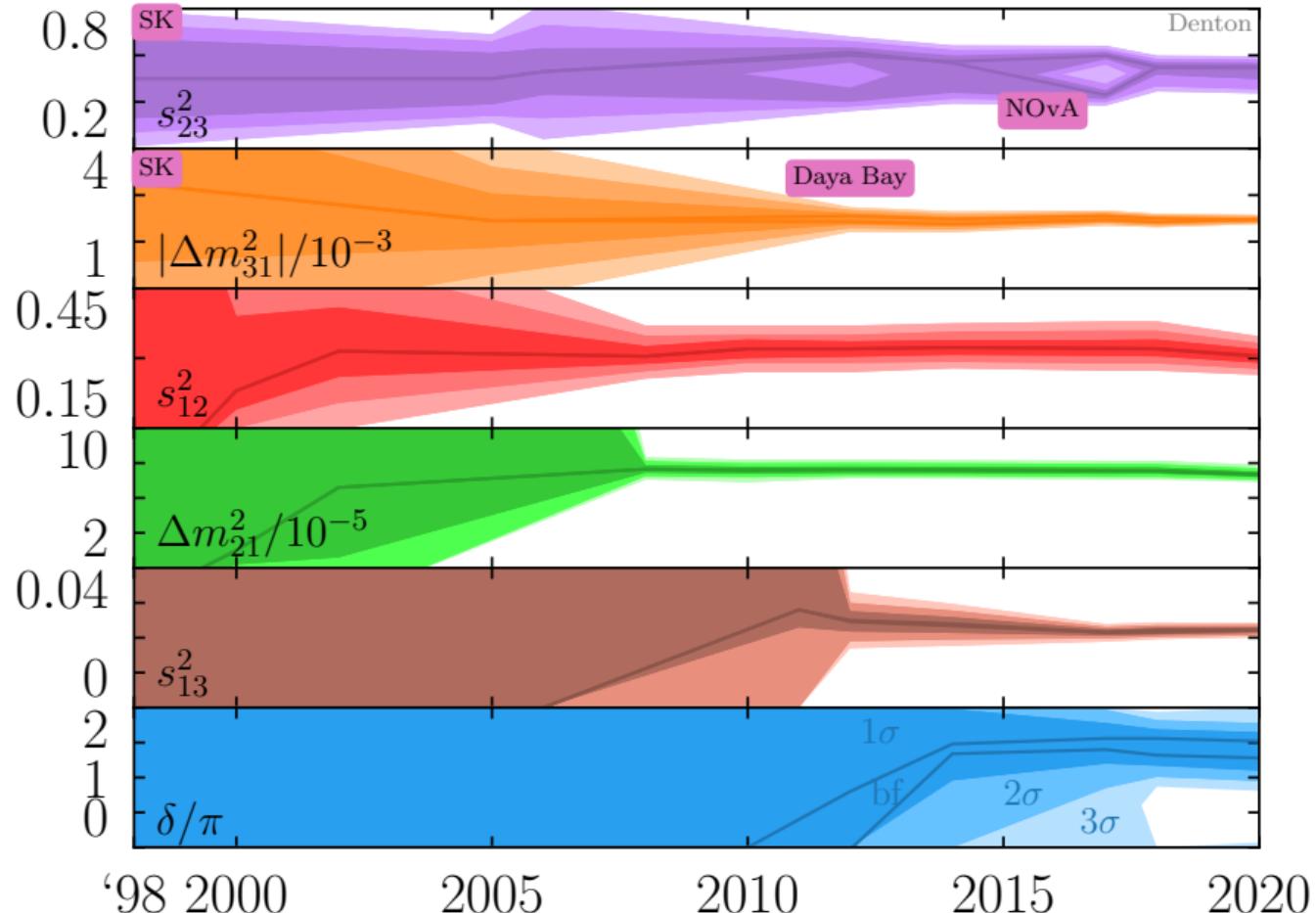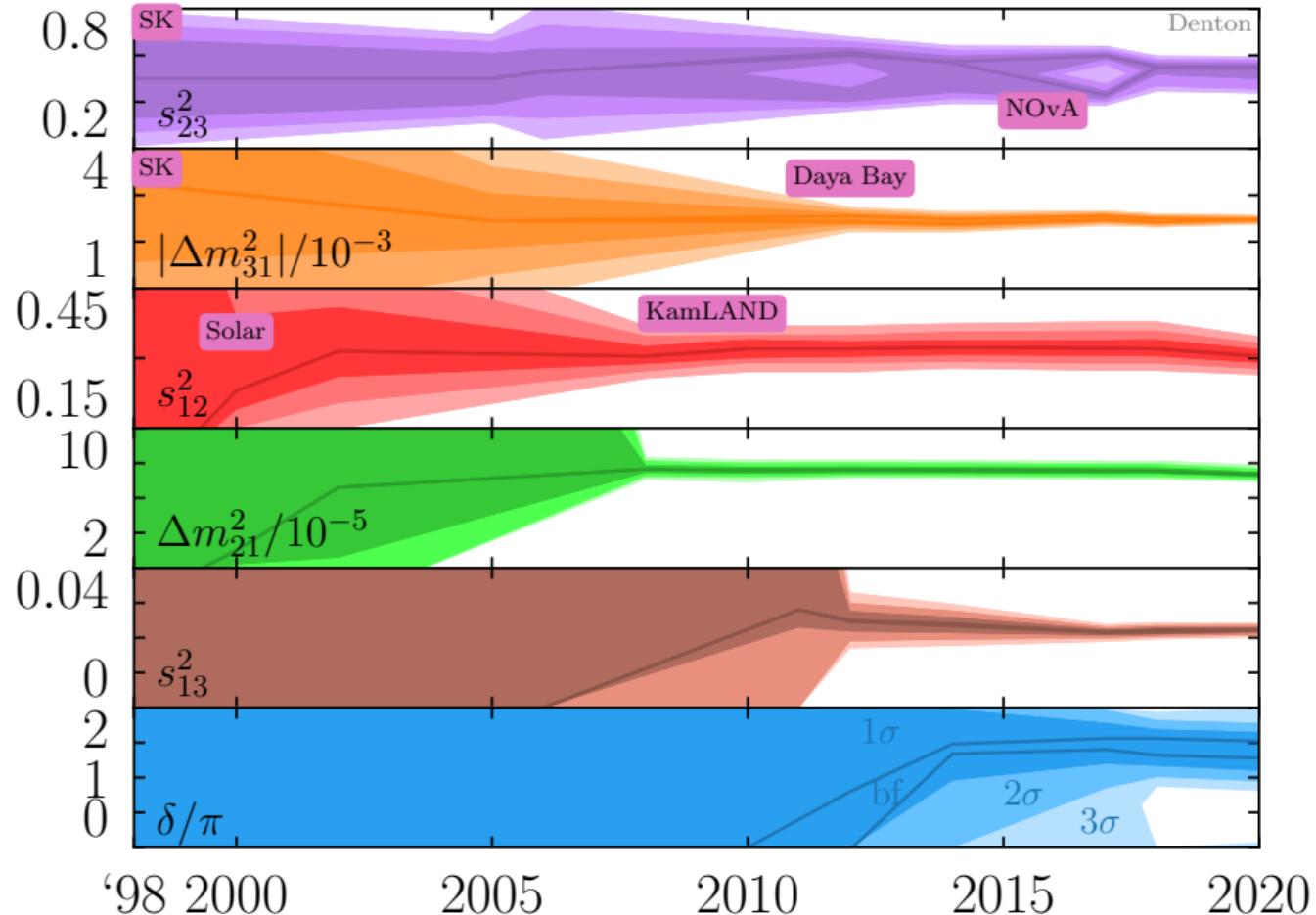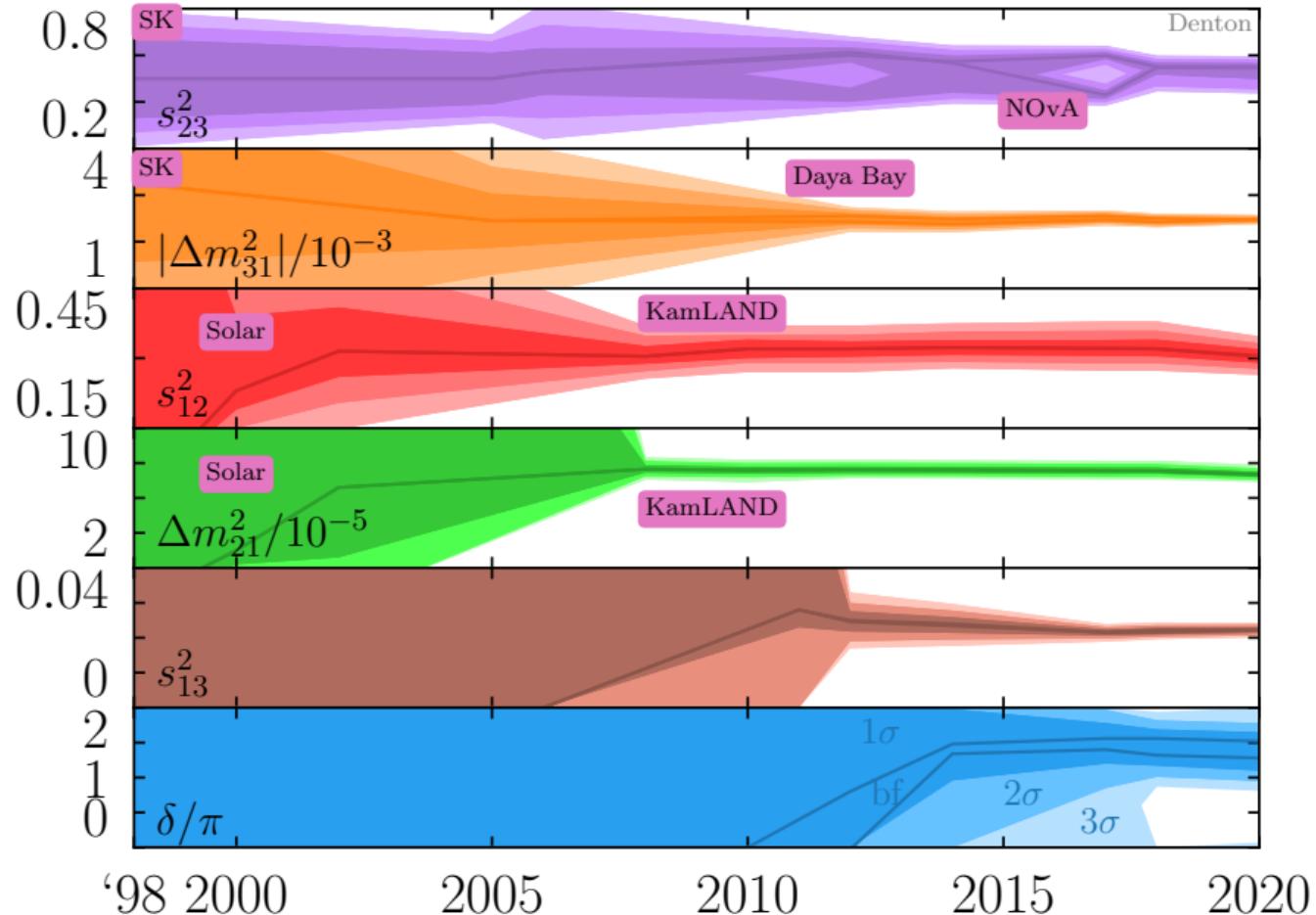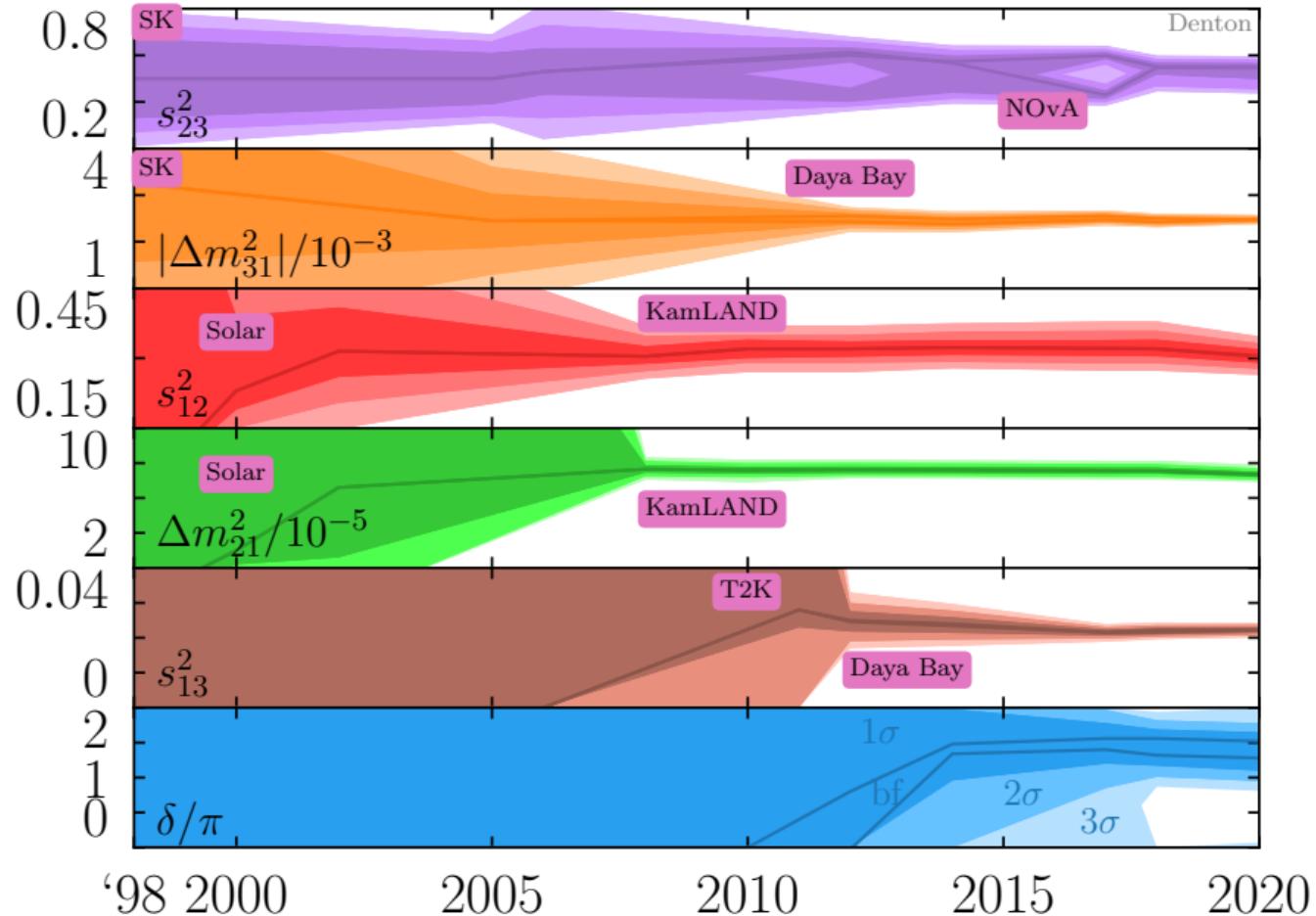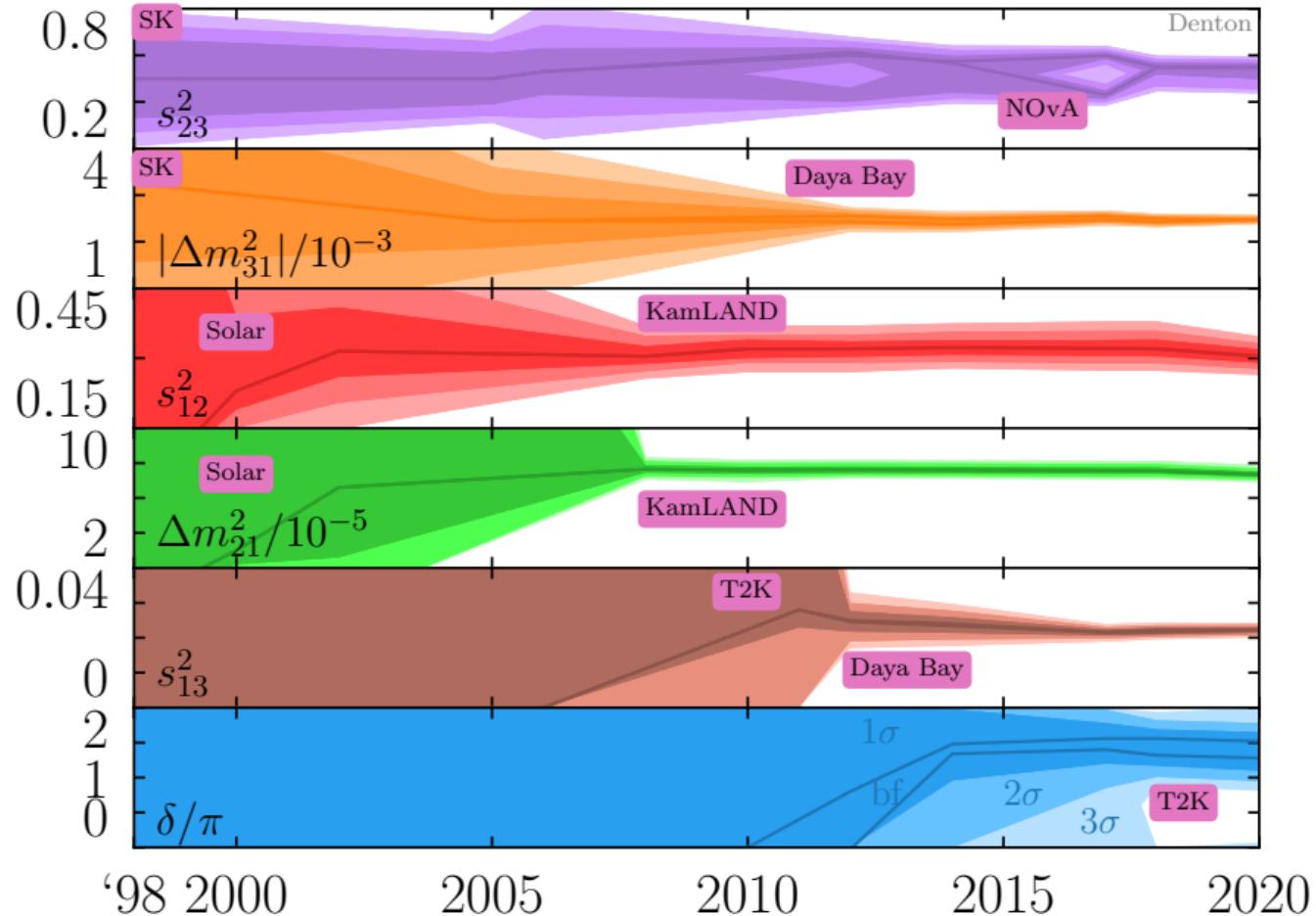
Atmospheric mass ordering

θ_{23} octant

Complex phase

Absolute mass scale

Atmospheric mass ordering

θ_{23} octant

See J. Gehrlein's talk yesterday

Complex phase

Absolute mass scale

Overview

- δ vs. CPV
- Different parameterizations lead to different conclusions

- NOvA and T2K slightly disagree
- New physics can resolve this

δ and CP violation

$$J_{CP} = s_{12}c_{12}s_{13}c_{13}^2 s_{23}c_{23}\sin\delta$$

C. Jarlskog PRL 55, 1039 (1985)

δ and CP violation

$$J_{CP} = s_{12}c_{12}s_{13}c_{13}^2 s_{23}c_{23} \sin\delta$$

<div align="right">C. Jarlskog PRL 55, 1039 (1985)</div>

1. Strong interaction: no observed EDM $\Rightarrow$ CP (nearly) **conserved**

$$\frac{\bar{\theta}}{2\pi} < 10^{-11}$$

<div align="right">J. Pendlebury, et al. 1509.04411</div>

2. Quark mass matrix: non-zero but **small** CP violation

$$\frac{|J_{\mathrm{CKM}}|}{J_{\mathrm{max}}} = 3 \times 10^{-4}$$

<div align="right">CKMfitter 1501.05013</div>

3. Lepton mass matrix: **?**

$$\frac{|J_{\mathrm{PMNS}}|}{J_{\mathrm{max}}} < 0.34$$

<div align="right">PBD, J. Gehrlein, R. Pestes 2008.01110</div>

$$J_{\mathrm{max}} = \frac{1}{6\sqrt{3}} \approx 0.096$$

δ: what is it really?
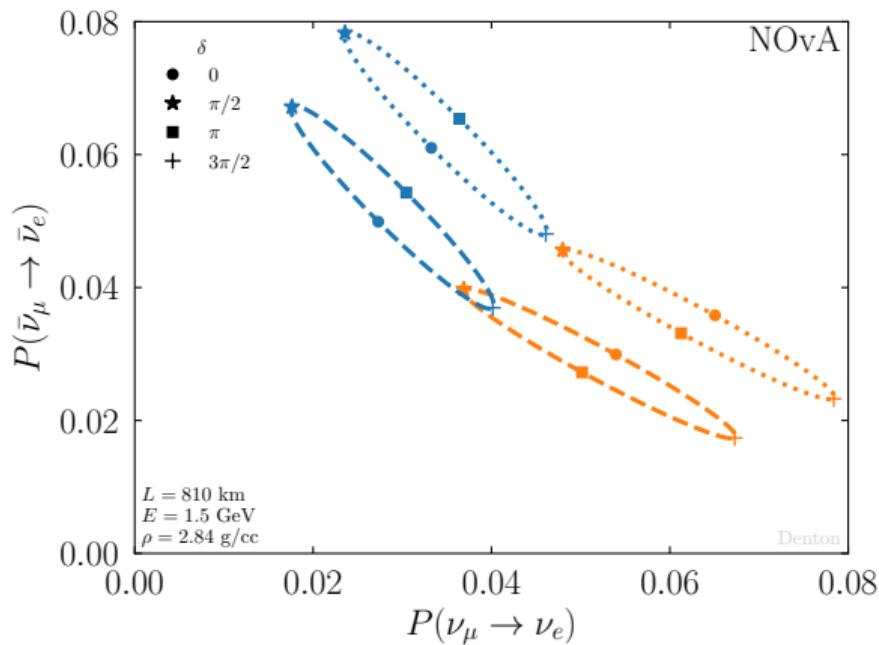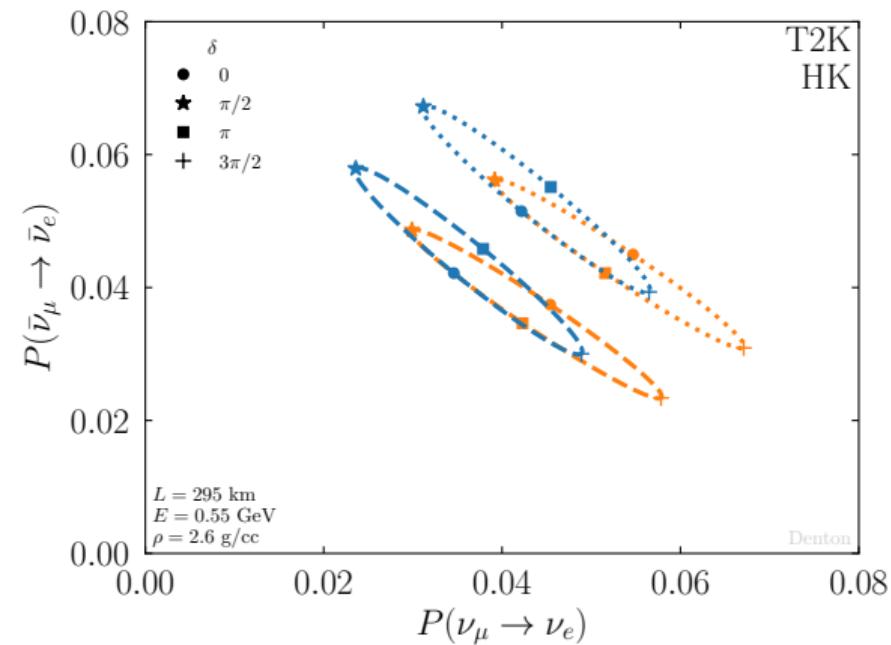
$\delta \neq$ Baryogenesis

The amount of leptogenesis is a function of:

1. the heavy mass scale
2. δ
3. α, β (Majorana phases)
4. CP phases in the RH neutrinos
5. ...

C. Hagedorn, et al 1711.02866

K. Moffat, et al 1809.08251

Measuring $\delta = 0, \pi$	$\neq$	no leptogenesis
Measuring $\delta \neq 0, \pi$	$\neq$	leptogenesis

δ, J: current status

Maximal CP violation is already ruled out:

1. $\theta_{12} \neq 45°$ at $\sim 15\sigma$
2. $\theta_{13} \neq \tan^{-1} \frac{1}{\sqrt{2}} \approx 35°$ at many (100) σ
3. $\theta_{23} = 45°$ allowed at $\sim 1\sigma$
4. $|\sin \delta| = 1$ allowed

$J_{max} = \frac{1}{6\sqrt{3}} \approx 0.096$

Unitarity 100% allowed

+SNO +KamLAND 91% allowed

+Daya Bay +RENO 34% allowed

+NOvA +T2K 26% allowed

CP conserving

Denton 2020
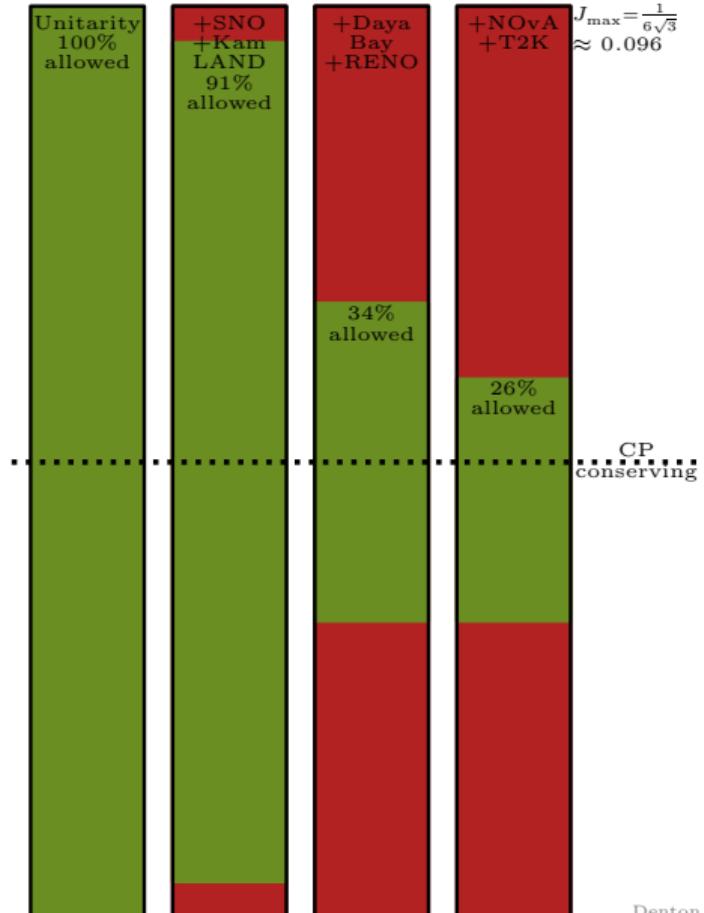
Parameterization of the PMNS matrix

A matrix takes us from mass states to flavor states and back

1. 3×3 $\mathbb{C}$: 18 dof
2. +Unitary: n^2 constraints: 9 dof
3. +Charged lepton rephasing: 6 dof
4. +Neutrino rephasing: 4 dof

Focused on oscillations not $0\nu\beta\beta$

Parameterization of the PMNS matrix

Many possible parameterizations in the literature

1. Product of three rotations and a complex phase on one rotation
   - ▶ Possibly including the same axis twice

   H. Fritzsch, Z.-z. Xing hep-ph/0103242

2. Gell-Mann matrices

   K. Merfeld, D. Latimer 1412.2728

   D. Boriero, D. Schwarz, H. Velten 1704.06139

   A. Davydova, K. Zhukovsky PAN 82, 281 (2019)

3. Four complex phases

   R. Aleksan, B. Kayser, D. London hep-ph/9403341

4. Perturbative

   L. Wolfenstein PRL 51 1945 (1983)

5. ⋮

Sequence of rotations

$$U_1 \equiv \begin{pmatrix} 1 & 0 & 0 \\ 0 & c_{23} & s_{23} \\ 0 & -s_{23} & c_{23} \end{pmatrix} \qquad U_2 \equiv \begin{pmatrix} c_{13} & 0 & s_{13}e^{-i\delta} \\ 0 & 1 & 0 \\ -s_{13}e^{i\delta} & 0 & c_{13} \end{pmatrix} \qquad U_3 \equiv \begin{pmatrix} c_{12} & s_{12} & 0 \\ -s_{12} & c_{12} & 0 \\ 0 & 0 & 1 \end{pmatrix}$$

Location of $e^{i\delta}$ on $\pm s_{ij}$ has no impact*

Standard parameterization is $U_{\text{PDG}} \equiv U_{123} = U_1 U_2 U_3$.

$$U_{\text{PDG}} \equiv U_{123} = \begin{pmatrix} c_{12}c_{13} & s_{12}c_{13} & s_{13}e^{-i\delta} \\ -s_{12}c_{23} - c_{12}s_{13}s_{23}e^{i\delta} & c_{12}c_{23} - s_{12}s_{13}s_{23}e^{i\delta} & c_{13}s_{23} \\ s_{12}s_{23} - c_{12}s_{13}c_{23}e^{i\delta} & -c_{12}s_{23} - s_{12}s_{13}c_{23}e^{i\delta} & c_{13}c_{23} \end{pmatrix}$$

Sequence of rotations

$$U_1 \equiv \begin{pmatrix} 1 & 0 & 0 \\ 0 & c_{23} & s_{23} \\ 0 & -s_{23} & c_{23} \end{pmatrix} \qquad U_2 \equiv \begin{pmatrix} c_{13} & 0 & s_{13}e^{-i\delta} \\ 0 & 1 & 0 \\ -s_{13}e^{i\delta} & 0 & c_{13} \end{pmatrix} \qquad U_3 \equiv \begin{pmatrix} c_{12} & s_{12} & 0 \\ -s_{12} & c_{12} & 0 \\ 0 & 0 & 1 \end{pmatrix}$$

Location of $e^{i\delta}$ on $\pm s_{ij}$ has no impact[*]

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What about other orders?

$$U_{123}, U_{132}, U_{213}, U_{231}, U_{312}, U_{321}$$

What about repeated rotations?

$$U_{121}, U_{131}, U_{212}, U_{232}, U_{313}, U_{323}$$

Complex phase in different parameterizations

- Can relate the complex phase in one parameterization to that in another
- U_{132} and U_{213} similar to U_{123}
- δ constrained to $\sim [150°, 210°]$ in $U_{231}, U_{312}, U_{321}$
- Bands indicate 3σ uncertainty on $\theta_{12}, \theta_{13}, \theta_{23}$
- "50% of possible values of δ"
  $\Rightarrow$ parameterization dependent

DUNE TDR II 2002.03005
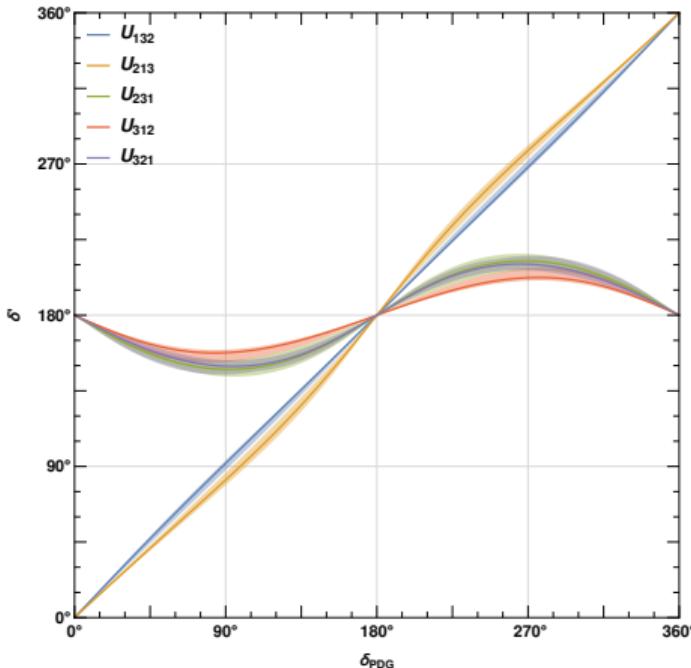
Repeated rotations in backups

Quark mixing

From the PDG, V_{CKM} in the V_{123} parameterization is

$$\theta_{12} = 13.09° \quad \theta_{13} = 0.2068° \quad \theta_{23} = 2.323° \quad \delta_{\text{PDG}} = 68.53°$$

Looks like "large" CPV:

$$\sin \delta_{\text{PDG}} = 0.93 \sim 1$$

yet $J_{\text{CKM}}/J_{\text{max}} = 3 \times 10^{-4}$.

Switch to V_{212} parameterization, $\Rightarrow \delta' = 1°$ and $\sin \delta' = 0.02$.

One caveat in support of δ

If the goal is **CP violation** the Jarlskog should be used

however

If the goal is **measuring the parameters** one must use δ

Given θ_{12}, θ_{13}, θ_{23}, and J, I can't determine the sign of $\cos\delta$ which is physical

e.g. $P(\nu_\mu \to \nu_\mu)$ depends on $\cos\delta$ a tiny bit

- ▶ T2K/HK are mostly sensitivity to $\sin\delta$; they should focus on J
- ▶ NOvA/DUNE has modest $\cos\delta$ sensitivity; both J and δ should be reported

The importance of $\cos\delta$

- If only $\sin\delta$ is measured $\Rightarrow$ sign degeneracy: $\cos\delta = \pm\sqrt{1 - \sin^2\delta}$

- Most models predict $\cos\delta$
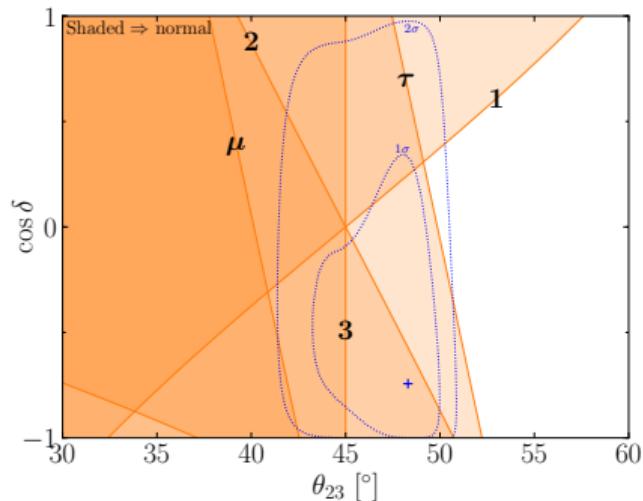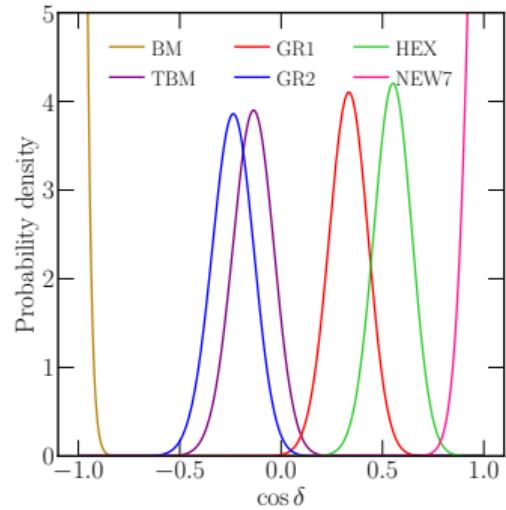
CP violation at NOvA and T2K?

Excitement at Neutrino2020!
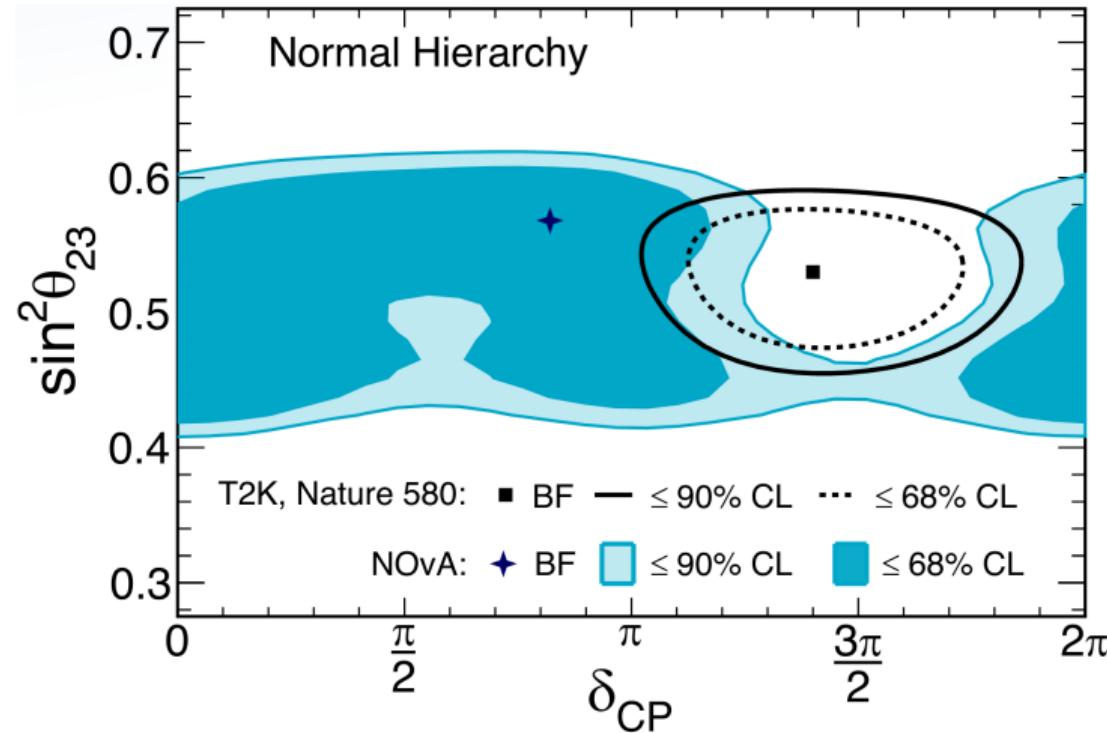
New physics

If this is new physics what could lead to this kind of effect?

- ▶ Steriles?   S. Chatterjee, A. Palazzo 2005.10338
- ▶ Decay?
- ▶ Decoherence?
- ▶ Dark matter interaction?
- ▶ LIV/CPT?   U. Rahaman 2103.04576
- ▶ Unitary violation?   D. Forero, et al. 2103.01998
- ▶ NSI with complex CP violating phases
  1. Different matter effects ⇒ different NSI effect
  2. New phases partially degenerate with standard phase
  3. T2K is closer to vacuum so they measure the vacuum parameters
  4. NOvA measures "vacuum" + "NSI"

See also S. Chatterjee, A. Palazzo 2008.04161

NSI review

$$\mathcal{L}_{\text{NSI}} = -2\sqrt{2}G_F \sum_{\alpha,\beta,f,P} \epsilon_{\alpha\beta}^{f,P} (\bar{\nu}_\alpha \gamma^\mu \nu_\beta)(\bar{f}\gamma_\mu f)$$

Models with large NSIs consistent with CLFV:

Y. Farzan, I. Shoemaker 1512.09147   Y. Farzan, J. Heeck 1607.07616   D. Forero and W. Huang 1608.04719
K. Babu, A. Friedland, P. Machado, I. Mocioiu 1705.01822   PBD, Y. Farzan, I. Shoemaker 1804.03660
U. Dey, N. Nath, S. Sadhukhan 1804.05808   Y. Farzan 1912.09408

Affects oscillations via new matter effect

$$H = \frac{1}{2E}\left[UM^2U^\dagger + a \begin{pmatrix} 1+\epsilon_{ee} & \epsilon_{e\mu} & \epsilon_{e\tau} \\ \epsilon_{e\mu}^* & \epsilon_{\mu\mu} & \epsilon_{\mu\tau} \\ \epsilon_{e\tau}^* & \epsilon_{\mu\tau}^* & \epsilon_{\tau\tau} \end{pmatrix} \right]$$

Matter potential $a \propto G_F \rho E$

B. Dev, K. Babu, PBD, P. Machado, et al. 1907.00991

Estimate size of effect: magnitude

$$|\epsilon_{e\beta}| \approx \frac{s_{12}c_{12}c_{23}\pi\Delta m_{21}^2}{2s_{23}w_\beta} \left| \frac{\sin\delta_{\text{T2K}} - \sin\delta_{\text{NOvA}}}{a_{\text{NOvA}} - a_{\text{T2K}}} \right| \approx \begin{cases} 0.22 & \text{for } \beta = \mu \\ 0.24 & \text{for } \beta = \tau \end{cases}$$

$a \propto \rho E$

$w_\beta = s_{23}, c_{23}$ for $\beta = \mu, \tau$

Assumed upper octant $\theta_{23} > 45°$

Consistency checks:

▶ $\sin\delta_{\text{NOvA}} = \sin\delta_{\text{T2K}} \Rightarrow |\epsilon| = 0$

▶ $\sin\delta_{\text{NOvA}} \neq \sin\delta_{\text{T2K}}$ and $a_{\text{NOvA}} = a_{\text{T2K}} \Rightarrow |\epsilon| \to \infty$

▶ Octant:
  1. LBL is governed by ν_3
  2. Upper octant $\Rightarrow \nu_3$ is more ν_μ
  3. More $\nu_\mu \Rightarrow$ need less new physics coupling to ν_μ to produce a given effect

Approximate the experiments

Appearance:

$$n(\nu_e) = xP(\nu_\mu \to \nu_e) + yP(\bar{\nu}_\mu \to \bar{\nu}_e) + z$$

Fit to all points on bievent plots for ν, $\bar{\nu}$, NOvA, T2K

Wrong sign leptons are non-zero at high significance

Disappearance:

NOvA:

$$|\Delta m_{32}^2| = (2.41 \pm 0.07) \times 10^{-3} \text{ eV}^2 \quad \text{and} \quad 4|U_{\mu3}|^2(1 - |U_{\mu3}|^2) = 0.99 \pm 0.02$$

K. Kelly, et al. 2007.08526

T2K: Δm_{32}^2 and θ_{23} likelihoods

Assume that $P_{\mu\mu} \approx \bar{P}_{\mu\mu}$ and that most info comes from disappearance

NOvA: $E \sim 1.9$ GeV, $\rho = 2.84$ g/cc, $L = 810$ km

T2K: $E \sim 0.6$ GeV, $\rho = 2.60$ g/cc, $L = 295$ km

Completing the fit

Appearance measurements depend on all six parameters
Can NOT use input from global fits

T2K affects θ_{13}

Use other vacuum experiments to constrain other parameters independent of NSI:

▶ Daya Bay: Constrains θ_{13} and Δm_{32}^2 for each atmospheric mass ordering

Daya Bay 1809.02261

▶ KamLAND: Constrains θ_{12} and $|\Delta m_{21}^2|$

KamLAND 1303.4667

SNO tells us $\Delta m_{21}^2 > 0$

or $\theta_{12} < 45°$ depending on definition, see PBD 2003.04319

This depends on NSI but LBL parameters don't cancel

NSI parameters
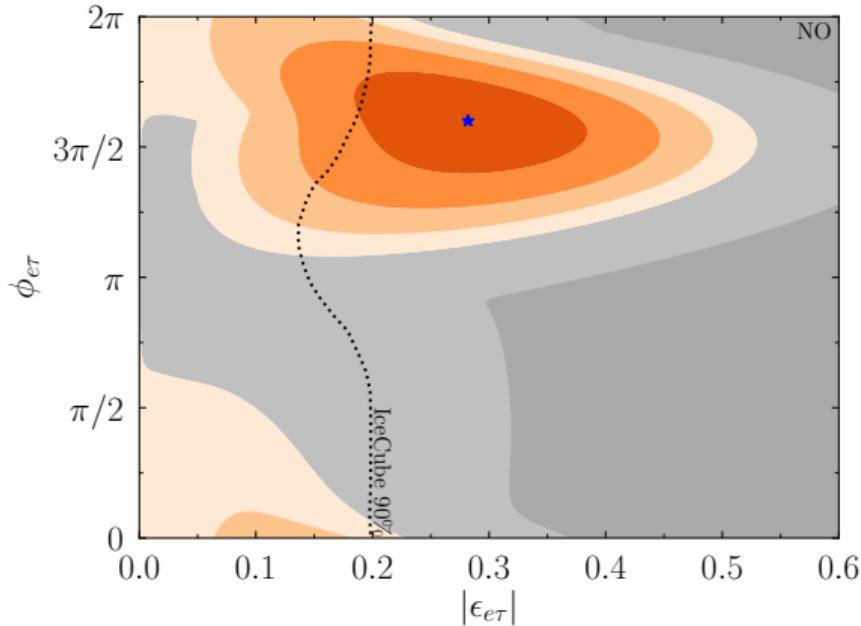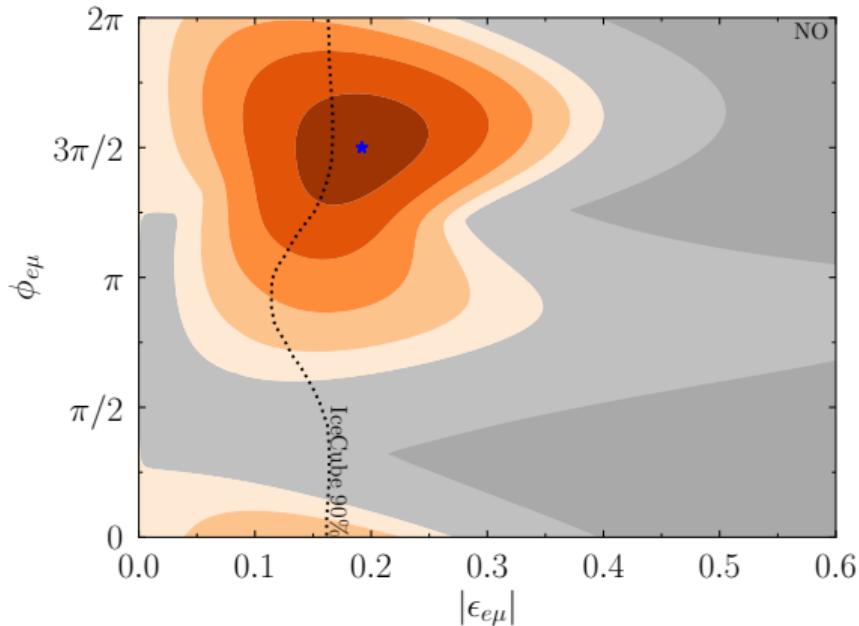
Orange is preferred over SM at integer values of $\Delta\chi^2$, dark gray is disfavored at 4.61

T. Ehrhardt, IceCube PPNT (2019)

$\epsilon_{\mu\tau}$, IO in backups

Other CP violating NSI constraints

NSI effects grow with energy, density, and distance

Best probes:

- ▶ $\epsilon_{\mu\tau}$: atmospheric
- ▶ $\epsilon_{e\mu}$, $\epsilon_{e\tau}$: LBL appearance, atmospheric

- ▶ IceCube
  - ▶ Constraint is at LBL best fit with 3 yrs
    10 yrs of data in the bank
  - ▶ Prefers non-zero $|\epsilon_{e\mu}|$ at $\sim 1\sigma$

- ▶ Super-K
  - ▶ Only consider real NSI
  - ▶ Comparable sensitivity as IceCube

- ▶ COHERENT
  - ▶ Only applies to NSI models with $M_{Z'} \gtrsim 10$ MeV
  - ▶ NSI u, d, e configuration matters
  - ▶ Comparable constraints
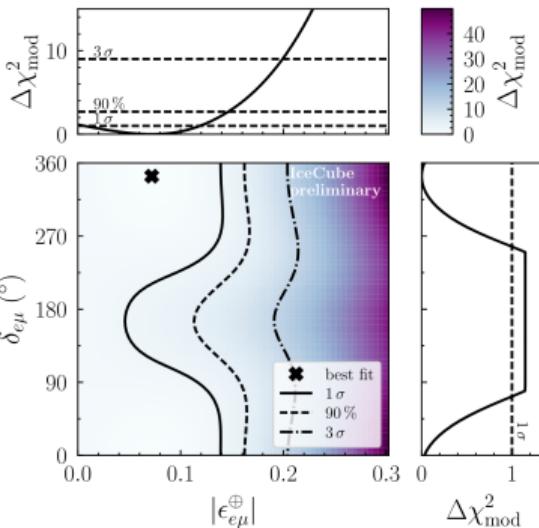
T. Ehrhardt, IceCube PPNT (2019)

Super-K 1109.1889

COHERENT 1708.01294
PBD, Y. Farzan, I. Shoemaker 1804.03660
PBD, J. Gehrlein 2008.06062

Summary

- Care is required in choice of parameterizations
- Jarlskog invariant is best for CP violation

- NOvA and T2K tension can be mitigated by NO $\to$ IO
- Tension can be fully resolved by NSI
- Easy to approximate magnitude and phase of NSI
- NSI introduces more CP violation
- Consistent with, and soon tested by, other experiments

Thanks!

Rice & Texas A&M:

Backups

References
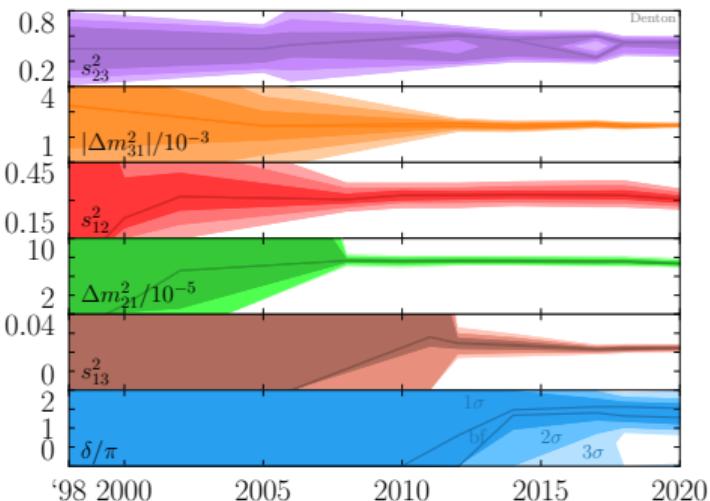
SK `hep-ex/9807003`

M. Gonzalez-Garcia, et al. `hep-ph/0009350`

M. Maltoni, et al. `hep-ph/0207227`

SK `hep-ex/0501064`

SK `hep-ex/0604011`

T. Schwetz, M. Tortola, J. Valle `0808.2016`

M. Gonzalez-Garcia, M. Maltoni, J. Salvado `1001.4524`

T2K `1106.2822`

D. Forero, M. Tortola, J. Valle `1205.4018`

D. Forero, M. Tortola, J. Valle `1405.7540`

P. de Salas, et al. `1708.01186`

F. Capozzi et al. `2003.08511`

The importance of $\cos\delta$
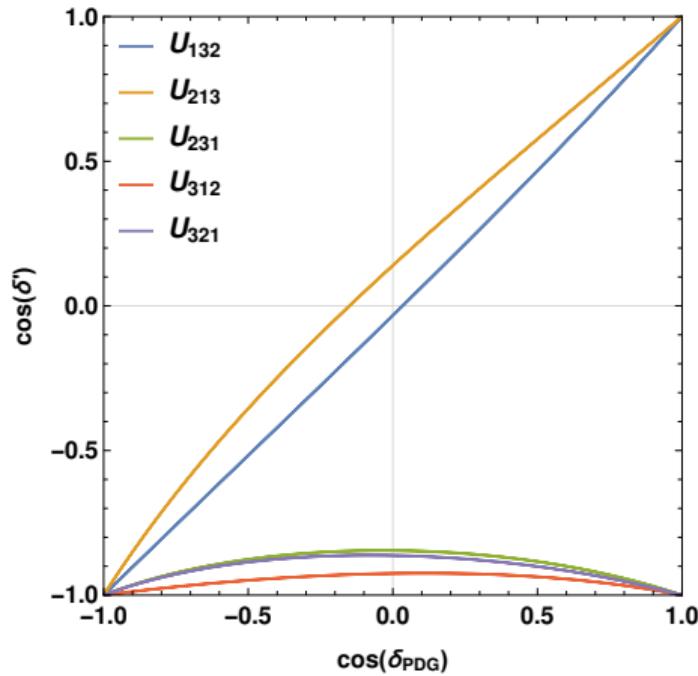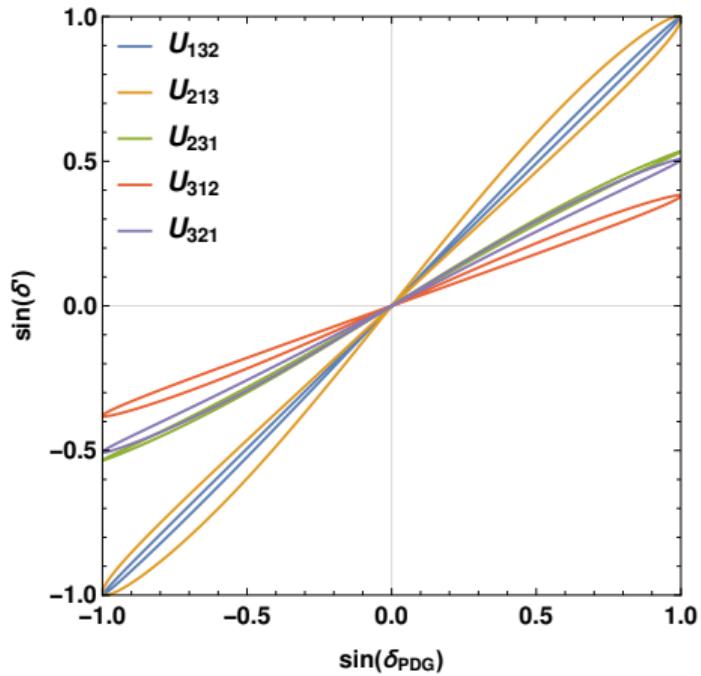
In these parameterizations $\cos\delta \lesssim -0.8$

$|U_{e3}|$ is small

Given θ_{12}, θ_{13}, θ_{23}:

$$|U| = \begin{pmatrix} 0.822 & 0.550 & 0.150 \\ \sqrt{0.138 + 0.068\cos(\delta_{\text{PDG}})} & \sqrt{0.293 - 0.068\cos(\delta_{\text{PDG}})} & 0.754 \\ \sqrt{0.186 - 0.068\cos(\delta_{\text{PDG}})} & \sqrt{0.405 + 0.068\cos(\delta_{\text{PDG}})} & 0.640 \end{pmatrix}$$

$$|U_{\alpha i}| > 0.23 \qquad \text{except} \qquad |U_{e3}| = 0.15$$

In U_{231}, U_{312}, U_{321}:

$$|U_{e3}| = \sqrt{A + B\cos(\delta')}$$

$$A, B > 0$$

Requires a partial cancellation $\Rightarrow \cos(\delta') \sim -1$

Terms with sums or differences are "complicated"

Terms without are "simple"

Quick approximation

Can easily related $\delta_{\mathrm{PDG}} \to \delta'$:

- $\delta' \approx \delta_{\mathrm{PDG}}$ in U_{132} and U_{213}
- $\sin(\delta') \approx d_{ijk} \sin(\delta_{\mathrm{PDG}})$

$$d_{231} \approx s_{13} \frac{1 - s_{12}^2 c_{23}^2}{s_{12} c_{12} s_{23} c_{23}} \approx 0.57$$

$$d_{312} \approx s_{13} \frac{1 - c_{12}^2 s_{23}^2}{s_{12} c_{12} s_{23} c_{23}} \approx 0.39$$

$$d_{321} \approx s_{13} \frac{1 - s_{12}^2 s_{23}^2}{s_{12} c_{12} s_{23} c_{23}} \approx 0.54$$
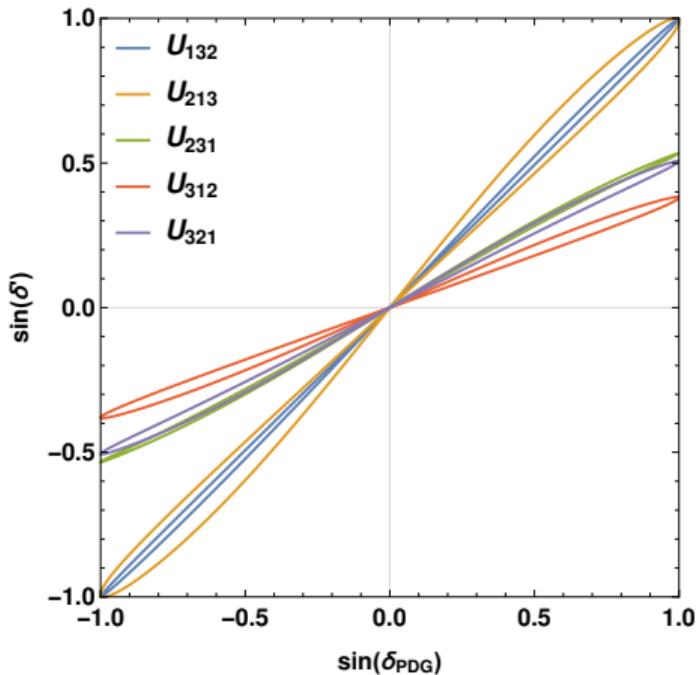
$\theta_{23} > 45°$ plotted

Precision on δ
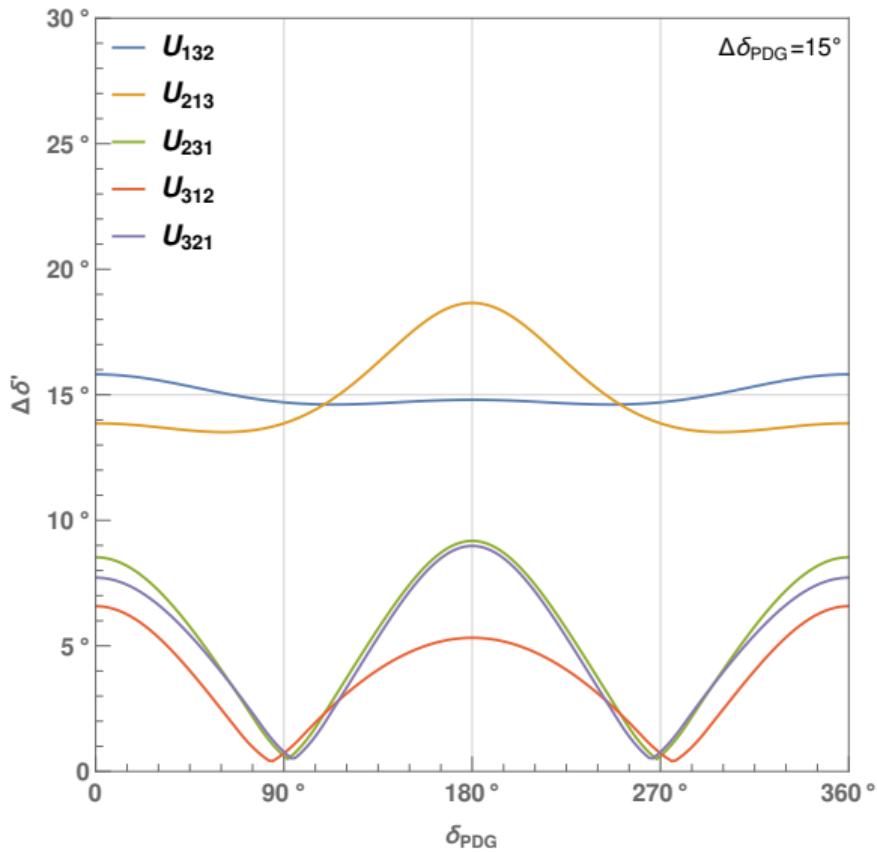
"For instance, the CKM angle γ, which is a very close analog of δ in the neutrino sector, is determined to $70.4^{+4.3}_{-4.4}$ and thus, a precision target for δ of roughly $5°$ would follow."

"A 3σ distinction between models translates into a target precision for δ of $5°$."

A. de Gouvea, et al. Snowmass 2013
Neutrino Working Group 1310.4340

Precision on δ is parameterization dependent

CP violation in oscillations

In vacuum at first maximum:

$$P_{\mu e} - \bar{P}_{\mu e} \approx 8\pi J \frac{\Delta m_{21}^2}{\Delta m_{32}^2}$$

$$J \equiv s_{12} c_{12} s_{13} c_{13}^2 s_{23} c_{23} \sin \delta$$

C. Jarlskog PRL 55, 1039 (1985)

► Extracting δ from data requires every other oscillation parameter
► J requires only Δm_{21}^2 (up to matter effects)

Matter effects are easily accounted for

$$\hat{J} \simeq \frac{J}{\sqrt{(c_{212} - c_{13}^2 a/\Delta m_{21}^2)^2 + s_{212}^2}\sqrt{(c_{213} - a/\Delta m_{ee}^2)^2 + s_{213}^2}}$$

PBD, S. Parke 1902.07185

PBD, H. Minakata, S. Parke 1604.08167

Optimal Parameterization

Want to be able to write

$$P \approx \sin^2 2\theta \sin^2 \left(\frac{\Delta m^2 L}{4E} \right)$$

1. Solar/long-baseline reactor: U_{e2}
2. Medium-baseline reactor: U_{e3}
3. Atmospheric/long-baseline accelerator disappearance: $U_{\mu3}$

Want these "simple" not the sum/difference of trig functions

	U_{123}	U_{132}	U_{213}	U_{231}	U_{312}	U_{321}
$|U_{e2}|$	✓	✓	✗	✗	✓	✗
$|U_{e3}|$	✓	✓	✓	✗	✗	✗
$|U_{\mu3}|$	✓	✗	✓	✓	✗	✗

Other priorities (theoretical, computational, . . .) may prefer different parameterizations

Optimal Parameterization

<div style="text-align: center">

Location of the phase?

</div>

Conventional:

$$U_{23}(\theta_{23})U_{13}(\theta_{13}, \delta)U_{12}(\theta_{12})$$

Sometimes useful when dealing with matter effect:

$$U_{23}(\theta_{23}, \delta)U_{13}(\theta_{13})U_{12}(\theta_{12})$$

<div style="text-align: center">

δ is the same (up to $\pm$) in each case

</div>

Optimal Parameterization

Location of the phase?

Conventional:

$$U_{23}(\theta_{23})U_{13}(\theta_{13}, \delta)U_{12}(\theta_{12}) \qquad \checkmark$$

Sometimes useful when dealing with matter effect:

$$U_{23}(\theta_{23}, \delta)U_{13}(\theta_{13})U_{12}(\theta_{12})$$

δ is the same (up to $\pm$) in each case

Repeated rotations
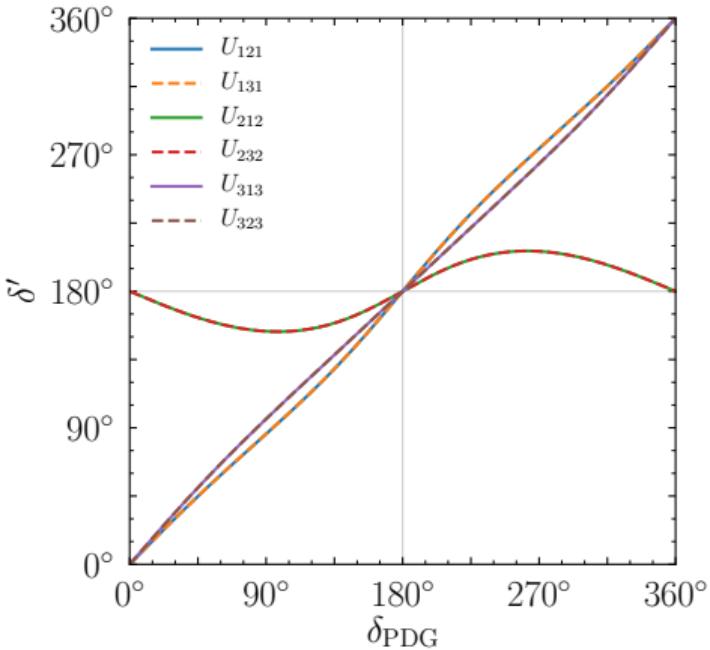
	U_{121}	U_{131}	U_{212}	U_{232}	U_{313}	U_{323}
$|U_{e2}|$	✓	✓	✓	✓	✗	✗
$|U_{e3}|$	✓	✓	✗	✗	✓	✓
$|U_{\mu3}|$	✗	✗	✓	✓	✓	✓

Note that $e^{i\delta}$ must be on first or third rotation

Allowed δ' range
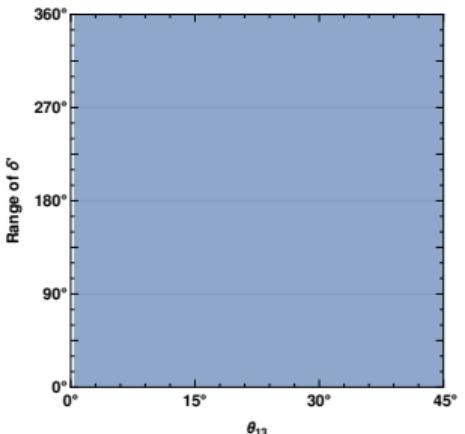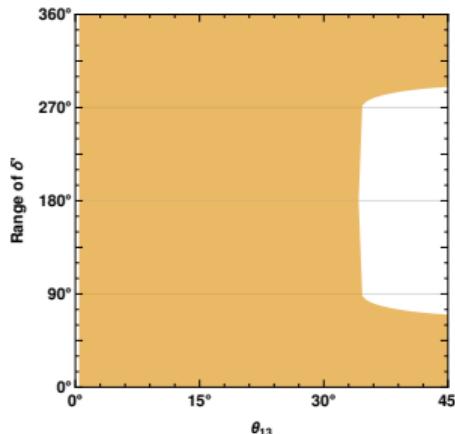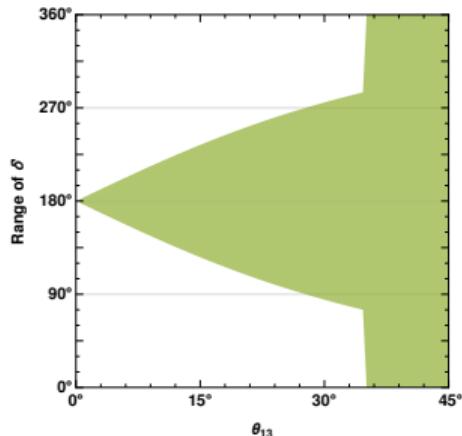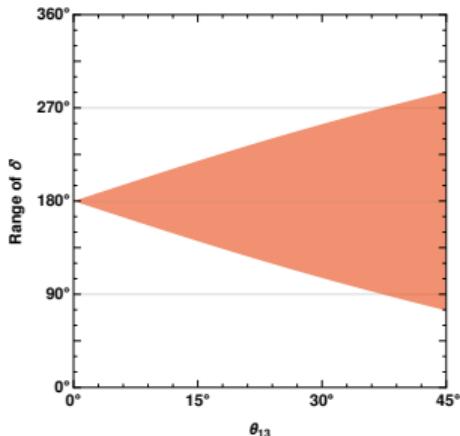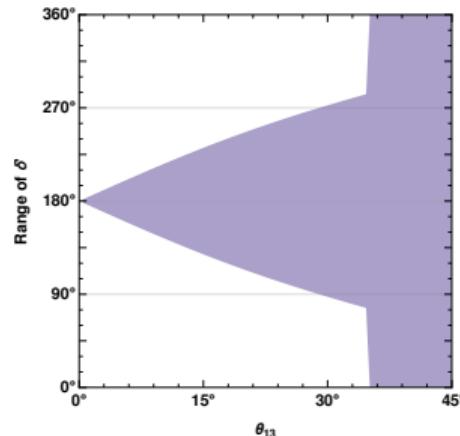
Mass ordering?

Measuring the mass ordering is important in of itself

Phenomenological implications:

- ▶ Affects cosmology
- ▶ Affects $0\nu\beta\beta$
- ▶ Affects end point measurements
- ▶ Affects CνB

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The NOvA+T2K issue is *slightly* resolved by swapping the mass ordering

1. NOvA and T2K both prefer NO over IO
2. NOvA+T2K prefers IO over NO
3. SK still prefers NO over IO
4. NOvA+T2K+SK still prefers NO over IO
5. Daya Bay & RENO provide some information

K. Kelly, et al. 2007.08526

I. Esteban, et al. 2007.14792

PBD, J. Gehrlein, R. Pestes 2008.01110

Is the mass ordering robust?

In the presence of new physics such as NSI we have:

$$[\text{NO}] + [\epsilon = 0] \quad \equiv \quad [\text{IO}] + [\epsilon_{ee} = -2]$$

$$[\text{IO}] + [\epsilon = 0] \quad \equiv \quad [\text{NO}] + [\epsilon_{ee} = -2]$$

Equivalences hold even if all oscillation probabilities are *perfectly* measured

<div align="right">

P. Bakhti, Y. Farzan 1403.0744

P. Coloma, T. Schwetz 1604.05772

PBD, S. Parke 2106.12436

</div>

This is known as the **LMA-Dark** solution

<div align="right">

A. de Gouvêa, A. Friedland, H. Murayama hep-ph/0002064

</div>

Is the mass ordering robust?

Need **scattering** to break

CHARM and NuTeV for $M_{Z'} \gtrsim 10$ GeV

PBD, et al 1701.04828

COHERENT for $M_{Z'} \gtrsim 50$ MeV and cosmology for $M_{Z'} \lesssim 5$ MeV

PBD, Y. Farzan, I. Shoemaker 1804.03660

Dresden-II for any mediator mass

PBD, J. Gehrlein 2204.09060

Can still evade with $\epsilon_{\mu\mu} = \epsilon_{\tau\tau} = 2$ or certain u / d combinations

NSI parameters

Many parameters:

- ▶ Neutrino flavor: 3 diagonal + 3 × 2 flavor changing   9
- ▶ Matter fermion: u, d, e: 3   27
- ▶ V vs. A (or L vs. R): 2   54

If SPVAT then 135

Generally leads to $\nu\nu$ interactions in SNe and early universe: $\times 2 \to 270$

- ▶ For oscillations u, d, e doesn't matter (much)
- ▶ Focus on V for propagation effects
- ▶ Since we want CP violation, focus on flavor changing

6 parameters: $|\epsilon_{e\mu}|e^{i\phi_{e\mu}}$ $|\epsilon_{e\tau}|e^{i\phi_{e\tau}}$ $|\epsilon_{\mu\tau}|e^{i\phi_{\mu\tau}}$

Take one of these three at a time

Relate NSI to vacuum parameters

There is a mapping between vacuum parameters with and without NSI that depends on ρ, E:

$$UM^2U^\dagger + A + N = \widetilde{U}\widetilde{M}^2\widetilde{U}^\dagger + A$$

Vacuum   SM   NSI   apparent   SM
               matter matter vacuum matter

Works for off-axis experiments

Estimate size of effect

Ansatz:

- The data is well described by NSI
- NSI mainly modifies δ:

$$P(\epsilon, \delta_{\text{true}}) \approx P(\epsilon = 0, \delta_{\text{meas}})$$
$$\bar{P}(\epsilon, \delta_{\text{true}}) \approx \bar{P}(\epsilon = 0, \delta_{\text{meas}})$$

Leverage approximate expressions for NSI in LBL

T. Kikuchi, H. Minakata, S. Uchinami 0809.3312

Estimate size of effect: NSI phase

Under the ansatz, if $\delta_{\text{NOvA}} \neq \delta_{\text{T2K}}$

$$\sin(\delta_{\text{true}} + \phi_{e\beta}) \approx 0$$

Since $a_{\text{NOvA}} > a_{\text{T2K}}$ and the data suggests $\sin\delta_{\text{T2K}} \lesssim \sin\delta_{\text{NOvA}}$:

$$\cos(\delta_{\text{true}} + \phi_{e\beta}) \approx -1$$

$$\delta_{\text{true}} \approx \delta_{\text{T2K}} \qquad \Rightarrow \qquad \phi_{e\beta} \approx \frac{3}{2}\pi$$

Estimate size of effect: measured phases

$$\sin \delta_{\text{true}} \approx \frac{\sin \delta_{\text{NOvA}} a_{\text{T2K}} - \sin \delta_{\text{T2K}} a_{\text{NOvA}}}{a_{\text{T2K}} - a_{\text{NOvA}}}$$

Since $\sin \delta_{\text{T2K}} \sim -1$ this suggests
$\sin \delta_{\text{true}} < -1$

Alleviated by:

▶ Statistical fluctuations

▶ Relaxing the ansatz that only δ matters

Dashed: SM
Colored: NSI
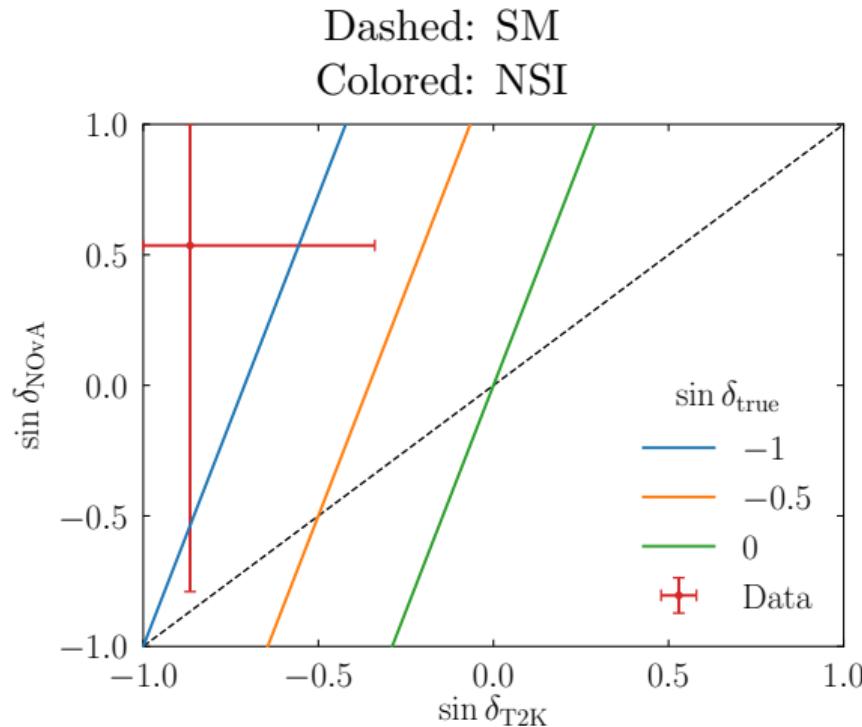
Standard oscillation parameters
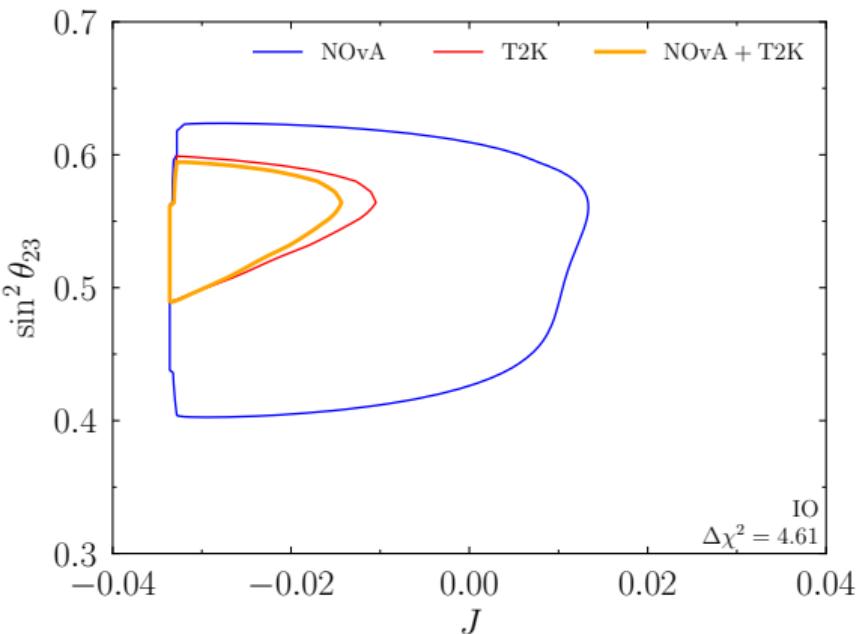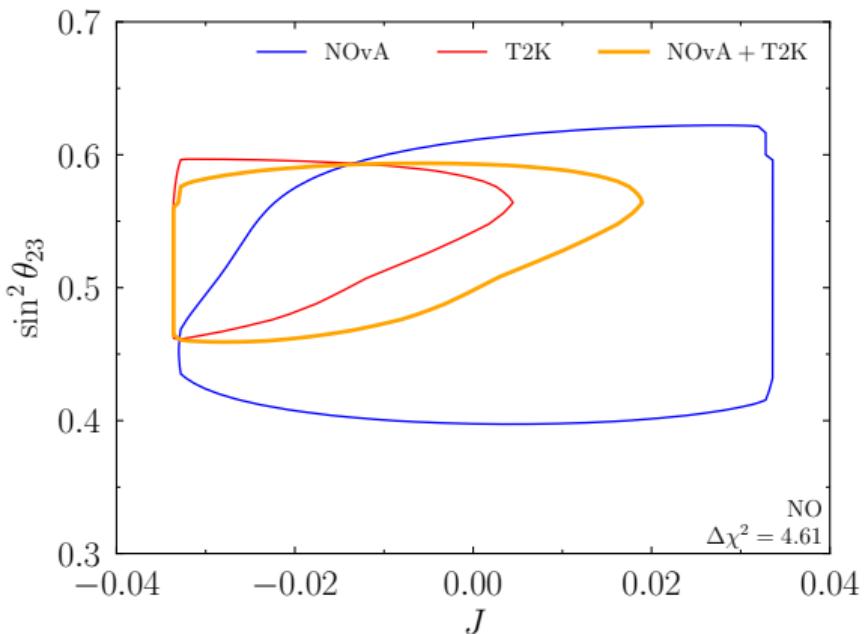
Can see that the combination doesn't like the NO while it does like the IO

IO preferred over NO at $\Delta\chi^2 = 2.3$

NSI parameters

Analytic estimations:

$$|\epsilon_{e\mu}| \approx 0.22 \qquad |\epsilon_{e\tau}| \approx 0.24 \qquad \phi_{e\beta}/\pi \approx 1.5 \qquad \delta/\pi \approx 1.5$$

Numerical fit:

MO	NSI	$|\epsilon_{\alpha\beta}|$	$\phi_{\alpha\beta}/\pi$	δ/π	$\Delta\chi^2$
NO	$\epsilon_{e\mu}$	0.19	1.50	1.46	4.44
	$\epsilon_{e\tau}$	0.28	1.60	1.46	3.65
	$\epsilon_{\mu\tau}$	0.35	0.60	1.83	0.90
IO	$\epsilon_{e\mu}$	0.04	1.50	1.52	0.23
	$\epsilon_{e\tau}$	0.15	1.46	1.59	0.69
	$\epsilon_{\mu\tau}$	0.17	0.14	1.51	1.03

$$\Delta\chi^2 = \chi^2_{\text{SM}} - \chi^2_{\text{NSI}}$$
For the SM: $\chi^2_{\text{NO}} - \chi^2_{\text{IO}} = 2.3$

NSI parameters: IO

NSI parameters: $\epsilon_{\mu\tau}$